Plane wave limits and the string/gauge theory correspondence

José Figueroa-O'Farrill Edinburgh Mathematical Physics Group



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['t Hooft (1974)]

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$$\sum_{g=0}^{\infty} N^{2-2g} f_g(\lambda)$$

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Which string theory?

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 $S_{\rm bulk}$

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IIB supergravity in bulk + IIB supergravity near brane horizon

[Maldacena (1997)]

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- gravity approximation is valid for $R^4 \gg \ell_s^4$

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 spectrum of, say, single trace operators; e.g., chiral primaries (in short multiplets)

 $\operatorname{Tr}' \phi \otimes \cdots \otimes \phi$

But...

But...further checks are hindered by lack of clear dictionary between *massive* string modes

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However...recent progress in this direction has been made using a different large N limit, the so-called plane wave limit.

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Boost along γ and zoom in on γ while rescaling g.

 $g = 2dUdV + AdV^2 + B_i dY^i dV + C_{ij} dY^i dY^j$

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• rescale

$$g_{\mathsf{plane wave}} = \lim_{\Omega \to 0} \Omega^{-2} g(\Omega)$$

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• "lightlike gauge"

 $\imath_{\partial/\partial U}A^{(p)} = 0$

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Note: field strength is *null*:

$$F_{\text{plane wave}} = dA_{\text{plane wave}} = dx^+ \wedge \Theta(x^+)$$

11

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It is useful in classifying the different plane wave limits of a given background.

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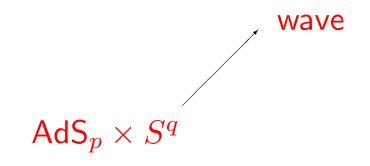
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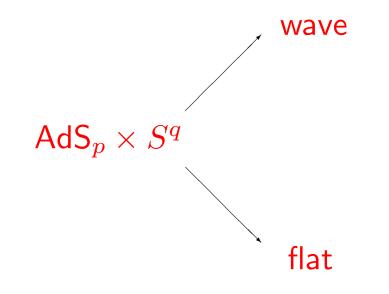
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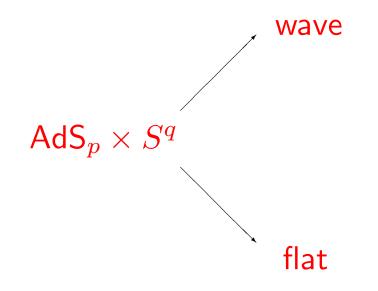
 $\mathsf{PWL}:\mathsf{Vacua}\to\mathsf{Vacua}$

Vacua of the form $\operatorname{AdS}_p \times S^q$

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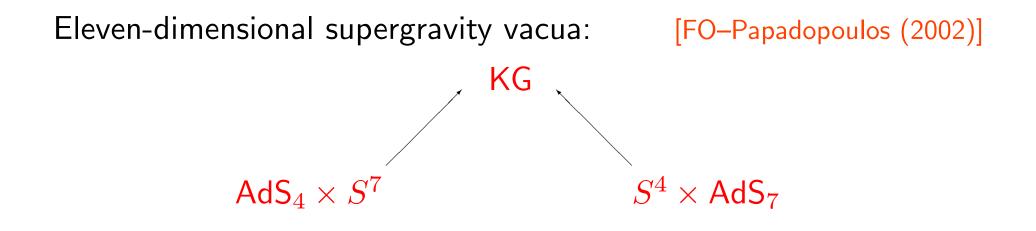
depending on whether or not the component of $\dot{\gamma}$ tangent to S^q vanishes. [Blau et al. (2002)]

