Quotienting string backgrounds

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1

• JHEP 12 (2001) 011, hep-th/0110170

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- $\mathbb{R}^{1,4}/\Gamma$, with $\Gamma \cong \mathbb{R}$, or $\mathbb{R}^{1,10}/\Gamma \implies$ IIA fluxbranes

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• we are interested in the orbit space M/Γ



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 - \star M supersymmetric, but M/Γ_L breaking all supersymmetry

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i.e., projectivised adjoint orbits of ${\mathfrak g}$

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$$D_X = \nabla_X + \frac{1}{6}\iota_X F - \frac{1}{12}X^{\flat} \wedge F$$



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which we need to put in normal form

 $\bullet \ X \ \in \ \mathfrak{so}(p,q)$

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• for each indecomposable block, if λ is an eigenvalue, then so are $-\lambda$, λ^* , and $-\lambda^*$

 \star $\lambda = 0$

$$\star \ \lambda = 0 \qquad \qquad \mu(x) = x^r$$

$$\begin{array}{l} \star \ \lambda = 0 \\ \star \ \lambda = \beta \in \mathbb{R} \end{array} \qquad \qquad \mu(x) = x^n \\ \end{array}$$

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 $\star \ \lambda = \beta + i \varphi$

 $\begin{array}{l} \star \ \lambda = 0 \qquad \qquad \mu(x) = x^n \\ \star \ \lambda = \beta \in \mathbb{R}, \qquad \qquad \mu(x) = (x^2 - \beta^2)^n \end{array}$

 $\star \lambda = \beta + i \varphi$, $\beta \varphi \neq 0$

 $\begin{array}{ll} \star \ \lambda = 0 & \mu(x) = x^n \\ \star \ \lambda = \beta \in \mathbb{R}, & \mu(x) = (x^2 - \beta^2)^n \end{array}$

*
$$\lambda = \beta + i\varphi, \ \beta\varphi \neq 0,$$

$$\mu(x) = \left(\left(x^2 + \beta^2 + \varphi^2 \right)^2 - 4\beta^2 x^2 \right)^n$$

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- example: $\mu(x) = x^3$

Signature Minimal polynomial Type





Signature	Minimal polynomial	Туре
(0,1)	x	trivial
Signature	Minimal polynomial	Туре
-----------	--------------------	---------
(0,1)	x	trivial
(0,2)		

Signature	Minimal polynomial	Туре
(0,1)	x	trivial
(0,2)	$x^2 + \varphi^2$	

Signature	Minimal polynomial	Туре
(0,1)	x	trivial
(0,2)	$x^2 + \varphi^2$	rotation

Signature	Minimal polynomial	Туре
(0,1)	x	trivial
(0,2)	$x^2 + \varphi^2$	rotation
(1, 0)		

Signature	Minimal polynomial	Туре
(0,1)	x	trivial
(0,2)	$x^2 + \varphi^2$	rotation
(1, 0)	x	

Signature	Minimal polynomial	Туре
(0,1)	x	trivial
(0,2)	$x^2 + \varphi^2$	rotation
(1, 0)	x	trivial

Signature	Minimal polynomial	Туре
(0,1)	x	trivial
(0,2)	$x^2 + \varphi^2$	rotation
(1, 0)	x	trivial
(1, 1)		

Signature	Minimal polynomial	Туре
(0,1)	x	trivial
(0,2)	$x^2 + \varphi^2$	rotation
(1,0)	x	trivial
(1, 1)	$x^2 - \beta^2$	

Signature	Minimal polynomial	Туре
(0,1)	x	trivial
(0,2)	$x^2 + \varphi^2$	rotation
(1, 0)	x	trivial
(1,1)	$x^2 - \beta^2$	boost

Signature	Minimal polynomial	Туре
(0,1)	x	trivial
(0,2)	$x^2 + \varphi^2$	rotation
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Signature	Minimal polynomial	Туре
(0,1)	x	trivial
(0,2)	$x^2 + \varphi^2$	rotation
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(1,1)	$x^2 - \beta^2$	boost
(1,2)	x^3	null rotation







• $R_{12}(\varphi_1) + R_{34}(\varphi_2) + R_{56}(\varphi_3) + R_{78}(\varphi_4) + R_{9\natural}(\varphi_5)$



