Why I like homogeneous manifolds

José Figueroa-O'Farrill



1 March 2012

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- $H \rightarrow G$ is a principal H-bundle \downarrow M

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- G/H reductive: the sequence splits (as H-modules); i.e.,
 g = ħ ⊕ m with m an Ad(H)-module
- there is a one-to-one correspondence

 $\left\{ \begin{matrix} \mathsf{Ad}(H)\text{-invariant} \\ \text{tensors on } \mathfrak{m} \end{matrix} \right\} \leftrightarrow \left\{ \begin{matrix} H\text{-invariant} \\ \text{tensors on } T_\mathfrak{m}M \end{matrix} \right\} \leftrightarrow \left\{ \begin{matrix} G\text{-invariant} \\ \text{tensor fields on } M \end{matrix} \right\}$

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- our (spatial) universe!

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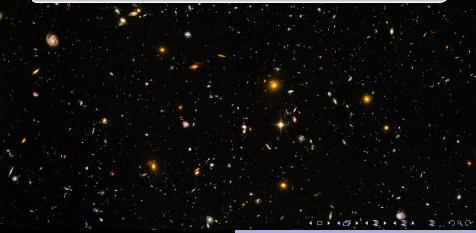
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Why I like homogeneous manifolds

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The cosmological principle

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As far as cosmology is concerned, the spatial universe is a 3-dimensional homogeneous manifold.



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As far as cosmology is concerned, the spatial universe is a 3-dimensional homogeneous manifold.

The big questions

- What (topological) manifold is it? Is it compact? Simply-connected?
- We know it's expanding: but will it do so forever? or will it eventually contract?

. . .

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- riemannian: H compact, so G/H reductive

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CARTAN (1926)

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• if R = 0 then (M, g) is either a Lie group with a bi-invariant metric or the round 7-sphere

CARTAN-SCHOUTEN (1926); WOLF (1971-2)

Examples

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- Fubini–Study metric on complex projective space $\mathbb{C}P^n = U(n+1)/U(n) \times U(1)$

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where $U:S^2\mathfrak{m}\to\mathfrak{m}$ is defined by

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• curvature conditions (e.g., Einstein) \implies algebraic

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But for more than century now, in Physics we work in *lorentzian* signature...

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The birth of spacetime

Hermann Minkowski (1908)

Die Anschauungen über Raum und Zeit, die ich Ihnen entwicklen möchte, sind **auf experimentell-physikalischem Boden erwachsen**. Darin liegt ihre Stärke. Ihre Tendenz ist eine radikale. Von Stund' an sollen Raum für sich und Zeit für sich völlig zu Schatten herabsinken und nur noch eine Art Union der beiden soll Selbständigkeit bewahren.



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The views of space and time that I wish to lay before you have **sprung from the soil of experimental physics**, and therein lies their strength. They are radical. Henceforth space by itself, and time by itself, are doomed to fade away into mere shadows, and only a kind of union of both will retain an independent reality.



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- relavity group = symmetry of trajectories of free particles

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- Minkowski spacetime: a very accurate model of the universe (at some scales)
- It is consistent with quantum theory (RQFT) and is spectacularly successful:

 $\left(\frac{g-2}{2}\right) = \begin{cases} 1\,159\,652\,182.79(7.71)\times 10^{-12} & \text{theory} \\ 1\,159\,652\,180.73(0.28)\times 10^{-12} & \text{experiment} \end{cases}$

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GR mantra

Spacetime tells matter how to move, matter tells spacetime how to curve

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- g, F, ... are subject to Einstein-like equations

• Unique supersymmetric theory in d = 11

NAHM (1979), CREMMER+JULIA+SCHERK (1980)

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- (bosonic) fields: lorentzian metric g, 3-form A

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- Field equations from action (with F = dA)

$$\underbrace{\frac{1}{2}\int R\,d\text{vol}}_{\text{Einstein-Hilbert}} - \underbrace{\frac{1}{4}\int F\wedge \star F}_{\text{Maxwell}} + \underbrace{\frac{1}{12}\int F\wedge F\wedge A}_{\text{Chern-Simons}}$$

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Explicitly,

$$d \star F = \frac{1}{2}F \wedge F$$
$$\operatorname{Ric}(X, Y) = \frac{1}{2} \langle \iota_X F, \iota_Y F \rangle - \frac{1}{6}g(X, Y)|F|^2$$

together with dF = 0

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 - ...
- Many of them are homogeneous or of low cohomogeneity.