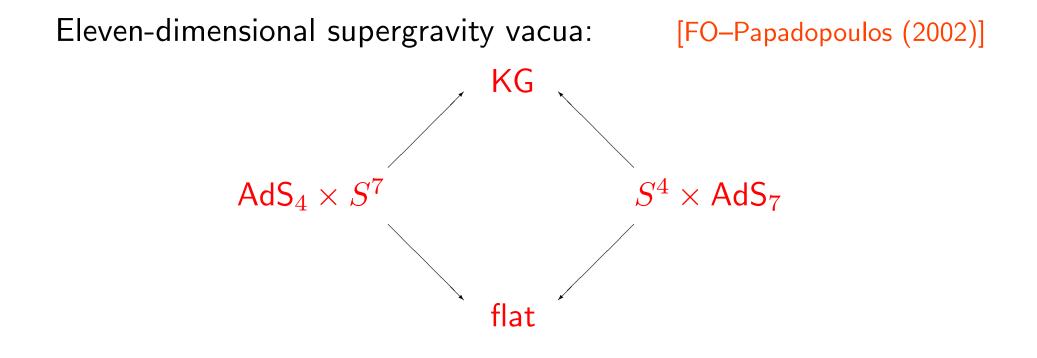
Eleven-dimensional supergravity vacua: [FO–Papadopoulos (2002)]

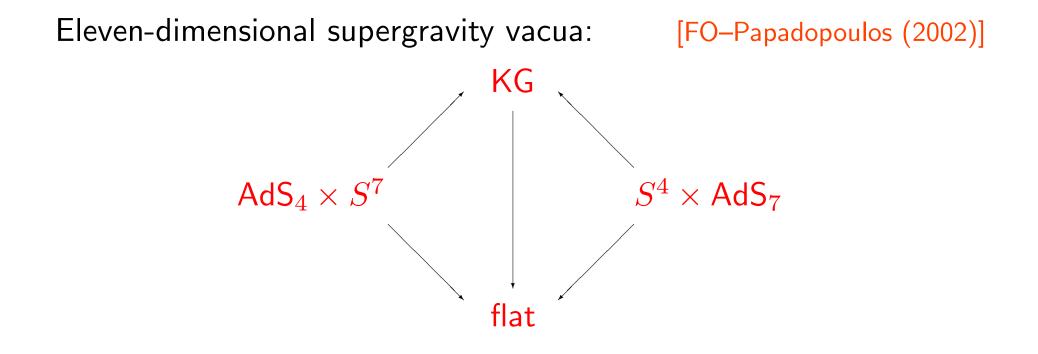
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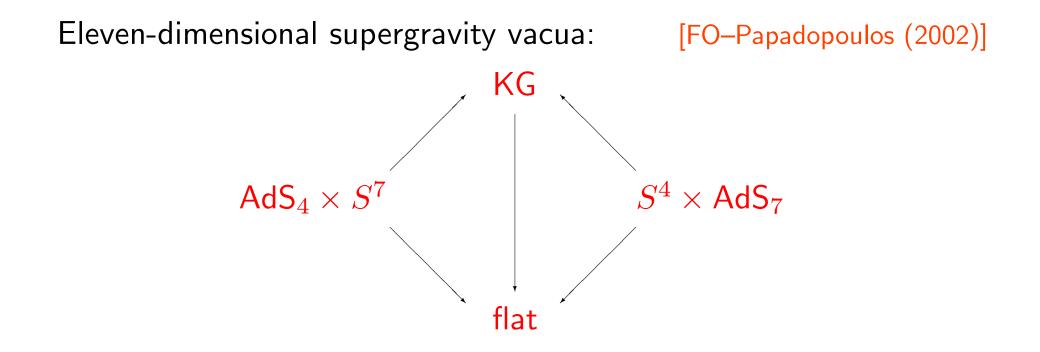
 $\mathsf{AdS}_4 \times S^7$

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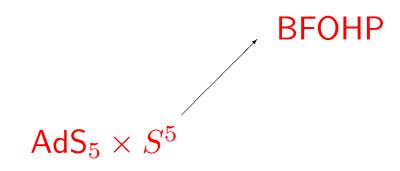
KG: 11-dimensional symmetric plane wave [Kowalski-Glikman (1984)]

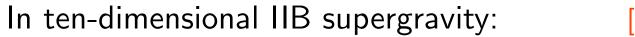
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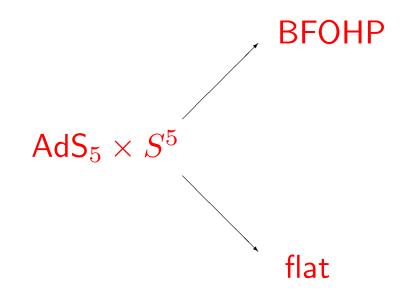


