José Figueroa-O'Farrill Edinburgh Mathematical Physics Group School of Mathematics



IAS, 12 May 2003

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Killing's identity \implies a Killing vector ξ is uniquely determined by:

• its value $\xi(p)$ at a point p; and

• the value $\nabla_{\mu}\xi_{\nu}(p) = -\nabla_{\nu}\xi_{\mu}(p)$ of its derivative at p.

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The flat and spherical cases are solved (culminating in the work of Wolf in the 1970s), but the hyperbolic and lorentzian cases remain largely open despite many partial results.

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Supersymmetric space forms

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In this talk I will report on progress towards answering this question.

Based on work in collaboration with

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George Papadopoulos (King's College, London)
 * hep-th/0211089 (JHEP 03 (2003) 048)
 * math.AG/0211170

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• Ali Chamseddine and Wafic Sabra (CAMS, Beirut)

 \star in preparation

Supergravities

	32			24		20	16		12	8	4	
11	М	M										
10	IIA	IIB						1				
9	N=2						N = 1					
8	N=2							N = 1				
7	N = 4							N=2				
6	(2,	2)	(3,1)	(4, 0)	(2,1)	(3,0)		(1,1)	(2,0)		(1, 0)	
5	N=8			N = 6			N=4			N = 2		
4		N = 8			N = 6		N = 5	N = 4		N=3	N=2	N = 1

[Van Proeyen, hep-th/0301005]

Let (M, g, Φ, S) be a supergravity background:

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(M, g, Φ, S) is supersymmetric if it admits Killing spinors

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defined by the supersymmetric variation of the gravitino:

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• spinors are Majorana; that is, associated to one of the two irreducible real 32-dimensional representations of $C\ell(1, 10)$. Therefore the gravitino also has 128 physical degrees of freedom.

$$D_{\mu} = \nabla_{\mu} - \frac{1}{288} F_{\nu\rho\sigma\tau} \left(\Gamma^{\nu\rho\sigma\tau}{}_{\mu} + 8\Gamma^{\nu\rho\sigma}\delta^{\tau}_{\mu} \right)$$

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where *I* is an index labeling the following elements

$$\Gamma_a \quad \Gamma_{ab} \quad \Gamma_{abc} \quad \Gamma_{abcd} \quad \Gamma_{abcde}$$

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with T quadratic in F. This means that $R_{\mu\nu\rho\sigma}$ is parallel; equivalently, that g is locally symmetric.

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- if the plane is null, we can choose a pseudo-orthonormal frame in which the only nonzero component of F is F_{-123}

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• F lorentzian: a one parameter R < 0 family of vacua

 $AdS_4(8R) \times S^7(-7R)$ $F = \sqrt{-6R} dvol(AdS_4)$

• F null: a one parameter $\mu \in \mathbb{R}$ family of symmetric plane waves:

$$g = 2dx^{+}dx^{-} - \frac{1}{36}\mu^{2} \left(4\sum_{i=1}^{3} (x^{i})^{2} + \sum_{i=4}^{9} (x^{i})^{2}\right) (dx^{-})^{2} + \sum_{i=1}^{9} (dx^{i})^{2}$$

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All vacua embed isometrically in $\mathbb{E}^{2,11}$ as the intersections of two quadrics.

