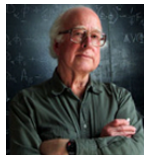


Homogeneous Lorentzian Manifolds in Supergravity

José Miguel Figueroa O'Farrill



HelgaFest
Greifswald, 18 March 2014



¡Feliz cumpleaños!

¡Feliz cumpleaños!

*Liebe Helga,
Alles Gute zum Geburtstag!*

Outline

1 Supergravity

Outline

- 1 Supergravity
- 2 A geometrical interlude

Outline

- 1 Supergravity
- 2 A geometrical interlude
- 3 Homogeneity in Supergravity

- 1 Supergravity
- 2 A geometrical interlude
- 3 Homogeneity in Supergravity

Supergravity

- result of ongoing effort to marry GR and quantum theory

Supergravity

- result of ongoing effort to marry GR and quantum theory
- four-dimensional supergravity “discovered” in 1975

FERRARA+FREEDMAN+VAN NIEUWENHUIZEN, DESER+ZUMINO

Supergravity

- result of ongoing effort to marry GR and quantum theory
- four-dimensional supergravity “discovered” in 1975
FERRARA+FREEDMAN+VAN NIEUWENHUIZEN, DESER+ZUMINO
- many more supergravity theories, painstakingly constructed in the 1970s and 1980s

Supergravity

- result of ongoing effort to marry GR and quantum theory
- four-dimensional supergravity “discovered” in 1975
FERRARA+FREEDMAN+VAN NIEUWENHUIZEN, DESER+ZUMINO
- many more supergravity theories, painstakingly constructed in the 1970s and 1980s
- “crown jewels of mathematical physics”

Supergravity

- result of ongoing effort to marry GR and quantum theory
- four-dimensional supergravity “discovered” in 1975
FERRARA+FREEDMAN+VAN NIEUWENHUIZEN, DESER+ZUMINO
- many more supergravity theories, painstakingly constructed in the 1970s and 1980s
- “crown jewels of mathematical physics”
- the formalism could use some improvement!

Supergravity

- result of ongoing effort to marry GR and quantum theory
- four-dimensional supergravity “discovered” in 1975
FERRARA+FREEDMAN+VAN NIEUWENHUIZEN, DESER+ZUMINO
- many more supergravity theories, painstakingly constructed in the 1970s and 1980s
- “crown jewels of mathematical physics”
- the formalism could use some improvement!
- The geometric set-up:

Supergravity

- result of ongoing effort to marry GR and quantum theory
- four-dimensional supergravity “discovered” in 1975
FERRARA+FREEDMAN+VAN NIEUWENHUIZEN, DESER+ZUMINO
- many more supergravity theories, painstakingly constructed in the 1970s and 1980s
- “crown jewels of mathematical physics”
- the formalism could use some improvement!
- The geometric set-up:
 - (M, g) a lorentzian, spin manifold of dimension ≤ 11

Supergravity

- result of ongoing effort to marry GR and quantum theory
- four-dimensional supergravity “discovered” in 1975
FERRARA+FREEDMAN+VAN NIEUWENHUIZEN, DESER+ZUMINO
- many more supergravity theories, painstakingly constructed in the 1970s and 1980s
- “crown jewels of mathematical physics”
- the formalism could use some improvement!
- The geometric set-up:
 - (M, g) a lorentzian, spin manifold of dimension ≤ 11
 - some extra geometric data, e.g., differential forms F, \dots

Supergravity

- result of ongoing effort to marry GR and quantum theory
- four-dimensional supergravity “discovered” in 1975
FERRARA+FREEDMAN+VAN NIEUWENHUIZEN, DESER+ZUMINO
- many more supergravity theories, painstakingly constructed in the 1970s and 1980s
- “crown jewels of mathematical physics”
- the formalism could use some improvement!
- The geometric set-up:
 - (M, g) a lorentzian, spin manifold of dimension ≤ 11
 - some extra geometric data, e.g., differential forms F, \dots
 - a connection $D = \nabla + \dots$ on the spinor (actually Clifford) bundle S

Poincaré supergravities

	32			24	20	16	12	8	4
11	M								
10	IIA	IIB				I			
9	$N = 2$					$N = 1$			
8	$N = 2$					$N = 1$			
7	$N = 4$					$N = 2$			
6	(2,2)	(3,1)	(4,0)	(2,1)	(3,0)	(1,1)	(2,0)	(1,0)	
5		$N = 8$		$N = 6$		$N = 4$		$N = 2$	
4		$N = 8$		$N = 6$	$N = 5$	$N = 4$	$N = 3$	$N = 2$	$N = 1$

Poincaré supergravities

	32		24	20	16	12	8	4
11	M							
10	IIA	IIB			I			
9	$N = 2$				$N = 1$			
8	$N = 2$				$N = 1$			
7	$N = 4$				$N = 2$			
6	$(2, 2)$	$(3, 1)$ $(4, 0)$	$(2, 1)$ $(3, 0)$		$(1, 1)$ $(2, 0)$		$(1, 0)$	
5	$N = 8$		$N = 6$		$N = 4$		$N = 2$	
4	$N = 8$		$N = 6$	$N = 5$	$N = 4$	$N = 3$	$N = 2$	$N = 1$

Eleven-dimensional supergravity

- Unique supersymmetric theory in $d = 11$

NAHM (1979), CREMMER+JULIA+SCHERK (1980)

Eleven-dimensional supergravity

- Unique supersymmetric theory in $d = 11$
NAHM (1979), CREMMER+JULIA+SCHERK (1980)
- (bosonic) fields: lorentzian metric g , 3-form A

Eleven-dimensional supergravity

- Unique supersymmetric theory in $d = 11$

NAHM (1979), CREMMER+JULIA+SCHERK (1980)

- (bosonic) fields: lorentzian metric g , 3-form A
- Field equations from action (with $F = dA$)

$$\underbrace{\frac{1}{2} \int R \, d\text{vol}}_{\text{Einstein-Hilbert}} - \underbrace{\frac{1}{4} \int F \wedge \star F}_{\text{Maxwell}} + \underbrace{\frac{1}{12} \int F \wedge F \wedge A}_{\text{Chern-Simons}}$$

Eleven-dimensional supergravity

- Unique supersymmetric theory in $d = 11$

NAHM (1979), CREMMER+JULIA+SCHERK (1980)

- (bosonic) fields: lorentzian metric g , 3-form A
- Field equations from action (with $F = dA$)

$$\underbrace{\frac{1}{2} \int R \, d\text{vol}}_{\text{Einstein-Hilbert}} - \underbrace{\frac{1}{4} \int F \wedge \star F}_{\text{Maxwell}} + \underbrace{\frac{1}{12} \int F \wedge F \wedge A}_{\text{Chern-Simons}}$$

- Explicitly,

$$d \star F = \frac{1}{2} F \wedge F$$

$$\text{Ric}(X, Y) = \frac{1}{2} \langle \iota_X F, \iota_Y F \rangle - \frac{1}{6} g(X, Y) |F|^2$$

together with $dF = 0$

Supergravity backgrounds

- A triple (M, g, F) where $dF = 0$ and (g, F) satisfying the above PDEs is called an **(eleven-dimensional) supergravity background**.

Supergravity backgrounds

- A triple (M, g, F) where $dF = 0$ and (g, F) satisfying the above PDEs is called an **(eleven-dimensional) supergravity background**.
- There is by now a huge catalogue of eleven-dimensional supergravity backgrounds:

Supergravity backgrounds

- A triple (M, g, F) where $dF = 0$ and (g, F) satisfying the above PDEs is called an **(eleven-dimensional) supergravity background**.
- There is by now a huge catalogue of eleven-dimensional supergravity backgrounds:
 - Freund–Rubin: $AdS_4 \times X^7$, $AdS_7 \times X^4$, ...