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- such spinor fields are called Killing spinors

 Not every manifold admits spinors: so an implicit condition on (M, g, F) is that M should be spin

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- a background is said to be ν -BPS, where $\nu = \frac{n}{32}$

• The Dirac current V of a Killing spinor ε is defined by

 $g(V,X)=(\epsilon,X\cdot\epsilon)$

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into a Lie superalgebra

JMF+MEESSEN+PHILIP (2004)

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 It is called the Killing superalgebra of the supersymmetric background (M, g, F)

Empirical Fact

Every known v-BPS background with $v > \frac{1}{2}$ is homogeneous.

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Homogeneity conjecture

Every *MhdWh* v-BPS background with $v > \frac{1}{2}$ is homogeneous. MEESSEN (2004)

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Homogeneity conjecture

Every $M/M/\nu$ -BPS background with $\nu > \frac{1}{2}$ is homogeneous. MEESSEN (2004)

Theorem

Every $\nu\text{-BPS}$ background of eleven-dimensional supergravity with $\nu>\frac{3}{4}$ is homogeneous.

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What we proved is that the ideal $[\mathfrak{g}_1, \mathfrak{g}_1]$ of \mathfrak{g}_0 generated by the Killing spinors spans the tangent space to every point of M: local homogeneity

Generalisations

Theorem

Every v-BPS background of type IIB supergravity with $v > \frac{3}{4}$ is homogeneous. Every v-BPS background of type I and heterotic supergravities with $v > \frac{1}{2}$ is homogeneous. JMF+Hackett-Jones+Moutsopoulos (2007)

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The theorems actually suggest a stronger version of the conjecture: that the symmetries which are generated from the supersymmetries already act (locally) transitively.

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What good is it?

If the homogeneity conjecture were true, then classifying homogeneous supergravity backgrounds would also classify v-BPS backgrounds for $v > \frac{1}{2}$.

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 the supergravity field equations for homogeneous backgrounds are algebraic and hence simpler to solve than PDEs

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This would be good because

- the supergravity field equations for homogeneous backgrounds are algebraic and hence simpler to solve than PDEs
- we have learnt **a lot** (about string theory) from supersymmetric supergravity backgrounds, so their classification could teach us even more

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A homogeneous eleven-dimensional supergravity background is described algebraically by the data $(\mathfrak{g}, \mathfrak{h}, \gamma, \phi)$, where

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- γ is an \mathfrak{h} -invariant lorentzian inner product on \mathfrak{m}
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subject to some algebraic equations which are given purely in terms of the structure constants of \mathfrak{g} (and \mathfrak{h}).

Skip technical details

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Choose a basis X_a for \mathfrak{h} and a basis Y_i for \mathfrak{m} .

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The metric and 4-forms are described by $\mathfrak{h}\text{-invariant}$ tensors γ_{ij} and $\phi_{ijkl}.$

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$$\begin{split} & [X_{a}, X_{b}] = f_{ab}{}^{c}X_{c} \\ & [X_{a}, Y_{i}] = f_{ai}{}^{j}Y_{j} + f_{ai}{}^{b}X_{b} \\ & [Y_{i}, Y_{j}] = f_{ij}{}^{a}X_{a} + f_{ij}{}^{k}Y_{k} \end{split}$$

If M is reductive, then $f_{\alpha i}{}^b = 0$. We will assume this in what follows.

The metric and 4-forms are described by $\mathfrak{h}\text{-invariant}$ tensors γ_{ij} and $\phi_{ijkl}.$

We raise and lower indices with γ_{ij} .

