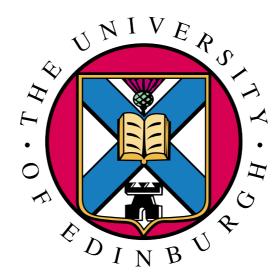
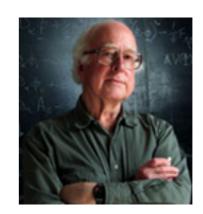


An Erlangen Programme for Supergravity

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References

- arXiv:1511.08737 [hep-th], with Andrea Santi
- arXiv:1511.09264 [hep-th], with Andrea Santi
- arXiv:1605.00881 [hep-th], with Paul de Medeiros + Andrea Santi
- arXiv:1608.05915 [hep-th], with Andrea Santi
- work in progress with Paul de Medeiros + Andrea Santi

Klein's Erlangen Programme

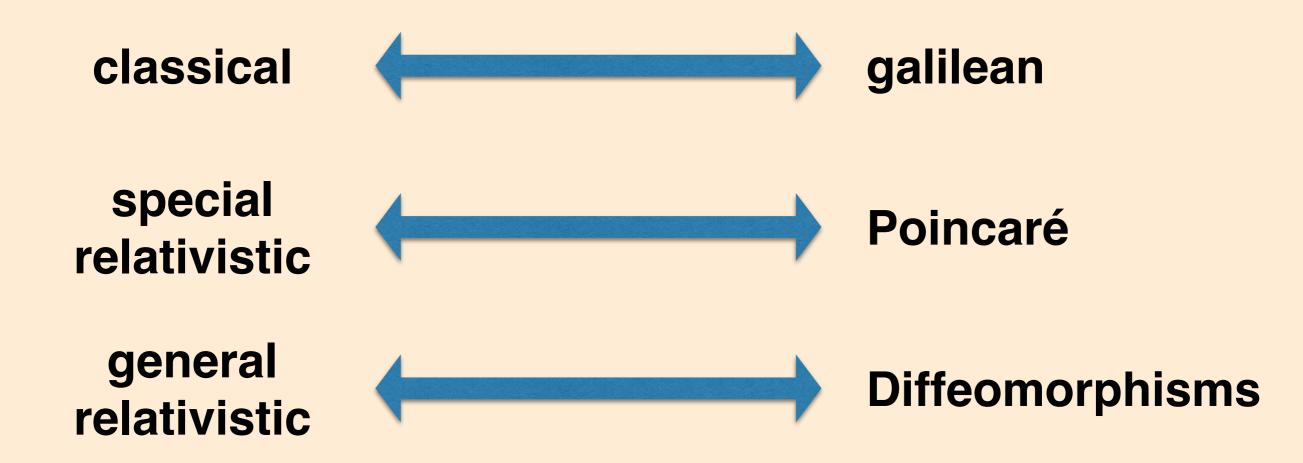


Felix Klein (1849-1925)

"geometry via symmetry"

Physics via symmetry

Physics via the relativity transformation group:



d=11 SUGRA backgrounds

- (M,g) lorentzian, spin, 11-dimensional
- $\$ \to M$ spinor bundle, real, rank 32, symplectic
- $F \in \Omega^4(M) \quad dF = 0$

Natural Problem (v1)

Classify supergravity backgrounds.

Hopelessly hard!

Supersymmetric backgrounds

A connection on spinors

$$D_X = \nabla_X - \frac{1}{24} (X \cdot F - 3F \cdot X)$$

A supersymmetric background admits Killing spinors

$$D\varepsilon = 0$$

$$\nu := \frac{\dim\{\varepsilon \in \$ \mid D\varepsilon = 0\}}{\operatorname{rank} \$} \in \{0, \frac{1}{32}, \frac{1}{16}, \frac{3}{32}, \dots, \frac{31}{32}, 1\}$$

Natural Problem (v2)

Classify supersymmetric backgrounds.

 $\nu > 0$

Still too hard!

Kaluza-Klein supergravity

Lots of work in the late 1970s and early-to-mid 1980s, much of it in Torino, on backgrounds which are metrically a product (e.g., Freund-Rubin, Englert,...).

The main tools are homogeneous geometry and riemannian holonomy.

Closely related to the classification of homogeneous Einstein manifolds.