Supergravity backgrounds

- A triple (M, g, F) where $dF = 0$ and (g, F) satisfying the above PDEs is called an **(eleven-dimensional) supergravity background**.
- There is by now a huge catalogue of eleven-dimensional supergravity backgrounds:
 - Freund–Rubin: $AdS_4 \times X^7$, $AdS_7 \times X^4, \dots$
 - pp-waves

Supergravity backgrounds

- A triple (M, g, F) where $dF = 0$ and (g, F) satisfying the above PDEs is called an **(eleven-dimensional) supergravity background**.
- There is by now a huge catalogue of eleven-dimensional supergravity backgrounds:
 - Freund–Rubin: $AdS_4 \times X^7$, $AdS_7 \times X^4$, ...
 - pp-waves
 - branes: elementary, intersecting, overlapping, wrapped, ...

Supergravity backgrounds

- A triple (M, g, F) where $dF = 0$ and (g, F) satisfying the above PDEs is called an **(eleven-dimensional) supergravity background**.
- There is by now a huge catalogue of eleven-dimensional supergravity backgrounds:
 - Freund–Rubin: $AdS_4 \times X^7$, $AdS_7 \times X^4$,...
 - pp-waves
 - branes: elementary, intersecting, overlapping, wrapped,...
 - Kaluza–Klein monopoles,...

Supergravity backgrounds

- A triple (M, g, F) where $dF = 0$ and (g, F) satisfying the above PDEs is called an **(eleven-dimensional) supergravity background**.
- There is by now a huge catalogue of eleven-dimensional supergravity backgrounds:
 - Freund–Rubin: $AdS_4 \times X^7$, $AdS_7 \times X^4$, ...
 - pp-waves
 - branes: elementary, intersecting, overlapping, wrapped, ...
 - Kaluza–Klein monopoles, ...
 - ...

Supergravity backgrounds

- A triple (M, g, F) where $dF = 0$ and (g, F) satisfying the above PDEs is called an **(eleven-dimensional) supergravity background**.
- There is by now a huge catalogue of eleven-dimensional supergravity backgrounds:
 - Freund–Rubin: $AdS_4 \times X^7$, $AdS_7 \times X^4$, ...
 - pp-waves
 - branes: elementary, intersecting, overlapping, wrapped, ...
 - Kaluza–Klein monopoles, ...
 - ...
- It is convenient to organise this information according to how much “supersymmetry” the background preserves.

Supersymmetry

- Eleven-dimensional supergravity has local supersymmetry

Supersymmetry

- Eleven-dimensional supergravity has local supersymmetry
- manifests itself as a connection \mathbb{D} on the spinor bundle S

Supersymmetry

- Eleven-dimensional supergravity has local supersymmetry
- manifests itself as a connection \mathbb{D} on the spinor bundle S
- \mathbb{D} is **not** induced from a connection on the spin bundle:
 $\text{hol}(\mathbb{D}) \subset \mathfrak{sl}(32, \mathbb{R})$

Supersymmetry

- Eleven-dimensional supergravity has local supersymmetry
- manifests itself as a connection \mathbf{D} on the spinor bundle S
- \mathbf{D} is **not** induced from a connection on the spin bundle:
 $\text{hol}(\mathbf{D}) \subset \mathfrak{sl}(32, \mathbb{R})$
- the field equations are encoded in the curvature of \mathbf{D} :

$$\sum_i e^i \cdot R^{\mathbf{D}}(e_i, -) = 0 \quad \text{in} \quad \Omega^1(M; \text{End } S)$$

Supersymmetry

- Eleven-dimensional supergravity has local supersymmetry
- manifests itself as a connection \mathbf{D} on the spinor bundle S
- \mathbf{D} is **not** induced from a connection on the spin bundle:
 $\text{hol}(\mathbf{D}) \subset \mathfrak{sl}(32, \mathbb{R})$
- the field equations are encoded in the curvature of \mathbf{D} :

$$\sum_i e^i \cdot R^{\mathbf{D}}(e_i, -) = 0 \quad \text{in} \quad \Omega^1(M; \text{End } S)$$

- geometric analogies (riemannian):

Supersymmetry

- Eleven-dimensional supergravity has local supersymmetry
- manifests itself as a connection \mathbb{D} on the spinor bundle S
- \mathbb{D} is **not** induced from a connection on the spin bundle:
 $\text{hol}(\mathbb{D}) \subset \mathfrak{sl}(32, \mathbb{R})$
- the field equations are encoded in the curvature of \mathbb{D} :

$$\sum_i e^i \cdot R^{\mathbb{D}}(e_i, -) = 0 \quad \text{in} \quad \Omega^1(M; \text{End } S)$$

- geometric analogies (riemannian):
 - $\nabla \varepsilon = 0 \implies \text{Ric} = 0$

Supersymmetry

- Eleven-dimensional supergravity has local supersymmetry
- manifests itself as a connection \mathbb{D} on the spinor bundle S
- \mathbb{D} is **not** induced from a connection on the spin bundle:
 $\text{hol}(\mathbb{D}) \subset \mathfrak{sl}(32, \mathbb{R})$
- the field equations are encoded in the curvature of \mathbb{D} :

$$\sum_i e^i \cdot R^{\mathbb{D}}(e_i, -) = 0 \quad \text{in} \quad \Omega^1(M; \text{End } S)$$

- geometric analogies (riemannian):
 - $\nabla \varepsilon = 0 \implies \text{Ric} = 0$
 - $\nabla_X \varepsilon = \frac{1}{2} X \cdot \varepsilon \implies \text{Einstein}$

Supersymmetry

- Eleven-dimensional supergravity has local supersymmetry
- manifests itself as a connection D on the spinor bundle S
- D is **not** induced from a connection on the spin bundle:
 $\text{hol}(D) \subset \mathfrak{sl}(32, \mathbb{R})$
- the field equations are encoded in the curvature of D :

$$\sum_i e^i \cdot R^D(e_i, -) = 0 \quad \text{in} \quad \Omega^1(M; \text{End } S)$$

- geometric analogies (riemannian):
 - $\nabla \varepsilon = 0 \implies \text{Ric} = 0$
 - $\nabla_X \varepsilon = \frac{1}{2} X \cdot \varepsilon \implies \text{Einstein}$
- a background (M, g, F) is **supersymmetric** if there exists a nonzero spinor field ε satisfying $D\varepsilon = 0$

Supersymmetry

- Eleven-dimensional supergravity has local supersymmetry
- manifests itself as a connection \mathbb{D} on the spinor bundle S
- \mathbb{D} is **not** induced from a connection on the spin bundle:
 $\text{hol}(\mathbb{D}) \subset \mathfrak{sl}(32, \mathbb{R})$
- the field equations are encoded in the curvature of \mathbb{D} :

$$\sum_i e^i \cdot R^{\mathbb{D}}(e_i, -) = 0 \quad \text{in} \quad \Omega^1(M; \text{End } S)$$

- geometric analogies (riemannian):
 - $\nabla \varepsilon = 0 \implies \text{Ric} = 0$
 - $\nabla_X \varepsilon = \frac{1}{2} X \cdot \varepsilon \implies \text{Einstein}$
- a background (M, g, F) is **supersymmetric** if there exists a nonzero spinor field ε satisfying $\mathbb{D}\varepsilon = 0$
- such spinor fields are called **Killing spinors**

Killing spinors

- Not every manifold admits spinors: so an implicit condition on (M, g, F) is that M should be **spin**

Killing spinors

- Not every manifold admits spinors: so an implicit condition on (M, g, F) is that M should be **spin**
- The spinor bundle of an eleven-dimensional lorentzian spin manifold is a real 32-dimensional symplectic vector bundle