In signature (1, 10):

- $R_{12}(\varphi_1) + R_{34}(\varphi_2) + R_{56}(\varphi_3) + R_{78}(\varphi_4) + R_{9\natural}(\varphi_5)$
- $B_{02}(\beta) + R_{34}(\varphi_1) + R_{56}(\varphi_2) + R_{78}(\varphi_3) + R_{9\natural}(\varphi_4)$



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- $N_{+2} + R_{34}(\varphi_1) + R_{56}(\varphi_2) + R_{78}(\varphi_3) + R_{9\natural}(\varphi_4)$



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where $\beta > 0$



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- $B_{02}(\beta) + R_{34}(\varphi_1) + R_{56}(\varphi_2) + R_{78}(\varphi_3) + R_{9\natural}(\varphi_4)$
- $N_{+2} + R_{34}(\varphi_1) + R_{56}(\varphi_2) + R_{78}(\varphi_3) + R_{9\natural}(\varphi_4)$

where $\beta > 0$, $\varphi_1 \ge \varphi_2 \ge \cdots \ge \varphi_{k-1} \ge \varphi_k \ge 0$

• $\lambda + \tau \in \mathfrak{so}(1, 10) \oplus \mathbb{R}^{1, 10}$

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to get rid of component of τ in the image of $[\lambda, -]$

• the subgroups with everywhere spacelike orbits

 $\star \partial_z + R_{12}(\varphi_1) + R_{34}(\varphi_2) + R_{56}(\varphi_3) + R_{78}(\varphi_4)$

- the subgroups with everywhere spacelike orbits are generated by either
 - * $\partial_z + R_{12}(\varphi_1) + R_{34}(\varphi_2) + R_{56}(\varphi_3) + R_{78}(\varphi_4)$; or * $\partial_z + N_{+2} + R_{34}(\varphi_2) + R_{56}(\varphi_3) + R_{78}(\varphi_4)$

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where $\varphi_1 \ge \varphi_2 \ge \varphi_3 \ge \varphi_4 \ge 0$

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$$\partial_z + R_{12}(\varphi_1) + R_{34}(\varphi_2) + R_{56}(\varphi_3) + R_{78}(\varphi_4)$$
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• both are $\cong \mathbb{R}$

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- both are $\cong \mathbb{R}$
- the former gives rise to fluxbranes

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- both are $\cong \mathbb{R}$
- the former gives rise to fluxbranes and the latter to nullbranes

Adapted coordinates

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• start with metric in flat coordinates y, z

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$$ds^2 = 2|d\boldsymbol{y}|^2 + dz^2$$
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$$\xi = \partial_z + \lambda$$

• start with metric in flat coordinates y, z

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$$\xi = \partial_z + \lambda = U \,\partial_z \, U^{-1}$$

• start with metric in flat coordinates y, z

$$ds^2 = 2|d\boldsymbol{y}|^2 + dz^2$$

$$\xi = \partial_z + \lambda = U \, \partial_z \, U^{-1}$$
 with $U = \exp(-z\lambda)$

 $\boldsymbol{x} = U \, \boldsymbol{y}$

$$\boldsymbol{x} = U \boldsymbol{y} = \exp(-zB)\boldsymbol{y}$$

$$oldsymbol{x} = U oldsymbol{y} = \exp(-zB)oldsymbol{y}$$
 where $\lambda oldsymbol{y} = Boldsymbol{y}$

 $\boldsymbol{x} = U \, \boldsymbol{y} = \exp(-zB) \boldsymbol{y}$ where $\lambda \boldsymbol{y} = B \boldsymbol{y}$

whence $\xi x = 0$

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whence $\xi x = 0$

• rewrite the metric in terms of *x*:

$$ds^{2} = \Lambda (dz + A)^{2} + |d\boldsymbol{x}|^{2} - \Lambda A^{2}$$

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• rewrite the metric in terms of \boldsymbol{x} :

$$ds^{2} = \Lambda (dz + A)^{2} + |d\boldsymbol{x}|^{2} - \Lambda A^{2}$$

where

 $\star \Lambda = 1 + |B\boldsymbol{x}|^2$

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• rewrite the metric in terms of *x*:

$$ds^{2} = \Lambda (dz + A)^{2} + |d\boldsymbol{x}|^{2} - \Lambda A^{2}$$

where

 $\star \Lambda = 1 + |B\boldsymbol{x}|^2$ $\star A = \Lambda^{-1} B\boldsymbol{x} \cdot d\boldsymbol{x}$

$$ds^{2} = e^{4\Phi/3}(dz + A)^{2} + e^{-2\Phi/3}h$$

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 \star dilaton

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* dilaton: $\Phi = \frac{3}{4}\log(1+|B\boldsymbol{x}|^2)$

$$ds^{2} = e^{4\Phi/3}(dz + A)^{2} + e^{-2\Phi/3}h$$

we read off the IIA fields * dilaton: $\Phi = \frac{3}{4} \log(1 + |Bx|^2)$ * RR 1-form potential

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we read off the IIA fields \star dilaton: $\Phi = \frac{3}{4} \log(1 + |B\boldsymbol{x}|^2)$

★ RR 1-form potential:

$$A = \frac{B\boldsymbol{x} \cdot d\boldsymbol{x}}{1 + |B\boldsymbol{x}|^2}$$

$$ds^{2} = e^{4\Phi/3}(dz + A)^{2} + e^{-2\Phi/3}h$$

we read off the IIA fields * dilaton: $\Phi = \frac{3}{4} \log(1 + |Bx|^2)$ * RR 1-form potential:

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***** string frame metric

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we read off the IIA fields * dilaton: $\Phi = \frac{3}{4} \log(1 + |Bx|^2)$ * RR 1-form potential:

$$A = \frac{B\boldsymbol{x} \cdot d\boldsymbol{x}}{1 + |B\boldsymbol{x}|^2}$$

***** string frame metric:

$$h = \Lambda^{1/2} |d\boldsymbol{x}|^2 - \Lambda^{-1/2} (B\boldsymbol{x} \cdot d\boldsymbol{x})^2$$

$$B = \begin{pmatrix} 0 & -\varphi_1 & & & 0 & u \\ \varphi_1 & 0 & & & 0 & 0 \\ & & 0 & -\varphi_2 & & & & \\ & & & \varphi_2 & 0 & & & & \\ & & & & 0 & -\varphi_3 & & & \\ & & & & & & \varphi_3 & 0 & & \\ & & & & & & & \varphi_4 & 0 & \\ & & & & & & & & \varphi_4 & 0 & \\ & & & & & & & & & \varphi_4 & 0 & \\ & & & & & & & & & & & \end{pmatrix}$$

where

where either

 $\star u = 0$

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 $\star u = 0$ (generalised fluxbranes)

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★ u = 0 (generalised fluxbranes); or ★ u = 1 and $\varphi_1 = 0$

where either

* u = 0 (generalised fluxbranes); or * u = 1 and $\varphi_1 = 0$ (generalised nullbranes)



e.g.,

$$\star$$
 $u=0$, $arphi_2=arphi_3=arphi_4=0$

$\star u = 0, \ \varphi_2 = \varphi_3 = \varphi_4 = 0 \implies \text{flux-sevenbrane}$ [Costa-Gutperle, hep-th/0012072]

e.g., $\star \ u = 0, \ \varphi_2 = \varphi_3 = \varphi_4 = 0 \implies \text{flux-sevenbrane}$ [Costa-Gutperle, hep-th/0012072] $\star \ u = 0$

e.g.,

* u = 0, $\varphi_2 = \varphi_3 = \varphi_4 = 0 \implies \text{flux-sevenbrane}$ [Costa-Gutperle, hep-th/0012072]