More recent approaches

- holonomy (of D) \Rightarrow classification for $\nu = 1$
- **spinorial geometry** \Rightarrow rules out $\nu = \frac{31}{32}, \frac{30}{32}$
- **G-structures** \Rightarrow general Ansätze for $\nu = \frac{1}{32}$

• **generalised geometry** \Rightarrow warped products $(\nu \leq \frac{1}{2})$

Natural Problem (v2.5)

Classify supersymmetric backgrounds which are >1/2-BPS

$\nu > \frac{1}{2}$

Still hard, but perhaps doable.

Why >1/2-BPS?

Conjecture (Meessen, '04)

>¹/₂-BPS supergravity backgrounds are homogeneous:

 $G \curvearrowright M$ transitively, preserving g, F

(e.g., M2 shows the conjecture is sharp.)

Theorem (JMF + Hustler, '12)

>¹/₂-BPS supergravity backgrounds are *locally* homogeneous

Key ingredient: the Killing superalgebra

Killing superalgebra

A Lie superalgebra generated by the Killing spinors

 $\mathfrak{k}=\mathfrak{k}_{\bar{0}}\oplus\mathfrak{k}_{\bar{1}}$

 $\mathfrak{k}_{\bar{0}} = \{\xi \in TM \mid \mathcal{L}_{\xi}g = \mathcal{L}_{\xi}F = 0\}$

 $\mathfrak{k}_{\overline{1}} = \{ \varepsilon \in \$ \mid D\varepsilon = 0 \}$

Killing superalgebra

- $[\mathfrak{k}_{\bar{0}},\mathfrak{k}_{\bar{0}}]\subset\mathfrak{k}_{\bar{0}}\quad \text{ Lie bracket of vector fields}$
- $[\mathfrak{k}_{\bar{0}},\mathfrak{k}_{\bar{1}}]\subset\mathfrak{k}_{\bar{1}}\quad \text{Lichnerowicz Lie derivative}$

 $[\mathfrak{k}_{\overline{1}},\mathfrak{k}_{\overline{1}}] \subset \mathfrak{k}_{\overline{0}}$ Dirac current

 $\kappa^a = \bar{\varepsilon} \gamma^a \varepsilon$ is a causal vector field

 $D\varepsilon = 0 \implies \mathcal{L}_{\kappa}g = \mathcal{L}_{\kappa}F = 0$

Local homogeneity

 $\dim \mathfrak{k}_{\bar{1}} > 16$



Every tangent space is spanned by the Dirac currents of Killing spinors.

Killing Lie algebras

Some (2-graded) Lie algebras

 $\mathfrak{so}(9)$ \mathfrak{f}_4 \mathfrak{e}_8

are also generated by (geometric) Killing spinors* on

$$S^7 \qquad S^8 \qquad S^{15}$$

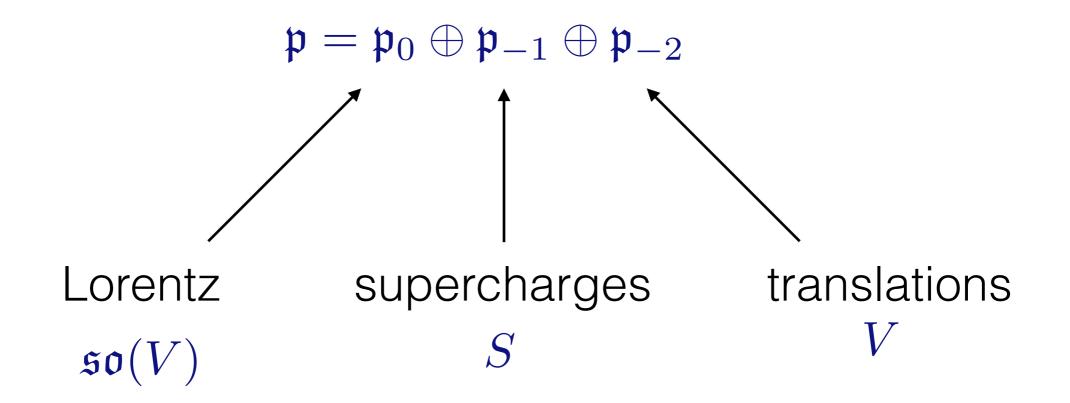
*
$$\nabla_X \varepsilon = \frac{1}{2} X \cdot \varepsilon$$

KSA as a Lie superalgebra

Theorem (JMF + Santi, '16)

The KSA is a filtered Lie superalgebra. It is a *filtered deformation* of a graded subalgebra of the Poincaré superalgebra.