Killing spinors

- Not every manifold admits spinors: so an implicit condition on (M, g, F) is that M should be **spin**
- The spinor bundle of an eleven-dimensional lorentzian spin manifold is a real 32-dimensional symplectic vector bundle
- The Killing spinor equation is

$$D_X \varepsilon = \nabla_X \varepsilon + \frac{1}{12} (X^b \wedge F) \cdot \varepsilon + \frac{1}{6} \iota_X F \cdot \varepsilon = 0$$

which is a linear, first-order PDE:

Killing spinors

- Not every manifold admits spinors: so an implicit condition on (M, g, F) is that M should be **spin**
- The spinor bundle of an eleven-dimensional lorentzian spin manifold is a real 32-dimensional symplectic vector bundle
- The Killing spinor equation is

$$D_X \varepsilon = \nabla_X \varepsilon + \frac{1}{12} (X^b \wedge F) \cdot \varepsilon + \frac{1}{6} \iota_X F \cdot \varepsilon = 0$$

which is a linear, first-order PDE:

- linearity: solutions form a vector space

Killing spinors

- Not every manifold admits spinors: so an implicit condition on (M, g, F) is that M should be **spin**
- The spinor bundle of an eleven-dimensional lorentzian spin manifold is a real 32-dimensional symplectic vector bundle
- The Killing spinor equation is

$$D_X \varepsilon = \nabla_X \varepsilon + \frac{1}{12} (X^b \wedge F) \cdot \varepsilon + \frac{1}{6} \iota_X F \cdot \varepsilon = 0$$

which is a linear, first-order PDE:

- linearity: solutions form a vector space
- first-order: solutions determined by their values at any point

Killing spinors

- Not every manifold admits spinors: so an implicit condition on (M, g, F) is that M should be **spin**
- The spinor bundle of an eleven-dimensional lorentzian spin manifold is a real 32-dimensional symplectic vector bundle
- The Killing spinor equation is

$$D_X \varepsilon = \nabla_X \varepsilon + \frac{1}{12} (X^b \wedge F) \cdot \varepsilon + \frac{1}{6} \iota_X F \cdot \varepsilon = 0$$

which is a linear, first-order PDE:

- linearity: solutions form a vector space
- first-order: solutions determined by their values at any point
- the dimension of the space of Killing spinors is $0 \leq n \leq 32$

Killing spinors

- Not every manifold admits spinors: so an implicit condition on (M, g, F) is that M should be **spin**
- The spinor bundle of an eleven-dimensional lorentzian spin manifold is a real 32-dimensional symplectic vector bundle
- The Killing spinor equation is

$$D_X \varepsilon = \nabla_X \varepsilon + \frac{1}{12}(X^b \wedge F) \cdot \varepsilon + \frac{1}{6} \iota_X F \cdot \varepsilon = 0$$

which is a linear, first-order PDE:

- linearity: solutions form a vector space
- first-order: solutions determined by their values at any point
- the dimension of the space of Killing spinors is $0 \leq n \leq 32$
- a background is said to be **v-BPS** if $n = 32v$

Which values of ν are known to appear?

- $\nu = 1$ backgrounds are classified

JMF+PAPADOPOULOS (2002)

Which values of ν are known to appear?

- $\nu = 1$ backgrounds are classified

JMF+PAPADOPOULOS (2002)

- $\nu = \frac{31}{32}$ has been ruled out

GRAN+GUTOWSKI+PAPADOPOULOS+ROEST (2006)

JMF+GADHIA (2007)

Which values of ν are known to appear?

- $\nu = 1$ backgrounds are classified
JMF+PAPADOPOULOS (2002)
- $\nu = \frac{31}{32}$ has been ruled out
GRAN+GUTOWSKI+PAPADOPOULOS+ROEST (2006)
JMF+GADHIA (2007)
- $\nu = \frac{15}{16}$ has been ruled out
GRAN+GUTOWSKI+PAPADOPOULOS (2010)

Which values of ν are known to appear?

- $\nu = 1$ backgrounds are classified
JMF+PAPADOPOULOS (2002)
- $\nu = \frac{31}{32}$ has been ruled out
GRAN+GUTOWSKI+PAPADOPOULOS+ROEST (2006)
JMF+GADHIA (2007)
- $\nu = \frac{15}{16}$ has been ruled out
GRAN+GUTOWSKI+PAPADOPOULOS (2010)
- No other values of ν have been ruled out

Which values of ν are known to appear?

- $\nu = 1$ backgrounds are classified

JMF+PAPADOPOULOS (2002)

- $\nu = \frac{31}{32}$ has been ruled out

GRAN+GUTOWSKI+PAPADOPOULOS+ROEST (2006)

JMF+GADHIA (2007)

- $\nu = \frac{15}{16}$ has been ruled out

GRAN+GUTOWSKI+PAPADOPOULOS (2010)

- No other values of ν have been ruled out

- The following values are known to appear:

$$0, \frac{1}{32}, \frac{1}{16}, \frac{3}{32}, \frac{1}{8}, \frac{5}{32}, \frac{3}{16}, \dots, \frac{1}{4}, \dots, \frac{3}{8}, \dots, \frac{1}{2}, \\ \dots, \frac{9}{16}, \dots, \frac{5}{8}, \dots, \frac{11}{16}, \dots, \frac{3}{4}, \dots, \frac{13}{16}, \dots, 1$$

Which values of ν are known to appear?

- $\nu = 1$ backgrounds are classified

JMF+PAPADOPOULOS (2002)

- $\nu = \frac{31}{32}$ has been ruled out

GRAN+GUTOWSKI+PAPADOPOULOS+ROEST (2006)

JMF+GADHIA (2007)

- $\nu = \frac{15}{16}$ has been ruled out

GRAN+GUTOWSKI+PAPADOPOULOS (2010)

- No other values of ν have been ruled out

- The following values are known to appear:

$$0, \frac{1}{32}, \frac{1}{16}, \frac{3}{32}, \frac{1}{8}, \frac{5}{32}, \frac{3}{16}, \dots, \frac{1}{4}, \dots, \frac{3}{8}, \dots, \frac{1}{2},$$

$$\dots, \frac{9}{16}, \dots, \frac{5}{8}, \dots, \frac{11}{16}, \dots, \frac{3}{4}, \dots, \frac{13}{16}, \dots, 1$$

- those in the 2nd row are now known to be **homogeneous!**

Supersymmetries generate isometries

- The **Dirac current** V_ε of a Killing spinor ε is defined by

$$g(V_\varepsilon, X) = (\varepsilon, X \cdot \varepsilon)$$

Supersymmetries generate isometries

- The **Dirac current** V_ε of a Killing spinor ε is defined by

$$g(V_\varepsilon, X) = (\varepsilon, X \cdot \varepsilon)$$

- More generally, if $\varepsilon_1, \varepsilon_2$ are Killing spinors,

$$g(V_{\varepsilon_1, \varepsilon_2}, X) = (\varepsilon_1, X \cdot \varepsilon_2)$$

Supersymmetries generate isometries

- The **Dirac current** V_ε of a Killing spinor ε is defined by

$$g(V_\varepsilon, X) = (\varepsilon, X \cdot \varepsilon)$$

- More generally, if $\varepsilon_1, \varepsilon_2$ are Killing spinors,

$$g(V_{\varepsilon_1, \varepsilon_2}, X) = (\varepsilon_1, X \cdot \varepsilon_2)$$