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Homogeneous Hodge/de Rham calculus

The G-invariant differential forms in M = G/H form a subcomplex of the de Rham complex:

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the codifferential is given by

$$\begin{split} (\delta\phi)_{ijk} = &-\tfrac{3}{2} f_{m[i}{}^n \phi^m{}_{jk]n} - 3 U_{m[i}{}^n \phi^m{}_{jk]n} - U_m{}^{mn} \phi_{nijk} \end{split}$$
 where $U_{ijk} = f_{i(jk)}$

Homogeneous Ricci curvature

Finally, the Ricci tensor for a homogeneous (reductive) manifold is given by

$$\begin{split} R_{ij} &= -\frac{1}{2} f_i{}^{k\ell} f_{jk\ell} - \frac{1}{2} f_{ik}{}^{\ell} f_{j\ell}{}^k + \frac{1}{2} f_{ik}{}^a f_{aj}{}^k \\ &+ \frac{1}{2} f_{jk}{}^a f_{ai}{}^k - \frac{1}{2} f_{k\ell}{}^{\ell} f^k{}_{ij} - \frac{1}{2} f_{k\ell}{}^{\ell} f^k{}_{ji} + \frac{1}{4} f_{k\ell i} f^{k\ell}{}_j \end{split}$$

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It is now a matter of assembling these ingredients to write down the supergravity field equations in a homogeneous Ansatz.

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Classifying homogeneous supergravity backgrounds of a certain type involves now the following steps:

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Classifying homogeneous supergravity backgrounds of a certain type involves now the following steps:

- Classify the desired homogeneous geometries
- For each such geometry parametrise the space of invariant lorentzian metrics $(\gamma_1, \gamma_2, ...)$ and invariant closed 4-forms $(\phi_1, \phi_2, ...)$
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- Solve the equations!

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Definition

The action of G on M is **proper** if the map $G \times M \to M \times M$, $(\gamma, m) \mapsto (\gamma \cdot m, m)$ is proper. In particular, proper actions have compact stabilisers.

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What if the action is not proper?

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Theorem (Kowalsky, 1996)

If a simple Lie group acts transitively and non-properly on a lorentzian manifold (M, g), then (M, g) is locally isometric to (anti) de Sitter spacetime.

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This means that we need only classify Lie subalgebras corresponding to *compact* Lie subgroups!

Some recent classification results

 Symmetric eleven-dimensional supergravity backgrounds JMF (2011)

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- Symmetric type IIB supergravity backgrounds JMF+HustLer (IN PREPARATION)

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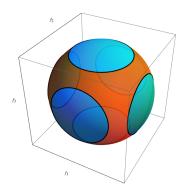
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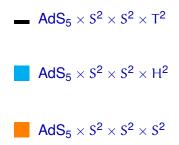
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- Homogeneous M2-duals: $\mathfrak{g} = \mathfrak{so}(3,2) \oplus \mathfrak{so}(N)$ for N > 4JMF+Ungureanu (in preparation)

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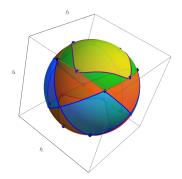
A moduli space of AdS₅ symmetric backgrounds





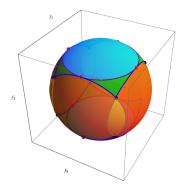
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A moduli space of AdS₃ symmetric backgrounds



- $AdS_3 \times S^2 \times T^6$
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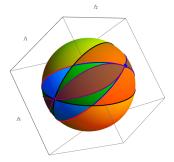
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- $AdS_2 \times S^2 \times T^7$
- $AdS_2 \times S^5 \times T^4$
- $AdS_2 \times T^5 \times S^2 \times S^2$
- $AdS_2 \times S^5 \times H^2 \times T^2$
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And one final gratuitous pretty picture



- $\bullet \quad AdS_2 \times S^3 \times T^6$
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Thank you for your attention

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