Poincaré superalgebra is graded



 $[\mathfrak{p}_i,\mathfrak{p}_j]=\mathfrak{p}_{i+j}$

 $\mathfrak{a} = \mathfrak{a}_0 \oplus \mathfrak{a}_{-1} \oplus \mathfrak{a}_{-2}$ graded subalgebra of \mathfrak{p}

$$\mathfrak{a}_0 \subset \mathfrak{so}(V) \qquad \mathfrak{a}_{-1} \subset S \qquad \mathfrak{a}_{-2} \subset V$$

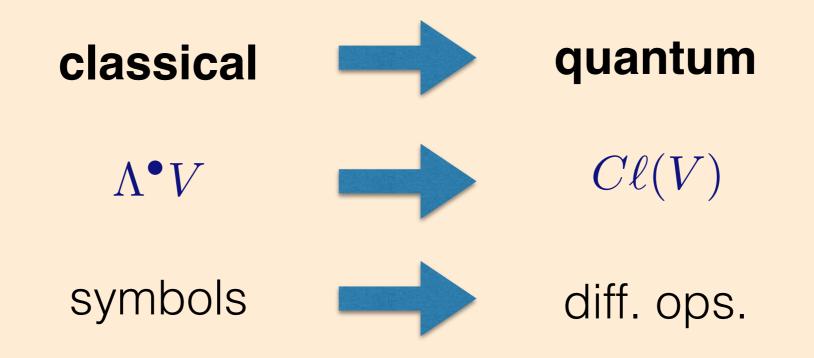
Modify the Lie brackets by terms of positive degree:

$$[\mathfrak{a}_{0},\mathfrak{a}_{0}] \subset \mathfrak{a}_{0}$$
$$[\mathfrak{a}_{0},\mathfrak{a}_{-1}] \subset \mathfrak{a}_{-1}$$
$$[\mathfrak{a}_{0},\mathfrak{a}_{-2}] \subset \mathfrak{a}_{-2} \oplus \mathfrak{a}_{0}$$
$$[\mathfrak{a}_{-1},\mathfrak{a}_{-1}] \subset \mathfrak{a}_{-2} \oplus \mathfrak{a}_{0}$$
$$[\mathfrak{a}_{-1},\mathfrak{a}_{-2}] \subset \mathfrak{a}_{-1}$$
$$[\mathfrak{a}_{-2},\mathfrak{a}_{-2}] \subset \mathfrak{a}_{0} \oplus \mathfrak{a}_{-2}$$

This is the structure of the Killing superalgebra

Filtered deformations

The passage from a *graded* to a *filtered* algebra is very familiar when passing to the very small:



The deformation parameter is \hbar

Filtered deformations

But also when passing to the very large:



The isometry Lie algebra of flat space is graded, but that of a curved manifold is filtered.

The deformation "parameter" is the curvature.

Spencer Cohomology

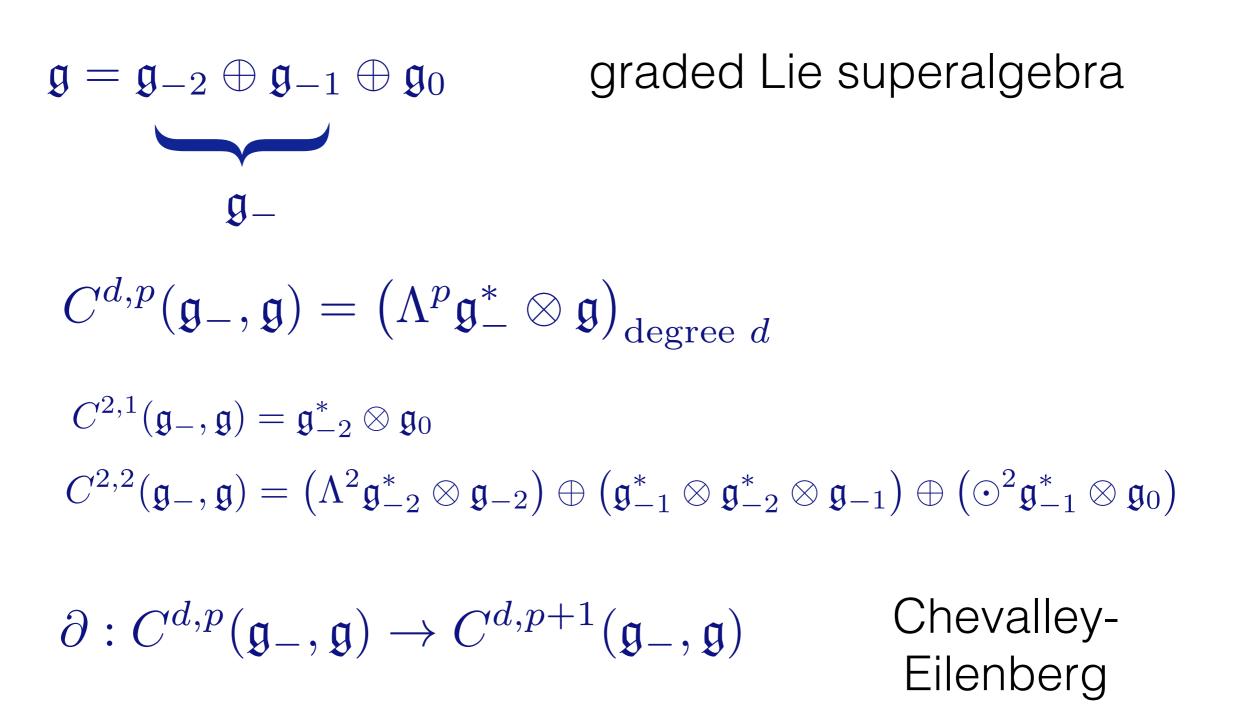
Deformations of algebraic structures are typically governed by a cohomology theory.