- $V := V_\varepsilon$ is **causal**: $g(V, V) \leq 0$

Supersymmetries generate isometries

- The **Dirac current** V_ε of a Killing spinor ε is defined by

$$g(V_\varepsilon, X) = (\varepsilon, X \cdot \varepsilon)$$

- More generally, if $\varepsilon_1, \varepsilon_2$ are Killing spinors,

$$g(V_{\varepsilon_1, \varepsilon_2}, X) = (\varepsilon_1, X \cdot \varepsilon_2)$$

- $V := V_\varepsilon$ is **causal**: $g(V, V) \leq 0$
- V is Killing: $\mathcal{L}_V g = 0$

Supersymmetries generate isometries

- The **Dirac current** V_ε of a Killing spinor ε is defined by

$$g(V_\varepsilon, X) = (\varepsilon, X \cdot \varepsilon)$$

- More generally, if $\varepsilon_1, \varepsilon_2$ are Killing spinors,

$$g(V_{\varepsilon_1, \varepsilon_2}, X) = (\varepsilon_1, X \cdot \varepsilon_2)$$

- $V := V_\varepsilon$ is **causal**: $g(V, V) \leq 0$
- V is Killing: $\mathcal{L}_V g = 0$
- $\mathcal{L}_V F = 0$

GAUNTLETT+PAKIS (2002)

Supersymmetries generate isometries

- The **Dirac current** V_ε of a Killing spinor ε is defined by

$$g(V_\varepsilon, X) = (\varepsilon, X \cdot \varepsilon)$$

- More generally, if $\varepsilon_1, \varepsilon_2$ are Killing spinors,

$$g(V_{\varepsilon_1, \varepsilon_2}, X) = (\varepsilon_1, X \cdot \varepsilon_2)$$

- $V := V_\varepsilon$ is **causal**: $g(V, V) \leq 0$
- V is Killing: $\mathcal{L}_V g = 0$
- $\mathcal{L}_V F = 0$
- $\mathcal{L}_V D = 0$

GAUNTLETT+PAKIS (2002)

Supersymmetries generate isometries

- The **Dirac current** V_ε of a Killing spinor ε is defined by

$$g(V_\varepsilon, X) = (\varepsilon, X \cdot \varepsilon)$$

- More generally, if $\varepsilon_1, \varepsilon_2$ are Killing spinors,

$$g(V_{\varepsilon_1, \varepsilon_2}, X) = (\varepsilon_1, X \cdot \varepsilon_2)$$

- $V := V_\varepsilon$ is **causal**: $g(V, V) \leq 0$

- V is Killing: $\mathcal{L}_V g = 0$

- $\mathcal{L}_V F = 0$

GAUNTLETT+PAKIS (2002)

- $\mathcal{L}_V D = 0$

- ε' Killing spinor \implies so is $\mathcal{L}_V \varepsilon' = \nabla_V \varepsilon' - \rho(\nabla V) \varepsilon'$

Supersymmetries generate isometries

- The **Dirac current** V_ε of a Killing spinor ε is defined by

$$g(V_\varepsilon, X) = (\varepsilon, X \cdot \varepsilon)$$

- More generally, if $\varepsilon_1, \varepsilon_2$ are Killing spinors,

$$g(V_{\varepsilon_1, \varepsilon_2}, X) = (\varepsilon_1, X \cdot \varepsilon_2)$$

- $V := V_\varepsilon$ is **causal**: $g(V, V) \leq 0$

- V is Killing: $\mathcal{L}_V g = 0$

- $\mathcal{L}_V F = 0$

GAUNTLETT+PAKIS (2002)

- $\mathcal{L}_V D = 0$

- ε' Killing spinor \implies so is $\mathcal{L}_V \varepsilon' = \nabla_V \varepsilon' - \rho(\nabla V) \varepsilon'$

- $\mathcal{L}_V \varepsilon = 0$

JMF+MEESSEN+PHILIP (2004)

The Killing superalgebra

- The **symmetry superalgebra** of a supersymmetric background (M, g, F) : $\mathfrak{g} = \mathfrak{g}_0 \oplus \mathfrak{g}_1$, where

The Killing superalgebra

- The **symmetry superalgebra** of a supersymmetric background (M, g, F) : $\mathfrak{g} = \mathfrak{g}_0 \oplus \mathfrak{g}_1$, where
 - \mathfrak{g}_0 is the space of F -preserving Killing vector fields, and

The Killing superalgebra

- The **symmetry superalgebra** of a supersymmetric background (M, g, F) : $\mathfrak{g} = \mathfrak{g}_0 \oplus \mathfrak{g}_1$, where
 - \mathfrak{g}_0 is the space of F -preserving Killing vector fields, and
 - \mathfrak{g}_1 is the space of Killing spinors

JMF+MEESSEN+PHILIP (2004)

The Killing superalgebra

- The **symmetry superalgebra** of a supersymmetric background (M, g, F) : $\mathfrak{g} = \mathfrak{g}_0 \oplus \mathfrak{g}_1$, where
 - \mathfrak{g}_0 is the space of F -preserving Killing vector fields, and
 - \mathfrak{g}_1 is the space of Killing spinors

JMF+MEESSEN+PHILIP (2004)

- The ideal $\mathfrak{k} = [\mathfrak{g}_1, \mathfrak{g}_1] \oplus \mathfrak{g}_1$ generated by \mathfrak{g}_1 is called the **Killing superalgebra**

The Killing superalgebra

- The **symmetry superalgebra** of a supersymmetric background (M, g, F) : $\mathfrak{g} = \mathfrak{g}_0 \oplus \mathfrak{g}_1$, where
 - \mathfrak{g}_0 is the space of F -preserving Killing vector fields, and
 - \mathfrak{g}_1 is the space of Killing spinors

JMF+MEESSEN+PHILIP (2004)

- The ideal $\mathfrak{k} = [\mathfrak{g}_1, \mathfrak{g}_1] \oplus \mathfrak{g}_1$ generated by \mathfrak{g}_1 is called the **Killing superalgebra**
- It behaves as expected: it deforms under geometric limits (e.g., Penrose) and it embeds under asymptotic limits.

The Killing superalgebra

- The **symmetry superalgebra** of a supersymmetric background (M, g, F) : $\mathfrak{g} = \mathfrak{g}_0 \oplus \mathfrak{g}_1$, where
 - \mathfrak{g}_0 is the space of F -preserving Killing vector fields, and
 - \mathfrak{g}_1 is the space of Killing spinors

JMF+MEESSEN+PHILIP (2004)

- The ideal $\mathfrak{k} = [\mathfrak{g}_1, \mathfrak{g}_1] \oplus \mathfrak{g}_1$ generated by \mathfrak{g}_1 is called the **Killing superalgebra**
- It behaves as expected: it deforms under geometric limits (e.g., Penrose) and it embeds under asymptotic limits.
- It is a very useful invariant of a supersymmetric supergravity background

The Killing superalgebra

- The **symmetry superalgebra** of a supersymmetric background (M, g, F) : $\mathfrak{g} = \mathfrak{g}_0 \oplus \mathfrak{g}_1$, where
 - \mathfrak{g}_0 is the space of F -preserving Killing vector fields, and
 - \mathfrak{g}_1 is the space of Killing spinors

JMF+MEESSEN+PHILIP (2004)

- The ideal $\mathfrak{k} = [\mathfrak{g}_1, \mathfrak{g}_1] \oplus \mathfrak{g}_1$ generated by \mathfrak{g}_1 is called the **Killing superalgebra**
- It behaves as expected: it deforms under geometric limits (e.g., Penrose) and it embeds under asymptotic limits.
- It is a very useful invariant of a supersymmetric supergravity background
- It “categorifies” the supersymmetry fraction ν

