

Singular Curves and Giant Magnons

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AdS/CFT and Integrability
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Outline

The Hofman-Maldacena limit

Reality conditions in finite-gap construction

Reality conditions

Elliptic finite-gap solutions

Elliptic solution

Conclusion

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$J_1 \rightarrow \infty$, $E \rightarrow \infty$, $E - J_1 = \text{fixed}$, $\lambda = \text{fixed}$, $J_2 = \text{fixed}$.

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String theory: $(X_0 = \kappa\tau) E = \frac{\sqrt{\lambda}}{2\pi} \int_{-\pi}^{\pi} \kappa d\sigma = \frac{\sqrt{\lambda}}{2\pi} \int_{-\pi\kappa}^{\pi\kappa} \sqrt{\lambda} dx$
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Gauge theory: $\text{tr}(Z^{J_1} W^{J_2})$ becomes **infinitely long**. Relax trace and consider

$$\mathcal{O}_p = \sum_l e^{ipl} (\dots ZZZ W^{J_2} ZZZ \dots).$$

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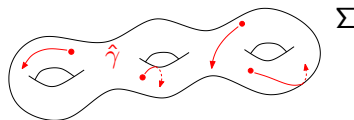
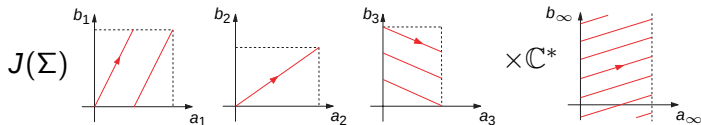
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Finite-Gap picture

The story so far,

$$\left(\begin{array}{l} \text{finite-gap solution to} \\ d * j = 0 \\ dj - j \wedge j = 0 \\ \frac{1}{2} \text{tr} j_{\pm}^2 = -\kappa^2 \end{array} \right)$$

 $\begin{array}{c} 1:1 \\ \longleftrightarrow \end{array}$

 $\vec{\mathcal{A}}$


Real curves

Recall $\Omega(x, \sigma, \tau) = P \overleftarrow{\exp} \int_{[\gamma(\sigma, \tau)]} \frac{1}{1-x^2} (j - x * j)$, hence
 $j \in \mathfrak{su}(2) \Rightarrow j^\dagger = -j \Rightarrow \Omega(x)^\dagger = \Omega(\bar{x})^{-1}$.

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$\Gamma : \det(\tilde{y}\mathbf{1} - \Omega(x)) = 0$ admits an antiholomorphic involution

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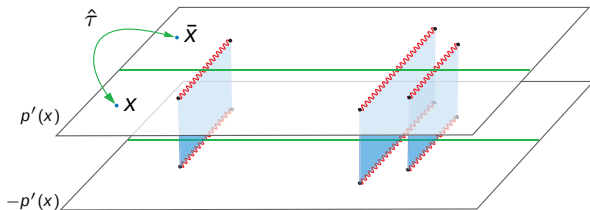
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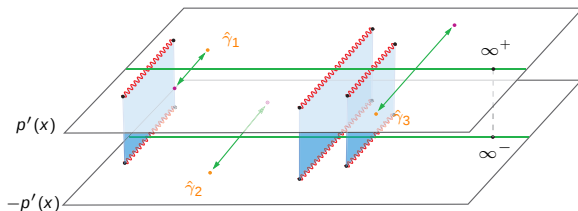
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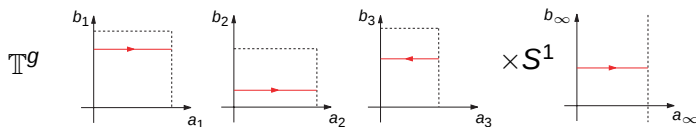
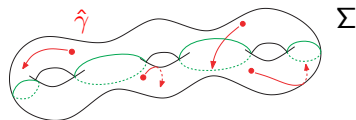
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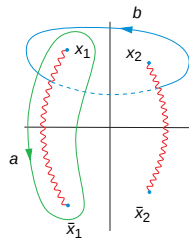
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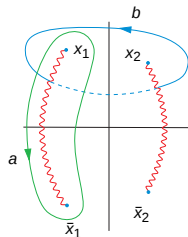