 \star u=0, $arphi_1=arphi_2$

* u = 0, $\varphi_2 = \varphi_3 = \varphi_4 = 0 \implies$ flux-sevenbrane [Costa-Gutperle, hep-th/0012072] * u = 0, $\varphi_1 = \varphi_2$, $\varphi_3 = \varphi_4 = 0 \implies$ half-BPS flux-fivebrane [Gutperle-Strominger, hep-th/0104136]
$\begin{array}{l} \star \ u = 0, \ \varphi_2 = \varphi_3 = \varphi_4 = 0 \implies {\rm flux-sevenbrane} \\ & \ \ \left[{\rm Costa-Gutperle, \ hep-th/0012072} \right] \\ \star \ u = 0, \ \varphi_1 = \varphi_2, \ \varphi_3 = \varphi_4 = 0 \implies {\rm half-BPS \ flux-fivebrane} \\ & \ \ \left[{\rm Gutperle-Strominger, \ hep-th/0104136} \right] \\ \star \ u = 0 \end{array}$

$$\begin{array}{l} \star \; u = 0, \; \varphi_2 = \varphi_3 = \varphi_4 = 0 \implies {\sf flux-sevenbrane} \\ & [{\sf Costa-Gutperle, \, hep-th/0012072}] \\ \star \; u = 0, \; \varphi_1 = \varphi_2, \; \varphi_3 = \varphi_4 = 0 \implies {\sf half-BPS \ flux-fivebrane} \\ & [{\sf Gutperle-Strominger, \, hep-th/0104136}] \\ \star \; u = 0, \; \varphi_1 = \varphi_2 + \varphi_3 \end{array}$$

 $\begin{array}{l} \star \; u = 0, \; \varphi_2 = \varphi_3 = \varphi_4 = 0 \; \Longrightarrow \; {\rm flux-sevenbrane} \\ & \ \ \left[{\rm Costa-Gutperle, \; hep-th/0012072} \right] \\ \star \; u = 0, \; \varphi_1 = \varphi_2, \; \varphi_3 = \varphi_4 = 0 \; \Longrightarrow \; {\rm half-BPS \; flux-fivebrane} \\ & \ \ \left[{\rm Gutperle-Strominger, \; hep-th/0104136} \right] \\ \star \; u = 0, \; \varphi_1 = \varphi_2 + \varphi_3, \; \varphi_4 = 0 \end{array}$

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$$\begin{array}{l} \star \; u = 0, \; \varphi_2 = \varphi_3 = \varphi_4 = 0 \implies {\rm flux-sevenbrane} \\ [{\rm Costa-Gutperle, \; hep-th/0012072}] \\ \star \; u = 0, \; \varphi_1 = \varphi_2, \; \varphi_3 = \varphi_4 = 0 \implies {\rm half-BPS \; flux-fivebrane} \\ [{\rm Gutperle-Strominger, \; hep-th/0104136}] \\ \star \; u = 0, \; \varphi_1 = \varphi_2 + \varphi_3, \; \varphi_4 = 0 \implies \frac{1}{4} \text{-BPS \; flux-threebrane} \\ \star \; u = 0 \end{array}$$

 $\begin{array}{l} \star \; u = 0, \; \varphi_2 = \varphi_3 = \varphi_4 = 0 \; \Longrightarrow \; {\rm flux-sevenbrane} \\ & [{\rm Costa-Gutperle, \; hep-th/0012072}] \\ \star \; u = 0, \; \varphi_1 = \varphi_2, \; \varphi_3 = \varphi_4 = 0 \; \Longrightarrow \; {\rm half-BPS \; flux-fivebrane} \\ & [{\rm Gutperle-Strominger, \; hep-th/0104136}] \\ \star \; u = 0, \; \varphi_1 = \varphi_2 + \varphi_3, \; \varphi_4 = 0 \; \Longrightarrow \; \frac{1}{4} \text{-BPS \; flux-threebrane} \\ \star \; u = 0, \; \varphi_1 - \varphi_2 = \varphi_3 \pm \varphi_4 \end{array}$

 $\begin{array}{l} \star \ u = 0, \ \varphi_2 = \varphi_3 = \varphi_4 = 0 \implies {\sf flux-sevenbrane} \\ [{\sf Costa-Gutperle, \ hep-th/0012072}] \\ \star \ u = 0, \ \varphi_1 = \varphi_2, \ \varphi_3 = \varphi_4 = 0 \implies {\sf half-BPS \ flux-fivebrane} \\ [{\sf Gutperle-Strominger, \ hep-th/0104136}] \\ \star \ u = 0, \ \varphi_1 = \varphi_2 + \varphi_3, \ \varphi_4 = 0 \implies \frac{1}{4} {\sf -BPS \ flux-threebrane} \\ \star \ u = 0, \ \varphi_1 - \varphi_2 = \varphi_3 \pm \varphi_4 \implies \frac{1}{8} {\sf -BPS \ flux-string} \\ [{\sf Uranga, \ hep-th/0108196}] \end{array}$