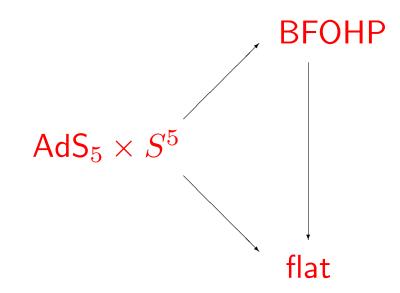
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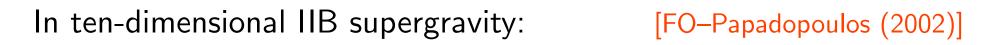


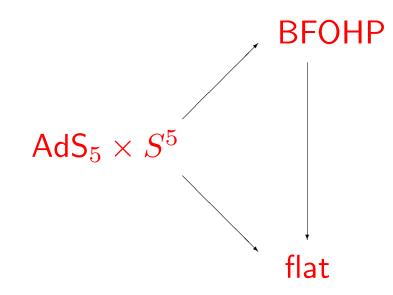














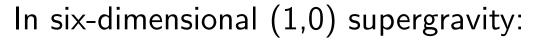
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[Chamseddine et al]

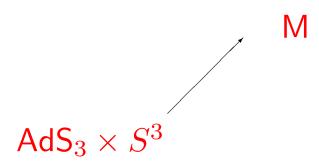
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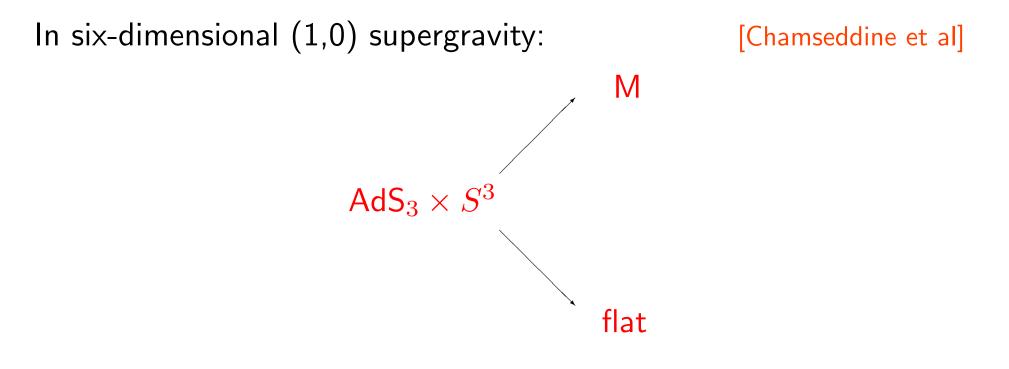
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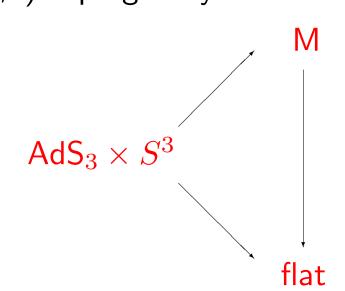
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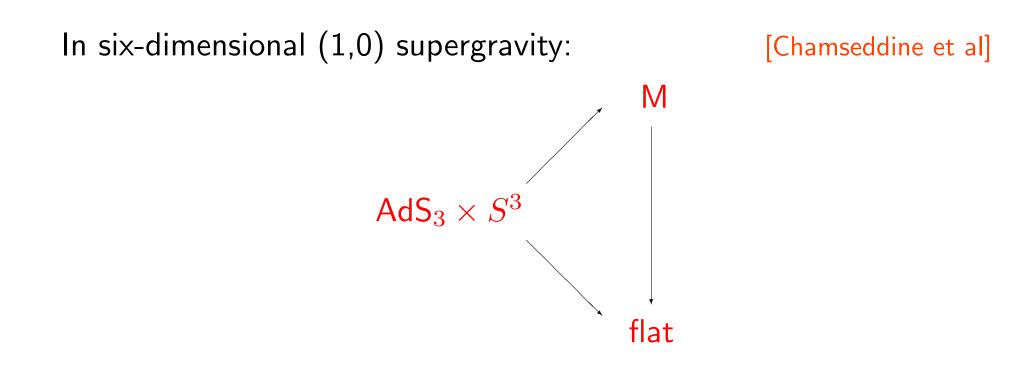