e.g., Lie algebra deformations are governed by Chevalley-Eilenberg cohomology.

Filtered deformations of graded Lie superalgebras are governed by generalised Spencer cohomology.

This is a bigraded refinement of Chevalley-Eilenberg cohomology.

Spencer complex



Infinitesimal deformations

Infinitesimal filtered deformations are controlled by $H^{2,2}(\mathfrak{g}_-,\mathfrak{g})$

whereas $H^2(\mathfrak{g},\mathfrak{g})$ controls *all* deformations

cocycles $\alpha + \beta + \gamma$

$$\begin{aligned} \alpha &: \Lambda^2 \mathfrak{g}_{-2} \to \mathfrak{g}_{-2} \\ \beta &: \mathfrak{g}_{-2} \otimes \mathfrak{g}_{-1} \to \mathfrak{g}_{-1} \\ \gamma &: \odot^2 \mathfrak{g}_{-1} \to \mathfrak{g}_0 \end{aligned}$$

Some calculations

 $\mathfrak{p} = V \oplus S \oplus \mathfrak{so}(V)$ d=11 Poincaré superalgebra

 $H^{2,2}(\mathfrak{p}_{-},\mathfrak{p}) = \Lambda^4 V$

The cocycle component $\ \ \beta:V\otimes S \to S$ of a class $F\in \Lambda^4 V$ is given

$$\beta(v,s) = \frac{1}{24}(v \cdot F - 3F \cdot v) \cdot s$$

cf. the connection in the gravitino variation!

p d=4 N=1 Poincaré superalgebra

$$H^{2,2}(\mathfrak{p}_{-},\mathfrak{p}) = \Lambda^0 V \oplus \Lambda^1 V \oplus \Lambda^4 V$$

bosonic field content of "old" minimal off-shell formulation of supergravity

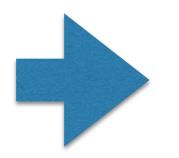
and we again recover the gravitino variation from the cocycle.

Reconstruction

Theorem (JMF + Santi, '16)

Every "admissible" filtered deformation of a graded subalgebra of the Poincaré superalgebra is (contained in) the KSA of a $>\frac{1}{2}$ -BPS supergravity background.

In particular, $>\frac{1}{2}$ -BPS implies the field equations.



Classify >1/2-BPS supergravity backgrounds via their KSAs.

Supergravity backgrounds via their (super)symmetries.

Some further results

- Recovered classification of maximally supersymmetric backgrounds of d=11supergravity.
- Recovered the Festuccia-Seiberg classification of four-dimensional geometries admitting rigid N=1 supersymmetry.
- N=(1,0) d=6 Poincaré: more than supergravity
- N>1 d=4 Poincaré
- Superconformal algebras in various low dimensions

Conclusions

- We have reformulated the classification problem of >½-BPS backgrounds of d=11 supergravity as the classification problem of certain Lie superalgebras ("admissible" filtered deformations of graded subalgebras of the Poincaré superalgebra)
- Generalised Spencer cohomology knows about supergravity and, more generally, about the Killing spinor equations which allow us to define rigidly supersymmetric field theories.