e.g.,

$$\begin{array}{l} \star \ u = 0, \ \varphi_2 = \varphi_3 = \varphi_4 = 0 \implies {\sf flux-sevenbrane} \\ [{\sf Costa-Gutperle, \ hep-th/0012072}] \\ \star \ u = 0, \ \varphi_1 = \varphi_2, \ \varphi_3 = \varphi_4 = 0 \implies {\sf half-BPS \ flux-fivebrane} \\ [{\sf Gutperle-Strominger, \ hep-th/0104136}] \\ \star \ u = 0, \ \varphi_1 = \varphi_2 + \varphi_3, \ \varphi_4 = 0 \implies \frac{1}{4} \cdot {\sf BPS \ flux-threebrane} \\ \star \ u = 0, \ \varphi_1 - \varphi_2 = \varphi_3 \pm \varphi_4 \implies \frac{1}{8} \cdot {\sf BPS \ flux-string} \\ [{\sf Uranga, \ hep-th/0108196}] \\ \star \ u = 0 \end{array}$$

$$\begin{array}{l} \star \; u = 0, \; \varphi_2 = \varphi_3 = \varphi_4 = 0 \implies {\rm flux-sevenbrane} \\ [{\rm Costa-Gutperle, \; hep-th/0012072}] \\ \star \; u = 0, \; \varphi_1 = \varphi_2, \; \varphi_3 = \varphi_4 = 0 \implies {\rm half-BPS \; flux-fivebrane} \\ [{\rm Gutperle-Strominger, \; hep-th/0104136}] \\ \star \; u = 0, \; \varphi_1 = \varphi_2 + \varphi_3, \; \varphi_4 = 0 \implies \frac{1}{4} \cdot {\rm BPS \; flux-threebrane} \\ \star \; u = 0, \; \varphi_1 - \varphi_2 = \varphi_3 \pm \varphi_4 \implies \frac{1}{8} \cdot {\rm BPS \; flux-string} \\ [{\rm Uranga, \; hep-th/0108196}] \\ \star \; u = 0, \; \varphi_1 = \varphi_2 \end{array}$$

$$\begin{array}{l} \star \ u = 0, \ \varphi_2 = \varphi_3 = \varphi_4 = 0 \implies {\sf flux-sevenbrane} \\ [{\sf Costa-Gutperle, \ hep-th/0012072}] \\ \star \ u = 0, \ \varphi_1 = \varphi_2, \ \varphi_3 = \varphi_4 = 0 \implies {\sf half-BPS \ flux-fivebrane} \\ [{\sf Gutperle-Strominger, \ hep-th/0104136}] \\ \star \ u = 0, \ \varphi_1 = \varphi_2 + \varphi_3, \ \varphi_4 = 0 \implies \frac{1}{4} - {\sf BPS \ flux-threebrane} \\ \star \ u = 0, \ \varphi_1 - \varphi_2 = \varphi_3 \pm \varphi_4 \implies \frac{1}{8} - {\sf BPS \ flux-string} \\ [{\sf Uranga, \ hep-th/0108196}] \\ \star \ u = 0, \ \varphi_1 = \varphi_2, \ \varphi_3 = \varphi_4 \end{array}$$

$$u=0$$
, $arphi_1=arphi_2$, $arphi_3=arphi_4$

 $\begin{array}{l} \star \ u = 0, \ \varphi_2 = \varphi_3 = \varphi_4 = 0 \implies {\sf flux-sevenbrane} \\ [{\sf Costa-Gutperle, \ hep-th/0012072}] \\ \star \ u = 0, \ \varphi_1 = \varphi_2, \ \varphi_3 = \varphi_4 = 0 \implies {\sf half-BPS \ flux-fivebrane} \\ [{\sf Gutperle-Strominger, \ hep-th/0104136}] \\ \star \ u = 0, \ \varphi_1 = \varphi_2 + \varphi_3, \ \varphi_4 = 0 \implies \frac{1}{4} \text{-BPS \ flux-threebrane} \\ \star \ u = 0, \ \varphi_1 - \varphi_2 = \varphi_3 \pm \varphi_4 \implies \frac{1}{8} \text{-BPS \ flux-string} \\ [{\sf Uranga, \ hep-th/0108196}] \\ \star \ u = 0, \ \varphi_1 = \varphi_2, \ \varphi_3 = \varphi_4 \implies \frac{1}{4} \text{-BPS \ flux-string} \end{array}$