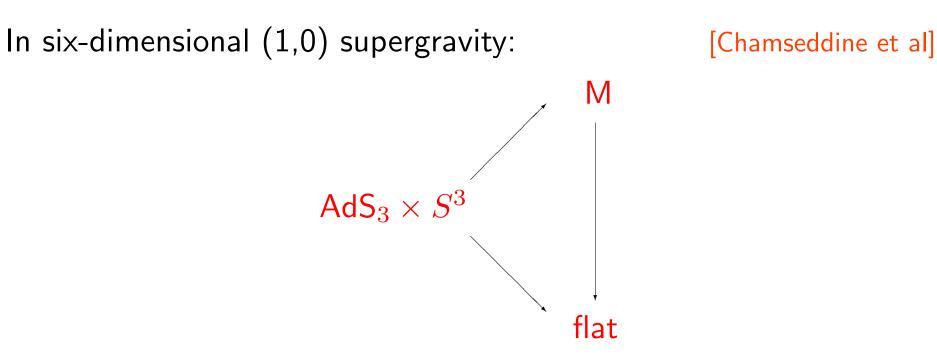
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M: 6-dimensional symmetric plane wave

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Symmetric plane waves

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where A_{ij} is *constant*, and is determined by its eigenvalues, up to order and scale.

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- all other fluxes vanish
- μ inessential, but $\mu \rightarrow 0$ recovers flat solution

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$$\frac{1}{2\pi\alpha'} \int d\tau \int_0^{2\pi\alpha' p^+} d\sigma \left[\frac{1}{2} \dot{x}^2 - \frac{1}{2} {x'}^2 - \frac{1}{2} \mu^2 x^2 + i\bar{\psi}(\partial \!\!\!/ + \mu M)\psi \right]$$

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where $x: W \to \mathbb{R}^8$

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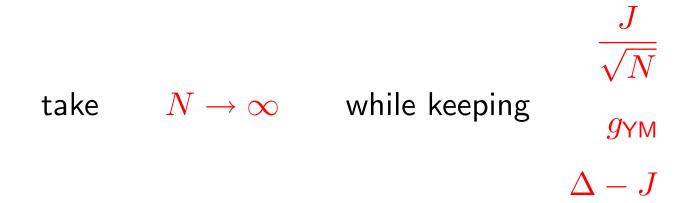
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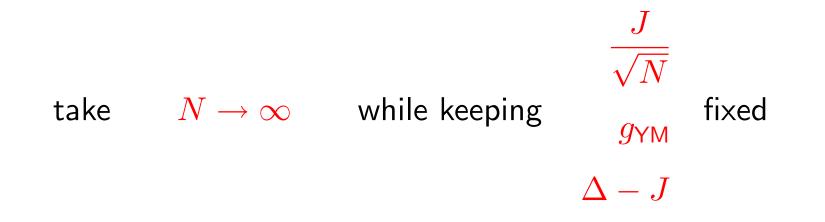


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gym

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Notice: BPS condition $\Delta \ge |J| \implies p^{\pm} \ge 0$

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Notice: As $N \to \infty$ we focus on states with larger and larger J. Thus observables are not held fixed in this limit.

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- they are dual to the *free* string excitations on the plane wave background

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BMN operators are single trace operators consisting in a string of J Z's and a finite number of *impurities*.

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How about interactions?

Interacting string theory

A detailed study of BMN correlators shows

Interacting string theory

A detailed study of BMN correlators shows: [Kristjansen et al., Gross et al., Constable et al. (2002)]

• theory develops a different effective coupling constant

$$\lambda' = \frac{g_{\rm YM}^2 N}{J^2} = \frac{1}{(\mu p^+ \alpha')^2}$$

• a different genus-counting parameter

$$g_2^2 = \left(\frac{J^2}{N}\right)^2 = 16\pi^2 g_s^2 (\mu p^+ \alpha')^4$$

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[..., Chu-Khoze, Gomis et al. (2003)]

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 the correspondence has been extended to theories with less supersymmetry; but as usual QCD remains elusive.

Thank you.