- 1 Supergravity
- 2 A geometrical interlude
- 3 Homogeneity in Supergravity

Geometric Killing spinors

- (M, g) a (pseudo-)riemannian spin manifold

Geometric Killing spinors

- (M, g) a (pseudo-)riemannian spin manifold
- $S \rightarrow M$ the spinor bundle (really, Clifford modules)

Geometric Killing spinors

- (M, g) a (pseudo-)riemannian spin manifold
- $S \rightarrow M$ the spinor bundle (really, Clifford modules)
- A spinor $\varepsilon \in \Gamma(S)$ is a **(geometric) Killing spinor** if for all $X \in \mathcal{X}(M)$

$$\nabla_X \varepsilon = \lambda X \cdot \varepsilon$$

for some $\lambda \in \mathbb{R} \cup i\mathbb{R}$

Geometric Killing spinors

- (M, g) a (pseudo-)riemannian spin manifold
- $S \rightarrow M$ the spinor bundle (really, Clifford modules)
- A spinor $\varepsilon \in \Gamma(S)$ is a **(geometric) Killing spinor** if for all $X \in \mathcal{X}(M)$

$$\nabla_X \varepsilon = \lambda X \cdot \varepsilon$$

for some $\lambda \in \mathbb{R} \cup i\mathbb{R}$

- $\lambda \in \mathbb{R}$: **real** Killing spinors

Geometric Killing spinors

- (M, g) a (pseudo-)riemannian spin manifold
- $S \rightarrow M$ the spinor bundle (really, Clifford modules)
- A spinor $\varepsilon \in \Gamma(S)$ is a **(geometric) Killing spinor** if for all $X \in \mathcal{X}(M)$

$$\nabla_X \varepsilon = \lambda X \cdot \varepsilon$$

for some $\lambda \in \mathbb{R} \cup i\mathbb{R}$

- $\lambda \in \mathbb{R}$: **real** Killing spinors
- $\lambda \in i\mathbb{R}$: **imaginary** Killing spinors

Geometric Killing spinors

- (M, g) a (pseudo-)riemannian spin manifold
- $S \rightarrow M$ the spinor bundle (really, Clifford modules)
- A spinor $\varepsilon \in \Gamma(S)$ is a **(geometric) Killing spinor** if for all $X \in \mathcal{X}(M)$

$$\nabla_X \varepsilon = \lambda X \cdot \varepsilon$$

for some $\lambda \in \mathbb{R} \cup i\mathbb{R}$

- $\lambda \in \mathbb{R}$: **real** Killing spinors
- $\lambda \in i\mathbb{R}$: **imaginary** Killing spinors
- they are special types of **twistor spinors**

BAUM+FRIEDRICH+GRUNEWALD+KATH (1991)

Riemannian manifolds admitting Killing spinors

- (M, g) **riemannian** and spin

Riemannian manifolds admitting Killing spinors

- (M, g) **riemannian** and spin
- \exists (nonzero) Killing spinors \implies Einstein with curvature proportional to λ^2

Riemannian manifolds admitting Killing spinors

- (M, g) **riemannian** and spin
- \exists (nonzero) Killing spinors \implies Einstein with curvature proportional to λ^2
 - $\lambda \in i\mathbb{R}$: negative curvature

Riemannian manifolds admitting Killing spinors

- (M, g) **riemannian** and spin
- \exists (nonzero) Killing spinors \implies Einstein with curvature proportional to λ^2
 - $\lambda \in i\mathbb{R}$: negative curvature
 - $\lambda \in \mathbb{R}$: positive curvature and compact

Riemannian manifolds admitting Killing spinors

- (M, g) **riemannian** and spin
- \exists (nonzero) Killing spinors \implies Einstein with curvature proportional to λ^2
 - $\lambda \in i\mathbb{R}$: negative curvature
 - $\lambda \in \mathbb{R}$: positive curvature and compact
- The smallest eigenvalue of the Dirac operator on a compact spin manifold is attained by Killing spinors

FRIEDRICH (1980)

Riemannian manifolds admitting Killing spinors

- (M, g) **riemannian** and spin
- \exists (nonzero) Killing spinors \implies Einstein with curvature proportional to λ^2
 - $\lambda \in i\mathbb{R}$: negative curvature
 - $\lambda \in \mathbb{R}$: positive curvature and compact
- The smallest eigenvalue of the Dirac operator on a compact spin manifold is attained by Killing spinors
FRIEDRICH (1980)
- $\lambda \in i\mathbb{R}$ case completely understood for complete manifolds
BAUM (1989)

Riemannian manifolds admitting Killing spinors

- (M, g) **riemannian** and spin
- \exists (nonzero) Killing spinors \implies Einstein with curvature proportional to λ^2
 - $\lambda \in i\mathbb{R}$: negative curvature
 - $\lambda \in \mathbb{R}$: positive curvature and compact

- The smallest eigenvalue of the Dirac operator on a compact spin manifold is attained by Killing spinors

FRIEDRICH (1980)

- $\lambda \in i\mathbb{R}$ case completely understood for complete manifolds

BAUM (1989)

- $\lambda \in \mathbb{R}$ case reduces to a holonomy problem: *which metric cones admit parallel spinors?*

BÄR (1993), ALEKSEEVESKY+CORTÉS+GALAEV+LEISTNER (2007)

Pseudo-riemannian manifolds admitting Killing spinors

- The pseudo-riemannian case is more complicated

Pseudo-riemannian manifolds admitting Killing spinors

- The pseudo-riemannian case is more complicated
- There are many partial results about the general pseudo-riemannian case

KATH (1998-2000)

Pseudo-riemannian manifolds admitting Killing spinors

- The pseudo-riemannian case is more complicated
- There are many partial results about the general pseudo-riemannian case

KATH (1998-2000)

- The lorentzian case was given particular attention in work of the “Baumschule”

BAUM, BOHLE, KATH, LEITNER, LEISTNER

Pseudo-riemannian manifolds admitting Killing spinors

- The pseudo-riemannian case is more complicated
- There are many partial results about the general pseudo-riemannian case

KATH (1998-2000)

- The lorentzian case was given particular attention in work of the “Baumschule”

BAUM, BOHLE, KATH, LEITNER, LEISTNER

- The cone construction may require the holonomy classification in signature $(2, d - 2)$

Pseudo-riemannian manifolds admitting Killing spinors

- The pseudo-riemannian case is more complicated
- There are many partial results about the general pseudo-riemannian case

KATH (1998-2000)

- The lorentzian case was given particular attention in work of the “Baumschule”

BAUM, BOHLE, KATH, LEITNER, LEISTNER

- The cone construction may require the holonomy classification in signature $(2, d - 2)$
- This is not yet solved, but partial progress has been made

GALAEV

Pseudo-riemannian manifolds admitting Killing spinors

- The pseudo-riemannian case is more complicated
- There are many partial results about the general pseudo-riemannian case

KATH (1998-2000)

- The lorentzian case was given particular attention in work of the “Baumschule”

BAUM, BOHLE, KATH, LEITNER, LEISTNER

- The cone construction may require the holonomy classification in signature $(2, d-2)$
- This is not yet solved, but partial progress has been made

GALAEV

- This is the obstacle to a classification of supersymmetric Freund–Rubin backgrounds

JMF+LEITNER+SIMÓN (20??)