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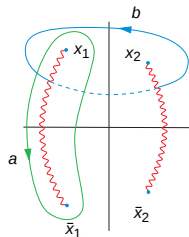


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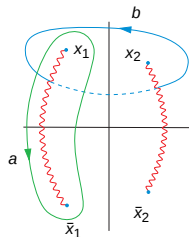
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$$\sigma\text{-periodicity: } \frac{2K(k)\sqrt{1-v^2}}{\kappa'} = \frac{2\pi}{n}, \quad \kappa' \equiv \kappa \frac{|x_1 - \bar{x}_2|}{\sqrt{y_+ y_-}}, \quad v \equiv \frac{y_+ - y_-}{y_+ + y_-}.$$

$k \rightarrow 0$ limit

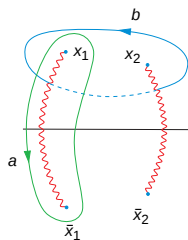
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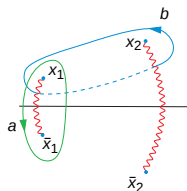
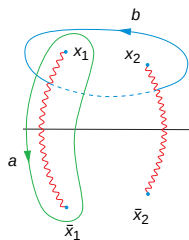
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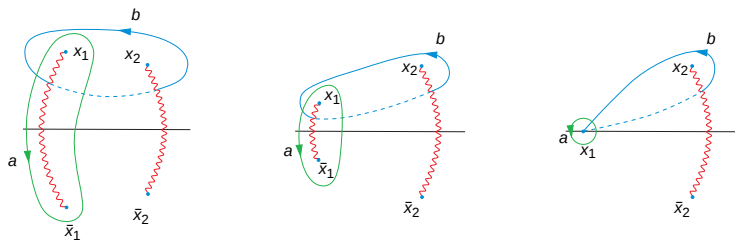
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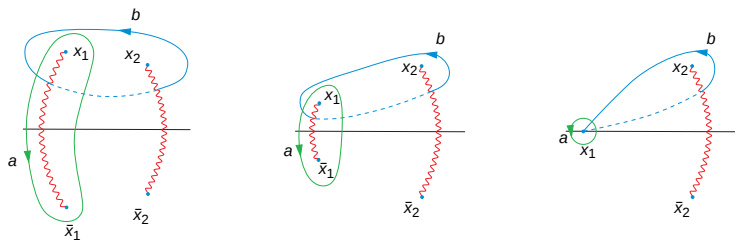
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Remark

- genus reduced by 1.