 $\begin{array}{l} \star \ u = 0, \ \varphi_2 = \varphi_3 = \varphi_4 = 0 \implies {\sf flux-sevenbrane} \\ [{\sf Costa-Gutperle, \ hep-th/0012072}] \\ \star \ u = 0, \ \varphi_1 = \varphi_2, \ \varphi_3 = \varphi_4 = 0 \implies {\sf half-BPS \ flux-fivebrane} \\ [{\sf Gutperle-Strominger, \ hep-th/0104136}] \\ \star \ u = 0, \ \varphi_1 = \varphi_2 + \varphi_3, \ \varphi_4 = 0 \implies \frac{1}{4} \text{-BPS \ flux-threebrane} \\ \star \ u = 0, \ \varphi_1 - \varphi_2 = \varphi_3 \pm \varphi_4 \implies \frac{1}{8} \text{-BPS \ flux-string} \\ [{\sf Uranga, \ hep-th/0108196}] \\ \star \ u = 0, \ \varphi_1 = \varphi_2, \ \varphi_3 = \varphi_4 \implies \frac{1}{4} \text{-BPS \ flux-string} \\ [{\sf Wanga, \ hep-th/0108196}] \\ \star \ u = 1 \end{array}$

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[Horowitz–Steif (1991)]

End of first lecture

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 \mathbb{CP}^2 is not even spin! [Duff-Lü-Pope, hep-th/9704186,9803061]



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35

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[e.g., Hull hep-th/0305039]

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 $\xi = \partial_z + R_{12}(\varphi_1) + R_{34}(\varphi_2) + R_{56}(\varphi_3) + R_{78}(\varphi_4)$

• e.g., fluxbranes

 $\boldsymbol{\varsigma}$

c)

JU



• for generic φ_i , there are no invariant Killing spinors

 $\star \varphi_1 - \varphi_2 - \varphi_3 \pm \varphi_4 = 0$

$$\star \varphi_1 - \varphi_2 - \varphi_3 \pm \varphi_4 = 0 \Longrightarrow \nu = \frac{1}{8}$$

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$$\varphi_1 = \varphi_2, \ \varphi_3 = \varphi_4 \implies \nu = \frac{1}{4}$$

$$\varphi_1 - \varphi_2 - \varphi_3 = 0 = \varphi_4 \implies \nu = \frac{1}{4}$$

$$\varphi_1 = \varphi_2 = \varphi_3 = \varphi_4$$

$$\begin{array}{l} \star \ \varphi_1 - \varphi_2 - \varphi_3 \pm \varphi_4 = 0 \Longrightarrow \ \nu = \frac{1}{8} \\ \star \ \varphi_1 = \varphi_2, \ \varphi_3 = \varphi_4 \Longrightarrow \ \nu = \frac{1}{4} \\ \star \ \varphi_1 - \varphi_2 - \varphi_3 = 0 = \varphi_4 \Longrightarrow \ \nu = \frac{1}{4} \\ \star \ \varphi_1 = \varphi_2 = \varphi_3 = \varphi_4 \Longrightarrow \ \nu = \frac{3}{8} \\ \star \ \varphi_1 = \varphi_2, \ \varphi_3 = \varphi_4 = 0 \end{array}$$

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• e.g., nullbranes

 $\xi = \partial_z + N_{+2} + R_{34}(\varphi_2) + R_{56}(\varphi_3) + R_{78}(\varphi_4)$

• e.g., nullbranes


• N_{+2} is nilpotent

$$\begin{split} \xi &= \partial_z + N_{+2} + R_{34}(\varphi_2) + R_{56}(\varphi_3) + R_{78}(\varphi_4) \\ \implies \\ \mathcal{L}_{\xi} &= \frac{1}{2}\Gamma_{+2} + \frac{1}{2}(\varphi_2\Gamma_{34} + \varphi_3\Gamma_{56} + \varphi_4\Gamma_{78}) \end{split}$$