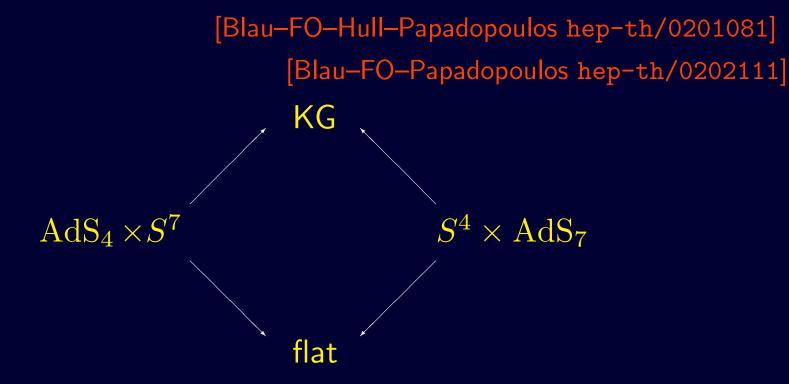
Solutions are related by Penrose limits

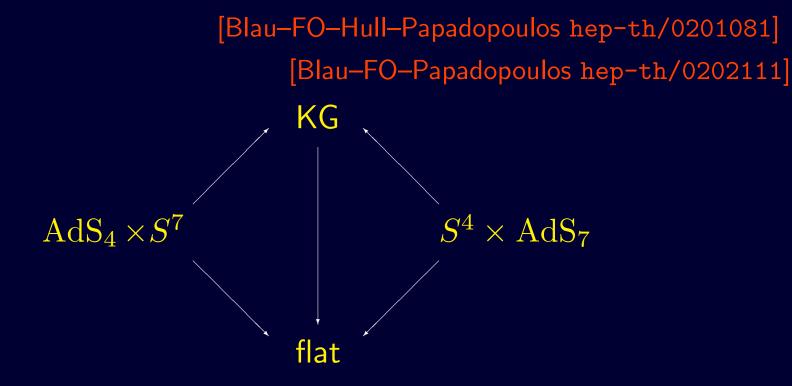
[Blau-FO-Hull-Papadopoulos hep-th/0201081] [Blau-FO-Papadopoulos hep-th/0202111]

 $AdS_4 \times S^7$

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[Back]

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Maximal supersymmetry $\implies D$ is flat.

 Theorem (Cartan–Schouten (1926), Wolf (1971/2),

 Cahen–Parker (1977)).

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The solution to this problem is known.

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But there is a more general construction.

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• since \mathfrak{h} preserves the metric on \mathfrak{g} , there is a linear map

 $\mathfrak{h} \to \Lambda^2 \mathfrak{g}$

$$\omega:\Lambda^2\mathfrak{g} o\mathfrak{h}^*$$

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relative to bases X_a , H_i and H^i for \mathfrak{g} , \mathfrak{h} and \mathfrak{h}^* , respectively.

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• \mathfrak{h} acts on $\mathfrak{g} \times_{\omega} \mathfrak{h}^*$ preserving the Lie bracket, so we can form the *double extension*

 $\mathfrak{d}(\mathfrak{g},\mathfrak{h})=\mathfrak{h}\ltimes(\mathfrak{g} imes_\omega\mathfrak{h}^*)$

$$egin{array}{cccc} X_b & H_j & H^j \ X_a & \left(egin{array}{cccc} g_{ab} & 0 & 0 \ 0 & B_{ij} & \delta^j_i \ 0 & \delta^j_j & 0 \end{array}
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This construction is due to Medina and Revoy who proved an important structure theorem.

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An indecomposable metric Lie algebra is either simple, onedimensional, or a double extension $\mathfrak{d}(\mathfrak{g},\mathfrak{h})$ where \mathfrak{h} is either simple or one-dimensional.

Every metric Lie algebra is obtained as an orthogonal direct sum of indecomposables.

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where \mathfrak{a} is abelian with euclidean metric and \mathfrak{h} is one-dimensional. (Any semisimple factors in \mathfrak{a} factor out of the double extension. [FO-Stanciu hep-th/9402035])



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$$\mathbb{R} \to \Lambda^2 \mathbb{R}^4 \cong \mathfrak{so}(4)$$



