

Outline

Based on the following papers:

1511.08737	(CMP)	w/ Andrea Santi
1511.09264	(JPA)	
1608.05915	(ATMP)	
1605.00881	(JHEP)	w/ Paul de Medeiros & Andrea Santi
180? —		

1. The Spencer cohomology of the Poincaré superalgebra
2. Killing superalgebras
3. Field equations on homogeneous Lorentzian manifolds

1. The Spencer cohomology of the Poincaré superalgebra

The ur-example of a Lorentzian manifold is (4-dim) Minkowski spacetime. In D dimensions it is a real affine space modelled on a vector space $V \cong \mathbb{R}^D$ with an indefinite inner product η with signature $(1, D-1)$ ("mostly minus").

The Lie algebra of isometries of Minkowski spacetime is the Poincaré Lie algebra $\mathfrak{p} \stackrel{\text{def}}{=} \underline{\mathfrak{so}}(V, \eta) \oplus V$. It is a \mathbb{Z} -graded Lie algebra:

$$[A, B] = A \cdot B - B \cdot A$$

$$[A, v] = Av$$

$$\forall A, B \in \underline{\mathfrak{so}}(V) = \text{End}(V)$$

$$[v, w] = 0$$

$$v, w \in V$$

Associated to (V, η) we have the Clifford algebra $\mathcal{Cl}(V, \eta)$, generated by V subject to $v^2 = -\eta(v, v) \mathbf{1} \quad \forall v \in V$.

The precise structure of $\mathcal{Cl}(V, \eta)$ depends on $\dim V$. If $\dim V=4$, $\mathcal{Cl}(V, \eta) \cong \text{Mat}(4, \mathbb{R})$ and if $\dim V=11$, $\mathcal{Cl}(V, \eta) \cong 2\text{Mat}(32, \mathbb{R})$. Let S be an irreducible real module of $\mathcal{Cl}(V, \eta)$. For example, if $\dim V=4$, $S \cong \mathbb{R}^4$ and if $\dim V=11$, $S \cong \mathbb{R}^{32}$ and we will choose the module on which the volume form in $\mathcal{Cl}(V, \eta)$ acts nontrivially. In both of these cases, Clifford action $V \times S \xrightarrow{\text{cl}} S$ is skewsymmetric relative to a $\underline{\mathfrak{so}}(V)$ -invariant symplectic structure $(-, -)$ on S . The transpose of cl rel. to $(-, -)$ and η defines a symmetric bilinear map $K: S \times S \rightarrow V$ known as the Dirac current. Explicitly, if $s \in S$, $K(s) \in V$ is defined by

$$\eta(K(s), v) = (s, v \cdot s) \quad \forall v \in V.$$

$$\text{Then } K(s_1, s_2) = \frac{1}{2} (K(s_1 + s_2) - K(s_1) - K(s_2)).$$

The action of $\underline{\mathfrak{so}}(V)$ on S is given by the embedding $\begin{aligned} \underline{\mathfrak{so}}(V) &\hookrightarrow \mathcal{Cl}(V) \\ v \wedge w &\mapsto \frac{1}{4} [v, w] \end{aligned}$

where $v \wedge w \in \underline{\mathfrak{so}}(V)$ is defined by $(v \wedge w)(x) = \eta(v \otimes x)w - \eta(w \otimes x)v$

The Poincaré superalgebra

$\overset{\text{so}}{\oplus}(\mathbb{V}) \oplus S \oplus V$ with brackets extending those of the Poincaré algebra above by
 $s \cdot s = 0$. It is a \mathbb{Z} -graded superalgebra $\overset{\text{so}}{\oplus}(\mathbb{V}) \oplus S \oplus V$

$$[A, s] = As$$

$$[s_1, s_2] = K(s_1, s_2)$$

S is a module over the subalgebra \mathbb{E}_- . Therefore we can consider the complex $\mathcal{J}: C^*(\mathbb{E}_-; S) \rightarrow C^{*+1}(\mathbb{E}_-; S)$

in the strict sense

where $C^*(\mathbb{E}_-; S) = \text{Hom}(\Lambda^p \mathbb{E}_-, S)$. Because S is \mathbb{Z} -graded, the differential has $\deg \partial = 0$ and the complex admits a second grading so we can talk about $\mathcal{J}: C^{p,d} \rightarrow C^{p+1,d}$ where $C^{p,d} \subset C^p$ consists of those p -chains of degree d . Also because S is \mathbb{Z} -graded, $H^{p,d}$ is an $\overset{\text{so}}{\oplus}(\mathbb{V}) = S_0$ -module. We are interested mostly in $H^{2,2}$, where $C^{2,2} = \{ \alpha: \Lambda^2 V \rightarrow V \} \oplus \{ \beta: V \times S \rightarrow S \} \oplus \{ \gamma: S \times S \rightarrow \overset{\text{so}}{\oplus}(\mathbb{V}) \}$