A geometric construction of exceptional Lie algebras

- The analogous construction to the Killing superalgebra applied to the real Killing spinors on the spheres S^7 , S^8 , S^{15} yields 2-graded Lie algebras:

A geometric construction of exceptional Lie algebras

- The analogous construction to the Killing superalgebra applied to the real Killing spinors on the spheres S^7 , S^8 , S^{15} yields 2-graded Lie algebras:
 - $\mathfrak{so}(9)$ for S^7

A geometric construction of exceptional Lie algebras

- The analogous construction to the Killing superalgebra applied to the real Killing spinors on the spheres S^7 , S^8 , S^{15} yields 2-graded Lie algebras:
 - $\mathfrak{so}(9)$ for S^7
 - \mathfrak{f}_4 for S^8

A geometric construction of exceptional Lie algebras

- The analogous construction to the Killing superalgebra applied to the real Killing spinors on the spheres S^7 , S^8 , S^{15} yields 2-graded Lie algebras:
 - $\mathfrak{so}(9)$ for S^7
 - \mathfrak{f}_4 for S^8
 - \mathfrak{e}_8 for S^{15}

A geometric construction of exceptional Lie algebras

- The analogous construction to the Killing superalgebra applied to the real Killing spinors on the spheres S^7 , S^8 , S^{15} yields 2-graded Lie algebras:
 - $\mathfrak{so}(9)$ for S^7
 - \mathfrak{f}_4 for S^8
 - \mathfrak{e}_8 for S^{15}
- One gets either the compact or split real forms of the algebras

A geometric construction of exceptional Lie algebras

- The analogous construction to the Killing superalgebra applied to the real Killing spinors on the spheres S^7 , S^8 , S^{15} yields 2-graded Lie algebras:
 - $\mathfrak{so}(9)$ for S^7
 - \mathfrak{f}_4 for S^8
 - \mathfrak{e}_8 for S^{15}
- One gets either the compact or split real forms of the algebras
- This is the geometrization of Frank Adams's algebraic construction

JMF (2007)

A geometric construction of exceptional Lie algebras

- The analogous construction to the Killing superalgebra applied to the real Killing spinors on the spheres S^7 , S^8 , S^{15} yields 2-graded Lie algebras:
 - $\mathfrak{so}(9)$ for S^7
 - \mathfrak{f}_4 for S^8
 - \mathfrak{e}_8 for S^{15}
- One gets either the compact or split real forms of the algebras
- This is the geometrization of Frank Adams's algebraic construction
- A similar construction exists for pseudo-riemannian manifolds admitting twistor spinors

JMF (2007)

DE MEDEIROS+HOLLANDS (2013)

An early homogeneity conjecture

- Early examples of manifolds admitting Killing spinors: spheres, hyperbolic spaces,... are homogeneous

An early homogeneity conjecture

- Early examples of manifolds admitting Killing spinors: spheres, hyperbolic spaces,... are homogeneous
- if (M, g) admits Killing spinors, is it homogeneous?

An early homogeneity conjecture

- Early examples of manifolds admitting Killing spinors: spheres, hyperbolic spaces,... are homogeneous
- if (M, g) admits Killing spinors, is it homogeneous?
- First known counterexamples: S^5/Γ for $\Gamma < \text{Spin}(6)$ finite
SULANKE (1980)

An early homogeneity conjecture

- Early examples of manifolds admitting Killing spinors: spheres, hyperbolic spaces,... are homogeneous
- if (M, g) admits Killing spinors, is it homogeneous?
- First known counterexamples: S^5/Γ for $\Gamma < \text{Spin}(6)$ finite
SULANKE (1980)
- Many counterexamples are now known (not space forms)

An early homogeneity conjecture

- Early examples of manifolds admitting Killing spinors: spheres, hyperbolic spaces,... are homogeneous
- if (M, g) admits Killing spinors, is it homogeneous?
- First known counterexamples: S^5/Γ for $\Gamma < \text{Spin}(6)$ finite
SULANKE (1980)
- Many counterexamples are now known (not space forms)
- Still, all known simply-connected 6-dimensional riemannian manifolds admitting real Killing spinors (**nearly-Kähler 6-manifolds**) are homogeneous; although there are non-homogeneous quotients
JMF+RITTER (2007)

An early homogeneity conjecture

- Early examples of manifolds admitting Killing spinors: spheres, hyperbolic spaces,... are homogeneous
- if (M, g) admits Killing spinors, is it homogeneous?
- First known counterexamples: S^5/Γ for $\Gamma < \text{Spin}(6)$ finite
SULANKE (1980)
- Many counterexamples are now known (not space forms)
- Still, all known simply-connected 6-dimensional riemannian manifolds admitting real Killing spinors (**nearly-Kähler 6-manifolds**) are homogeneous; although there are non-homogeneous quotients
JMF+RITTER (2007)
- There is hope of constructing non-homogeneous examples via deformations

MOROIANU+SEMMELMANN+NAGY (2006)

- 1 Supergravity
- 2 A geometrical interlude
- 3 Homogeneity in Supergravity

Homogeneous supergravity backgrounds

- A diffeomorphism $\varphi : M \rightarrow M$ is an **automorphism** of a supergravity background (M, g, F) if $\varphi^*g = g$ and $\varphi^*F = F$

Homogeneous supergravity backgrounds

- A diffeomorphism $\varphi : M \rightarrow M$ is an **automorphism** of a supergravity background (M, g, F) if $\varphi^*g = g$ and $\varphi^*F = F$
- Automorphisms form a Lie group $G = \text{Aut}(M, g, F)$

Homogeneous supergravity backgrounds

- A diffeomorphism $\varphi : M \rightarrow M$ is an **automorphism** of a supergravity background (M, g, F) if $\varphi^*g = g$ and $\varphi^*F = F$
- Automorphisms form a Lie group $G = \text{Aut}(M, g, F)$
- A background (M, g, F) is said to be **homogeneous** if G acts transitively on M

Homogeneous supergravity backgrounds

- A diffeomorphism $\varphi : M \rightarrow M$ is an **automorphism** of a supergravity background (M, g, F) if $\varphi^*g = g$ and $\varphi^*F = F$
- Automorphisms form a Lie group $G = \text{Aut}(M, g, F)$
- A background (M, g, F) is said to be **homogeneous** if G acts transitively on M
- Let \mathfrak{g} denote the Lie algebra of G : it consists of vector fields $X \in \mathcal{X}(M)$ such that $\mathcal{L}_X g = 0$ and $\mathcal{L}_X F = 0$

Homogeneous supergravity backgrounds

- A diffeomorphism $\varphi : M \rightarrow M$ is an **automorphism** of a supergravity background (M, g, F) if $\varphi^*g = g$ and $\varphi^*F = F$
- Automorphisms form a Lie group $G = \text{Aut}(M, g, F)$
- A background (M, g, F) is said to be **homogeneous** if G acts transitively on M
- Let \mathfrak{g} denote the Lie algebra of G : it consists of vector fields $X \in \mathcal{X}(M)$ such that $\mathcal{L}_X g = 0$ and $\mathcal{L}_X F = 0$
- (M, g, F) homogeneous \implies the evaluation maps $ev_p : \mathfrak{g} \rightarrow T_p M$ are surjective for all p

Homogeneous supergravity backgrounds

- A diffeomorphism $\varphi : M \rightarrow M$ is an **automorphism** of a supergravity background (M, g, F) if $\varphi^*g = g$ and $\varphi^*F = F$
- Automorphisms form a Lie group $G = \text{Aut}(M, g, F)$
- A background (M, g, F) is said to be **homogeneous** if G acts transitively on M
- Let \mathfrak{g} denote the Lie algebra of G : it consists of vector fields $X \in \mathcal{X}(M)$ such that $\mathcal{L}_X g = 0$ and $\mathcal{L}_X F = 0$
- (M, g, F) homogeneous \implies the evaluation maps $ev_p : \mathfrak{g} \rightarrow T_p M$ are surjective for all p
- The converse is not true in general: if ev_p are surjective, then (M, g, F) is **locally homogeneous**

Homogeneous supergravity backgrounds

- A diffeomorphism $\varphi : M \rightarrow M$ is an **automorphism** of a supergravity background (M, g, F) if $\varphi^*g = g$ and $\varphi^*F = F$
- Automorphisms form a Lie group $G = \text{Aut}(M, g, F)$
- A background (M, g, F) is said to be **homogeneous** if G acts transitively on M
- Let \mathfrak{g} denote the Lie algebra of G : it consists of vector fields $X \in \mathcal{X}(M)$ such that $\mathcal{L}_X g = 0$ and $\mathcal{L}_X F = 0$
- (M, g, F) homogeneous \implies the evaluation maps $ev_p : \mathfrak{g} \rightarrow T_p M$ are surjective for all p
- The converse is not true in general: if ev_p are surjective, then (M, g, F) is **locally homogeneous**
- This is the “right” working notion in supergravity

The homogeneity theorem

Empirical Fact

Every known ν -BPS background with $\nu > \frac{1}{2}$ is homogeneous.