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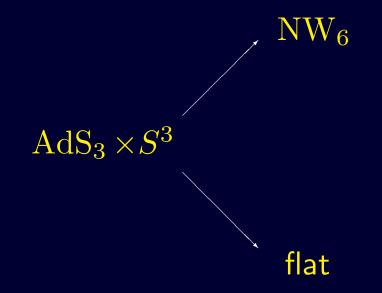
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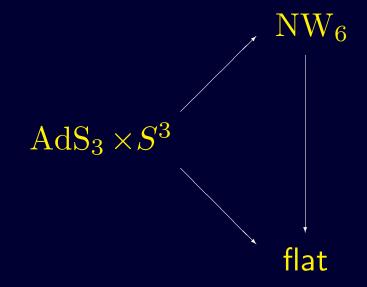
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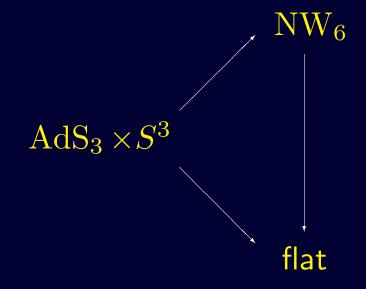
The third case is a six-dimensional version of the Nappi-Witten spacetime, NW_6 , discovered by Meessen. [Meessen hep-th/0111031]

$AdS_3 \times S^3$

 NW_6 $\mathrm{AdS}_3 imes S^3$







which in this case are *group contractions* à la Inönü–Wigner. [Stanciu–FO hep-th/0303212]

[Back]

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(1,0) vacua $\leftrightarrow (2,0)$ vacua up to $\mathrm{Sp}(2)$ R-symmetry.

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Expanding the curvature of D into antisymmetric products of Γ -matrices

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This equation defines a generalisation of a Lie algebra known as a4-Lie algebra (with an invariant metric).[Filippov (1985)]

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Again we can work in the tangent space at a point, where g gives rise to a lorentzian innner product and F defines a self-dual 5-form obeying a quadratic equation.

This equation defines a generalisation of a Lie algebra known as a4-Lie algebra (with an invariant metric).[Filippov (1985)]

(Unfortunate notation: a 2-Lie algebra is a Lie algebra.)

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 $ad_X[Y, Z] = [ad_X Y, Z] + [Y, ad_X Z]$

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If F is totally antisymmetric then $\langle -, - \rangle$ is an *invariant metric*. (*n*-Lie algebras also appear naturally in the context of Nambu dynamics. [Nambu (1973)]

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[FO-Papadopoulos math.AG/0211170]

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Notice that *g* is a bi-invariant metric on a Lie group: a ten-dimensional version of the Nappi–Witten spacetime.

[Stanciu-FO hep-th/0303212]

These vacua again embed isometrically in $\mathbb{E}^{2,10}$ as intersections of quadrics

[Blau-FO-Hull-Papadopoulos hep-th/0201081]

[Berenstein-Maldacena-Nastase hep-th/0202021]

 $AdS_5 \times S^5$

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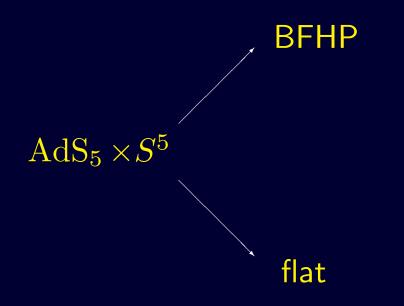
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BFHP

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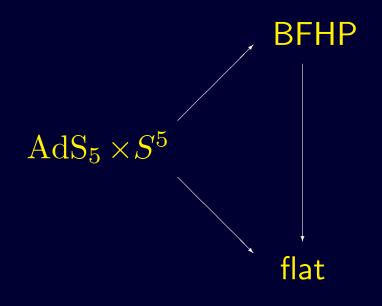
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[Back]

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[Gauntlett-Gutowsky-Hull-Pakis-Reall hep-th/0209114] [Lozano-Tellechea-Meessen-Ortín hep-th/0206200]

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[Gauntlett-Gutowsky-Hull-Pakis-Reall hep-th/0209114] [Lozano-Tellechea-Meessen-Ortín hep-th/0206200]

This now has a natural explanation.

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Let ξ be a (spacelike) left-invariant vector field

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Thank you.