e.g. $\dim V = 11$,

$$H^{2,2} \cong \Lambda^4 V \quad \text{as } \overset{\text{so}}{\oplus}(\mathbb{V})\text{-modules}$$

$$\dim V = 9$$

$$H^{2,2} \cong \Lambda^0 V \oplus \Lambda^1 V \oplus \Lambda^2 V \quad \text{if } (\alpha, A, b\omega) \in \Lambda^0 V \otimes V \otimes \Lambda^2 V, \quad \beta_{a,A,b}(v, s) = -v \cdot (a + b\omega) \cdot s + (A \wedge v) \cdot \text{rot. } s - 2\eta(A, v) \text{ vols}$$

Killing superalgebras ($D=11$ from now on)

Now let (M, g) be an 11-dim'l Lorentzian spin manifold and $F \in \Omega^1(M)$. Let $\$ \rightarrow M$ be the spinor bundle (actually a bundle of $(\mathfrak{sl}(TM, g))$ -modules) and define a connection D on $\$$ as follows:

$$D_X \varepsilon = \nabla_X \varepsilon + \beta_F(x, \varepsilon) \xrightarrow{\text{spin connection}} \frac{1}{24} (X^b F \cdot \varepsilon - 3F X^b \varepsilon)$$

We say $\varepsilon \in \Gamma(\$)$ is a Killing spinor if $D\varepsilon = 0$

If $\varepsilon \in \Gamma(\$)$ is a Killing spinor, then $K(\varepsilon)$ is

a Killing vector field & if $dF = 0$, also $\mathcal{L}_{K(\varepsilon)} F = 0$.

(From now on assume $dF = 0$.)

Let $k_{\bar{0}} = \{ X \in \mathfrak{X}(M) \mid \mathcal{L}_X g = 0 \text{ & } \mathcal{L}_X F = 0 \}$ and $k_{\bar{1}} = \{ \varepsilon \in \Gamma(\$) \mid D\varepsilon = 0 \}$.

Then (F+Maessen+Philip, '04) $k = k_{\bar{0}} \oplus k_{\bar{1}}$ is a Lie superalgebra called the Killing superalgebra $\mathfrak{so}_{10}(M, g, F)$.

Conjecture (Meessen, '04) If $\dim k_T > \frac{1}{2} \text{rank } \mathbb{F}$ ($= 16$) then \exists a lie group G acting transitively on M isometrically and preserving F .

Theorem (F+Hustler, '12) If $\dim k_T > \frac{1}{2} \text{rank } \mathbb{F}$, then the ideal $[k_T, k_T] \subset k_{\bar{T}}$ acts locally transitively on M :

$$\forall p \in M, \quad \text{span} \left\{ \xi(p) \mid \xi \in [k_T, k_T] \right\} = T_p M$$

Theorem (F+Sauli, '16) The Killing superalgebra $k = k_{\bar{0}} \oplus k_{\bar{1}}$ of (M, g, F) is a filtered lie superalgebra and its associated graded algebra is a \mathbb{Z} -graded subalgebra of \mathfrak{F} .

c.f. (M, g) riemannian and $\mathfrak{g} = \{ \xi \in \mathfrak{X}(M) \mid \nabla_{\xi}^{\bar{0}} g = 0 \}$ the lie algebra of isometries.

Killing transport (Kostant, Garoch) $\Rightarrow \mathfrak{g}$ is filtered and $\text{gr } \mathfrak{g}$ is a \mathbb{Z} -graded subalgebra of the euclidean lie algebra

3. Lorentzian field equations on homogeneous lorentzian manifolds

i.e.: isometry lie algebra
of the flat model

Theorem (F+Sauli, '16) If $\dim k_T > \frac{1}{2} \text{rank } \mathbb{F}$, then $g \& F$ obey the (bosonic) field equations of $d=11$ supergravity:

$$d * F = -\frac{1}{2} F \wedge F \quad \text{and} \quad \text{Ric}(g) = T(F, g)$$

"Maxwell" "Einstein"

Furthermore, we can reconstruct (M, g, F) up to local 'isometry' from k_T .

This 'reduces' the local classification of bosonic supergravity backgrounds preserving $> \frac{1}{2}$ of the supersymmetry to the classification of certain filtered deformations of \mathbb{Z} -graded subalgebras of \mathfrak{F} .

State of the art Let $n = \dim k_T$. Then $n \in [0, 32] \cap \mathbb{Z}$.