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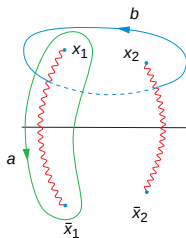
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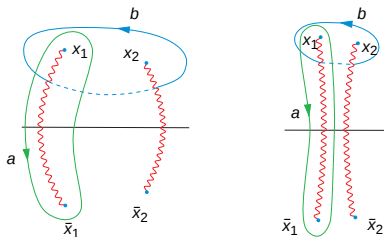
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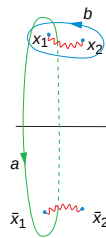
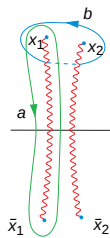
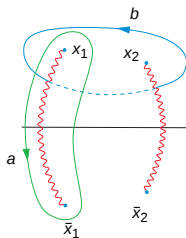
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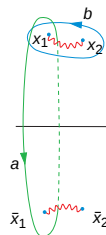
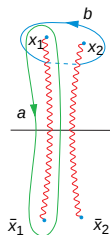
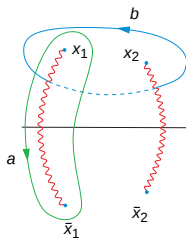
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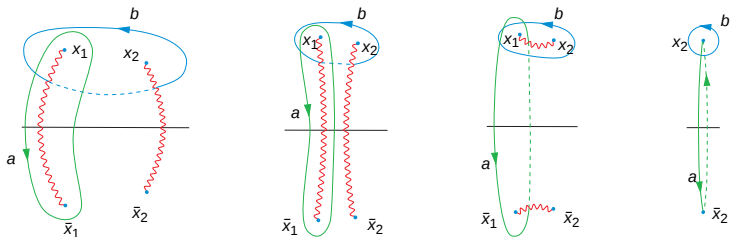
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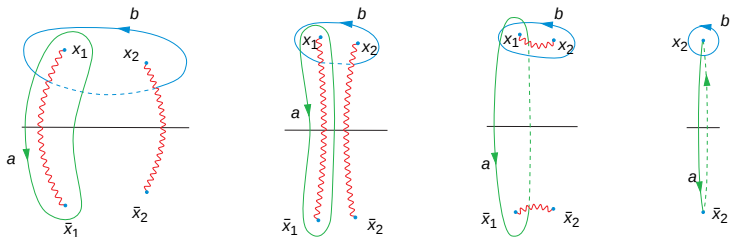
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In particular $E = \sqrt{\lambda} \kappa \rightarrow \infty$ but

$$\left\{ \begin{array}{c} E - J_1 \\ J_2 \end{array} \right\} \rightarrow n \frac{\sqrt{\lambda}}{4\pi} \left| \left(x_2 \mp \frac{1}{x_2} \right) - \left(\bar{x}_2 \mp \frac{1}{\bar{x}_2} \right) \right| < \infty.$$

Condensate 'cuts'

Recall that

$$\int_a dp = 0, \quad \int_b dp = 2\pi n.$$

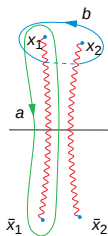
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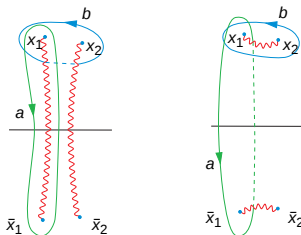
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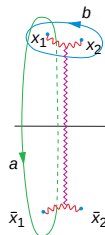
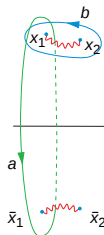
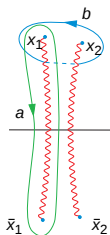
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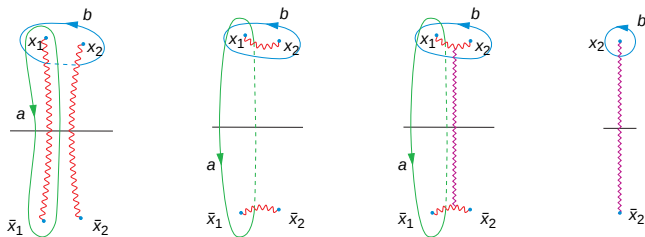
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$$Z_2 = C \frac{\Theta_1(\tilde{X} - i\tilde{\rho}_-)}{\Theta_0(i\tilde{\rho}_-) \Theta_0(\tilde{X})} \exp \left(Z_0(i\tilde{\rho}_-) \tilde{X} + iv_- \tilde{T} + i\varphi_2^0 \right),$$

where $\Theta_\mu(\cdot, k)$, $Z_\mu(\cdot, k)$ are **Jacobi Θ - and ζ -functions**, and