The homogeneity theorem

Homogeneity conjecture

Every ~~known~~ ν -BPS background with $\nu > \frac{1}{2}$ is homogeneous.

MEESSEN (2004)

The homogeneity theorem

Homogeneity conjecture

Every ~~known~~ ν -BPS background with $\nu > \frac{1}{2}$ is homogeneous.

MEESSEN (2004)

Theorem

Every ν -BPS background of eleven-dimensional supergravity with $\nu > \frac{1}{2}$ is locally homogeneous.

JMF+MEESSEN+PHILIP (2004), JMF+HUSTLER (2012)

The homogeneity theorem

Homogeneity conjecture

Every ~~known~~ ν -BPS background with $\nu > \frac{1}{2}$ is homogeneous.

MEESSEN (2004)

Theorem

Every ν -BPS background of eleven-dimensional supergravity with $\nu > \frac{1}{2}$ is locally homogeneous.

JMF+MEESSEN+PHILIP (2004), JMF+HUSTLER (2012)

The existence of non-homogeneous $\nu = \frac{1}{2}$ backgrounds shows that the theorem is sharp. But, **why** $\frac{1}{2}$?

Generalisations

Theorem

Every ν -BPS background of type IIB supergravity with $\nu > \frac{1}{2}$ is locally homogeneous.

Every ν -BPS background of type I and heterotic supergravities with $\nu > \frac{1}{2}$ is locally homogeneous.

JMF+HACKETT-JONES+MOUTSOPOULOS (2007)

JMF+HUSTLER (2012)

Every ν -BPS background of six-dimensional $(1,0)$ and $(2,0)$ supergravities with $\nu > \frac{1}{2}$ is locally homogeneous.

JMF + HUSTLER (2013)

Generalisations

Theorem

Every ν -BPS background of type IIB supergravity with $\nu > \frac{1}{2}$ is locally homogeneous.

Every ν -BPS background of type I and heterotic supergravities with $\nu > \frac{1}{2}$ is locally homogeneous.

JMF+HACKETT-JONES+MOUTSOPOULOS (2007)

JMF+HUSTLER (2012)

Every ν -BPS background of six-dimensional $(1,0)$ and $(2,0)$ supergravities with $\nu > \frac{1}{2}$ is locally homogeneous.

JMF + HUSTLER (2013)

The theorems actually prove the strong version of the conjecture: that the symmetries which are generated from the supersymmetries already act (locally) transitively.

Poincaré supergravities again

	32			24	20	16	12	8	4
11	M								
10	IIA	IIB				I			
9	N = 2					N = 1			
8	N = 2					N = 1			
7	N = 4					N = 2			
6	(2,2)	(3,1)	(4,0)	(2,1)	(3,0)	(1,1)	(2,0)	(1,0)	
5	N = 8			N = 6		N = 4		N = 2	
4	N = 8			N = 6	N = 5	N = 4	N = 3	N = 2	N = 1

Idea of proof

The proof consists of two steps:

Idea of proof

The proof consists of two steps:

- 1 One shows the existence of the Killing superalgebra

$$\mathfrak{k} = \mathfrak{k}_0 \oplus \mathfrak{k}_1$$

Idea of proof

The proof consists of two steps:

- 1 One shows the existence of the Killing superalgebra
 $\mathfrak{k} = \mathfrak{k}_0 \oplus \mathfrak{k}_1$
- 2 One shows that for all $p \in M$, $ev_p : \mathfrak{k}_0 \rightarrow T_p M$ is surjective
whenever $\dim \mathfrak{k}_1 > \frac{1}{2} \text{rank } S$

Idea of proof

The proof consists of two steps:

- 1 One shows the existence of the Killing superalgebra
 $\mathfrak{k} = \mathfrak{k}_0 \oplus \mathfrak{k}_1$
- 2 One shows that for all $p \in M$, $ev_p : \mathfrak{k}_0 \rightarrow T_p M$ is surjective
whenever $\dim \mathfrak{k}_1 > \frac{1}{2} \text{rank } S$

This actually only shows local homogeneity.

What good is it?

The homogeneity theorem implies that classifying homogeneous supergravity backgrounds also classifies ν -BPS backgrounds for $\nu > \frac{1}{2}$.

What good is it?

The homogeneity theorem implies that classifying homogeneous supergravity backgrounds also classifies ν -BPS backgrounds for $\nu > \frac{1}{2}$.

This is **good** because

- the supergravity field equations for homogeneous backgrounds are algebraic and hence simpler to solve than PDEs

What good is it?

The homogeneity theorem implies that classifying homogeneous supergravity backgrounds also classifies ν -BPS backgrounds for $\nu > \frac{1}{2}$.

This is **good** because

- the supergravity field equations for homogeneous backgrounds are algebraic and hence simpler to solve than PDEs
- we have learnt **a lot** (about string theory) from supersymmetric supergravity backgrounds, so their classification could teach us even more

Searching for homogeneous supergravity backgrounds

A homogeneous eleven-dimensional supergravity background is described algebraically by the data $(\mathfrak{g}, \mathfrak{h}, \gamma, \varphi)$, where

Searching for homogeneous supergravity backgrounds

A homogeneous eleven-dimensional supergravity background is described algebraically by the data $(\mathfrak{g}, \mathfrak{h}, \gamma, \varphi)$, where

- $\mathfrak{g} = \mathfrak{h} \oplus \mathfrak{m}$ with $\dim \mathfrak{m} = 11$

Searching for homogeneous supergravity backgrounds

A homogeneous eleven-dimensional supergravity background is described algebraically by the data $(\mathfrak{g}, \mathfrak{h}, \gamma, \varphi)$, where

- $\mathfrak{g} = \mathfrak{h} \oplus \mathfrak{m}$ with $\dim \mathfrak{m} = 11$
- γ is an \mathfrak{h} -invariant lorentzian inner product on \mathfrak{m}

Searching for homogeneous supergravity backgrounds

A homogeneous eleven-dimensional supergravity background is described algebraically by the data $(\mathfrak{g}, \mathfrak{h}, \gamma, \varphi)$, where

- $\mathfrak{g} = \mathfrak{h} \oplus \mathfrak{m}$ with $\dim \mathfrak{m} = 11$
- γ is an \mathfrak{h} -invariant lorentzian inner product on \mathfrak{m}
- φ is an \mathfrak{h} -invariant 4-form $\varphi \in \Lambda^4 \mathfrak{m}$

Searching for homogeneous supergravity backgrounds

A homogeneous eleven-dimensional supergravity background is described algebraically by the data $(\mathfrak{g}, \mathfrak{h}, \gamma, \varphi)$, where

- $\mathfrak{g} = \mathfrak{h} \oplus \mathfrak{m}$ with $\dim \mathfrak{m} = 11$
- γ is an \mathfrak{h} -invariant lorentzian inner product on \mathfrak{m}
- φ is an \mathfrak{h} -invariant 4-form $\varphi \in \Lambda^4 \mathfrak{m}$

subject to some algebraic equations which are given purely in terms of the structure constants of \mathfrak{g} (and \mathfrak{h}).