$n=32$ classification (F+Papadopoulos '01) recovered algebraically (F+Sauli, '15)

$n=31$ \nexists (Gran+Gutowski+Papadopoulos+Roest '07, F+Gratia '07)

$n=30$ \nexists (Gran+Gutowski+Papadopoulos '10)

$27 \leq n \leq 29$?

$n=26, 24, 23, 20, 18$ \exists but no classification

$n=17, 19, 21, 23, 25$?

4. Other dimensions

One can play this game in any dimension and with other (\mathbb{Z} -graded) lie superalgebras:

- (extended) Poincaré superalgebra
- "non-relativistic" superalgebras
- conformal superalgebras
- ...

We have already analysed this problem in $d=4$ and $d=6$.

Theorem (de Medeiros + F + Santi '16)

Let $\hat{\mathfrak{s}} = \mathfrak{so}(V) \oplus S \oplus V$ be the 4d $N=1$ Poincaré superalgebra. Then $H^{2,2}(\hat{\mathfrak{s}}_-; \hat{\mathfrak{s}}) \cong \overset{a}{\wedge^0} V \oplus \overset{\varphi}{\wedge^1} V \oplus \overset{b}{\wedge^2} V$

The corresponding Killing spinor equation is

$$D_X \varepsilon = \nabla_X \varepsilon - X \cdot (a+b \text{vol}) \cdot \varepsilon + (\varphi \cdot X) \cdot \text{vol} \cdot \varepsilon - 2g(X, Y) \text{vol} \cdot \varepsilon$$

and Killing spinors generate a lie superalgebra.
 ↴ minimal (4d) off-shell SUGRA

The geometries for which D is flat are:

- ① $a=b=\varphi=0 \Rightarrow$ Minkowski spacetime
- ② $\varphi=0, a^2+b^2>0 \Rightarrow AdS_4 \& Ric = -12(a^2+b^2)g$
- ③ $a=b=0, \varphi \neq 0$
 - (a) φ timelike $\Rightarrow R \times S^3$
 - (b) φ spacelike $\Rightarrow AdS_3 \times R$
 - (c) φ lightlike \Rightarrow Nappi-Witten ppwave

Theorem (de Medeiros + F + Santi, '18)

Let $\mathfrak{s} = \mathfrak{so}(V) \oplus S \oplus V$ be the $(1,0)$ $d=6$ Poincaré superalgebra, and let $\hat{\mathfrak{s}} = (\mathfrak{so}(V) \oplus \mathfrak{r}) \oplus S \oplus V$ be the extension by R-symmetry.

Here $S \otimes \mathbb{C} = \sum_{\pm} \otimes_{\mathfrak{r}} \Delta$
 ↪ fundamental rep of $Sp(1)$ R-symmetry
 half-spinor rep of $Spin(V) \cong SL(2, \mathbb{H})$

Then $H^{2,2}(\hat{\mathfrak{s}}_-; \hat{\mathfrak{s}})$ and $H^{2,2}(\hat{\mathfrak{s}}_-; \hat{\mathfrak{s}})$ are $\mathfrak{so}(V) \otimes \mathfrak{sp}(1)$ -mods

$$\text{Then } H^{2,2}(\hat{\mathfrak{s}}_-; \hat{\mathfrak{s}}) \cong (\wedge^0 V \otimes \wedge^2 \Delta) \oplus (V \otimes \text{Sym}^2 \Delta) \quad \text{and} \quad H^{2,2}(\hat{\mathfrak{s}}_-; \hat{\mathfrak{s}}) \cong (\wedge^0 V \otimes \wedge^2 \Delta) \oplus (V \otimes \text{Sym}^2 \Delta)$$

↑ self-dual 3-form ← difference → general 3-form

Killing spinor equation:

$$D_X \varepsilon := \nabla_X \varepsilon - t_X H \cdot \varepsilon + 3 \varphi(X) \cdot \varepsilon - X \cdot \varphi \cdot \varepsilon \quad (\text{this goes beyond SUGRA})$$

Hamburg talk (summary)

- Minkowski spacetime \rightarrow Poincaré algebra \rightarrow Poincaré superalgebra ($d=4$, $d=11$)
- Spencer cohomology \rightarrow $\beta: V \otimes S \rightarrow S$ ($d=4$, $d=11$)
- Killing spinors \rightarrow Killing superalgebra \rightarrow Homogeneity \rightarrow Algebraic structure
- $> \frac{1}{2}$ -BPS \Rightarrow field equations
- Applications to classification of $> \frac{1}{2}$ -BPS sugra backgrounds
- Application to constructing rigidly supersymmetric field theories