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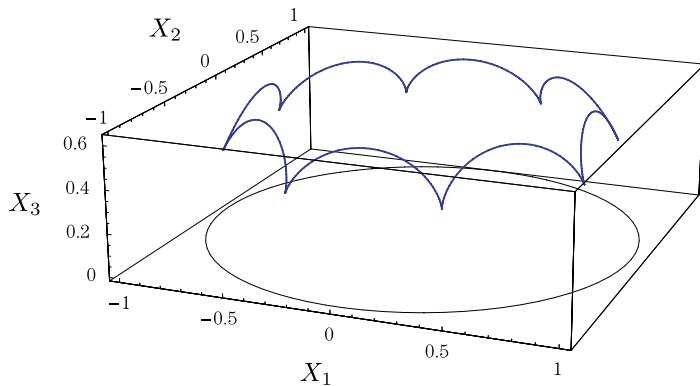
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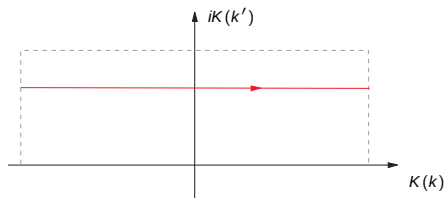
Focus on a **'single hop'** region $\sigma \in [-\frac{\pi}{n}, \frac{\pi}{n}] \Leftrightarrow \tilde{X} \in [-K, K]$.

Profile



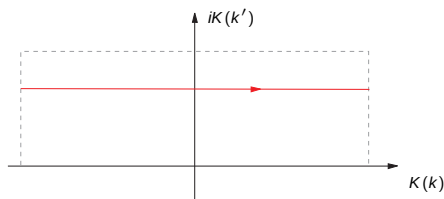
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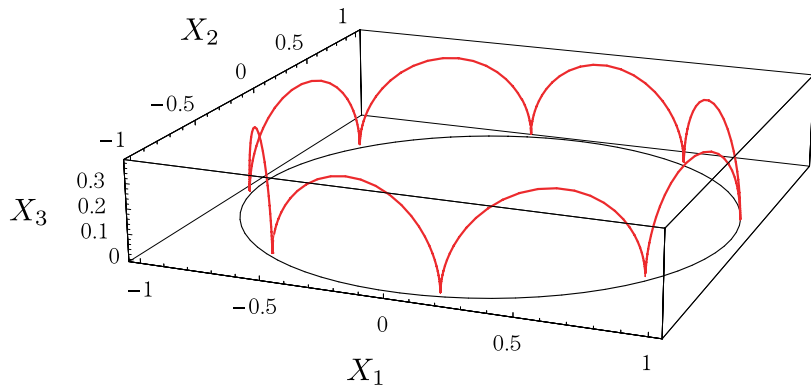
In the limit $k \rightarrow 1$

$$Z_1 = \frac{\cos(\tilde{\rho}_-)}{\cosh(\tilde{X})} \exp\left(iv_+ \tilde{T} + i\varphi_1^0\right),$$

$$Z_2 = \frac{\sinh(\tilde{X} - i\tilde{\rho}_-)}{\cosh(\tilde{X})} \exp\left(i \tan(\tilde{\rho}_-) \tilde{X} + iv_- \tilde{T} + i\varphi_2^0\right).$$

which is known as the **dyonic giant magnon** on S^3 .

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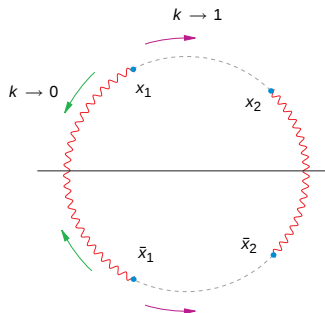
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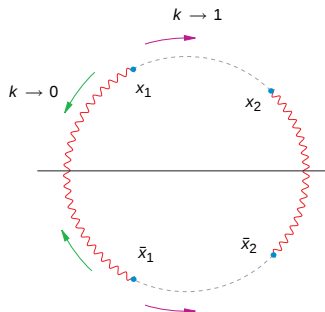
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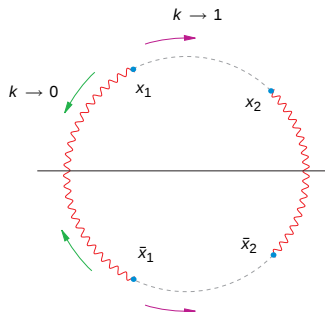
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