Methodology

Classifying homogeneous supergravity backgrounds of a certain type involves now the following steps:

Methodology

Classifying homogeneous supergravity backgrounds of a certain type involves now the following steps:

- 1 Classify the desired homogeneous geometries

Methodology

Classifying homogeneous supergravity backgrounds of a certain type involves now the following steps:

- 1 Classify the desired homogeneous geometries
- 2 For each such geometry parametrise the space of invariant lorentzian metrics and invariant closed 4-forms

Methodology

Classifying homogeneous supergravity backgrounds of a certain type involves now the following steps:

- 1 Classify the desired homogeneous geometries
- 2 For each such geometry parametrise the space of invariant lorentzian metrics and invariant closed 4-forms
- 3 Plug them into the supergravity field equations to get (nonlinear) algebraic equations for the parameters

Methodology

Classifying homogeneous supergravity backgrounds of a certain type involves now the following steps:

- 1 Classify the desired homogeneous geometries
- 2 For each such geometry parametrise the space of invariant lorentzian metrics and invariant closed 4-forms
- 3 Plug them into the supergravity field equations to get (nonlinear) algebraic equations for the parameters
- 4 Solve the equations!

Homogeneous lorentzian manifolds I

- Their classification can seem daunting!

Homogeneous lorentzian manifolds I

- Their classification can seem daunting!
- We wish to classify d -dimensional lorentzian manifolds (M, g) homogeneous under a Lie group G .

Homogeneous lorentzian manifolds I

- Their classification can seem daunting!
- We wish to classify d -dimensional lorentzian manifolds (M, g) homogeneous under a Lie group G .
- Then $M \cong G/H$ with H a closed subgroup.

Homogeneous lorentzian manifolds I

- Their classification can seem daunting!
- We wish to classify d -dimensional lorentzian manifolds (M, g) homogeneous under a Lie group G .
- Then $M \cong G/H$ with H a closed subgroup.
- One starts by classifying Lie subalgebras $\mathfrak{h} \subset \mathfrak{g}$ with

Homogeneous lorentzian manifolds I

- Their classification can seem daunting!
- We wish to classify d -dimensional lorentzian manifolds (M, g) homogeneous under a Lie group G .
- Then $M \cong G/H$ with H a closed subgroup.
- One starts by classifying Lie subalgebras $\mathfrak{h} \subset \mathfrak{g}$ with
 - codimension d

Homogeneous lorentzian manifolds I

- Their classification can seem daunting!
- We wish to classify d -dimensional lorentzian manifolds (M, g) homogeneous under a Lie group G .
- Then $M \cong G/H$ with H a closed subgroup.
- One starts by classifying Lie subalgebras $\mathfrak{h} \subset \mathfrak{g}$ with
 - codimension d
 - Lie subalgebras of closed subgroups

Homogeneous lorentzian manifolds I

- Their classification can seem daunting!
- We wish to classify d -dimensional lorentzian manifolds (M, g) homogeneous under a Lie group G .
- Then $M \cong G/H$ with H a closed subgroup.
- One starts by classifying Lie subalgebras $\mathfrak{h} \subset \mathfrak{g}$ with
 - codimension d
 - Lie subalgebras of closed subgroups
 - leaving invariant a lorentzian inner product on $\mathfrak{g}/\mathfrak{h}$

Homogeneous lorentzian manifolds I

- Their classification can seem daunting!
- We wish to classify d -dimensional lorentzian manifolds (M, g) homogeneous under a Lie group G .
- Then $M \cong G/H$ with H a closed subgroup.
- One starts by classifying Lie subalgebras $\mathfrak{h} \subset \mathfrak{g}$ with
 - codimension d
 - Lie subalgebras of closed subgroups
 - leaving invariant a lorentzian inner product on $\mathfrak{g}/\mathfrak{h}$
- Hopeless except in low dimension or if G is semisimple

Homogeneous lorentzian manifolds I

- Their classification can seem daunting!
- We wish to classify d -dimensional lorentzian manifolds (M, g) homogeneous under a Lie group G .
- Then $M \cong G/H$ with H a closed subgroup.
- One starts by classifying Lie subalgebras $\mathfrak{h} \subset \mathfrak{g}$ with
 - codimension d
 - Lie subalgebras of closed subgroups
 - leaving invariant a lorentzian inner product on $\mathfrak{g}/\mathfrak{h}$
- Hopeless except in low dimension or if G is semisimple

Definition

The action of G on M is **proper** if the map $G \times M \rightarrow M \times M$, $(\gamma, m) \mapsto (\gamma \cdot m, m)$ is proper (i.e., inverse image of compact is compact). In particular, proper actions have compact stabilisers.

Homogeneous lorentzian manifolds II

What if the action is not proper?

Homogeneous lorentzian manifolds II

What if the action is not proper?

Theorem (Kowalsky, 1996)

If a simple Lie group acts transitively and non-properly on a lorentzian manifold (M, g) , then (M, g) is locally isometric to (anti) de Sitter spacetime.

Homogeneous lorentzian manifolds II

What if the action is not proper?

Theorem (Kowalsky, 1996)

If a simple Lie group acts transitively and non-properly on a lorentzian manifold (M, g) , then (M, g) is locally isometric to (anti) de Sitter spacetime.

Theorem (Deffaf–Melnick–Zeghib, 2008)

If a semisimple Lie group acts transitively and non-properly on a lorentzian manifold (M, g) , then (M, g) is locally isometric to the product of (anti) de Sitter spacetime and a riemannian homogeneous space.

Homogeneous lorentzian manifolds II

What if the action is not proper?

Theorem (Kowalsky, 1996)

If a simple Lie group acts transitively and non-properly on a lorentzian manifold (M, g) , then (M, g) is locally isometric to (anti) de Sitter spacetime.

Theorem (Deffaf–Melnick–Zeghib, 2008)

If a semisimple Lie group acts transitively and non-properly on a lorentzian manifold (M, g) , then (M, g) is locally isometric to the product of (anti) de Sitter spacetime and a riemannian homogeneous space.

This means that we need only classify Lie subalgebras corresponding to *compact* Lie subgroups!

Some recent classification results

- Symmetric eleven-dimensional supergravity backgrounds
JMF (2011)

Some recent classification results

- Symmetric eleven-dimensional supergravity backgrounds
JMF (2011)
- Symmetric type IIB supergravity backgrounds
JMF+HUSTLER (2012)

Some recent classification results

- Symmetric eleven-dimensional supergravity backgrounds
JMF (2011)
- Symmetric type IIB supergravity backgrounds
JMF+HUSTLER (2012)
- Homogeneous M2-duals: $\mathfrak{g} = \mathfrak{so}(3, 2) \oplus \mathfrak{so}(N)$ for $N > 4$
JMF+UNGUREANU (IN PREPARATION)

Outlook

- With patience and optimism, some classes of homogeneous backgrounds can be classified

Outlook

- With patience and optimism, some classes of homogeneous backgrounds can be classified
- In particular, we can “dial up” a semisimple G and hope to solve the homogeneous supergravity equations with symmetry G

Outlook

- With patience and optimism, some classes of homogeneous backgrounds can be classified
- In particular, we can “dial up” a semisimple G and hope to solve the homogeneous supergravity equations with symmetry G
- Checking supersymmetry is an additional problem, but there is an efficient algorithm which has already discarded many of the symmetric eleven-dimensional backgrounds.
LISCHEWSKI (2014), HUSTLER+LISCHEWSKI (IN PROGRESS)