• N_{+2} is nilpotent, whereas $\frac{1}{2}(\varphi_2\Gamma_{34}+\varphi_3\Gamma_{56}+\varphi_4\Gamma_{78})$ is semisimple and commutes with it

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- $\ker N_{+2} = \ker \Gamma_+$, and this simply halves the number of supersymmetries



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• $A(r), B(r), C(r) \rightarrow 0$ as $r \rightarrow \infty \implies$ asymptotic to $(\mathbb{R}^{1,D-1}, F = 0)$



 $G = \left(\mathrm{SO}(1, p) \ltimes \mathbb{R}^{1, p} \right)$

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 $\operatorname{dvol}(\mathbb{R}^{1,p}) \cdot \varepsilon_{\infty} = \varepsilon_{\infty}$





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 \implies the M2-brane is half-BPS

• action of G on Killing spinors also induced from asymptotic limit
$\star \xi$ a Killing vector

- $\star \xi$ a Killing vector
- $\star \varepsilon$ a Killing spinor

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(independent of r) \star compute it at $r \to \infty$ \implies the classification problem reduces to the flat case

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• e.g., ξ may be everywhere spacelike in the brane metric

- \star G is <u>not</u> the full asymptotic symmetry group, and
- \star brane metric is <u>not</u> flat, so causal properties of Killing vectors may differ from $\mathbb{R}^{1,D-1}$
- e.g., ξ may be everywhere spacelike in the brane metric, but its asymptotic value may not be everywhere spacelike in the flat metric

• metric

47

• metric:

 $V^{-2/3}ds^2(\mathbb{R}^{1,2}) + V^{1/3}ds^2(\mathbb{R}^8)$

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$$\xi = \tau_{\parallel}$$

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• Killing vector:

 $\xi = \tau_{\parallel} + \rho_{\parallel}$

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• Killing vector:

 $\xi = \tau_{\parallel} + \rho_{\parallel} + \rho_{\perp}$

mutually orthogonal



$$\|\xi\|^{2} = V^{-2/3} \left(|\tau|^{2} + |\rho_{\parallel}|^{2} \right) + V^{1/3} |\rho_{\perp}|^{2}$$

$$\begin{aligned} \|\xi\|^2 &= V^{-2/3} \left(|\tau|^2 + |\rho_{\parallel}|^2 \right) + V^{1/3} |\rho_{\perp}|^2 \\ &= V^{-2/3} \left(|\tau|^2 + |\rho_{\parallel}|^2 \right) + V^{1/3} r^2 |\rho_{\perp}|_S^2 \end{aligned}$$

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• odd-dimensional spheres can be combed $\implies m^2$ can be > 0

• f(r) has a minimum at $r_0 > 0$

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$$\star$$
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$$\star$$
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$$V(r_0)r_0^8 = -rac{2|Q|}{m^2}| au|^2$$

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$$\|\xi\|^2 \ge V(r_0)^{-2/3} |\tau|^2 + V(r_0)^{1/3} r_0^2 m^2$$

which is > 0 provided that

$$|\tau|^2 > -\frac{3}{2}m^2(2|Q|)^{1/3}$$

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End of second lecture

purely geometric backgrounds

purely geometric backgrounds, with product geometry

 $(M^4 imes N^7, g \oplus h)$

purely geometric backgrounds, with product geometry

 $(M^4 imes N^7, g \oplus h)$ and $F \propto \operatorname{dvol}_g$

• purely geometric backgrounds, with product geometry $(M^4 imes N^7, g \oplus h)$ and $F \propto {
m dvol}_g$

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- supersymmetry $\iff (M,g)$ and (N,h) admit geometric Killing spinors

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- field equations $\iff (M,g)$ and (N,h) are Einstein
- supersymmetry $\iff (M,g)$ and (N,h) admit geometric Killing spinors:

 $\nabla_a \varepsilon = \lambda \Gamma_a \varepsilon$ where $\lambda \in \mathbb{R}^{\times}$

• (M,g) admits geometric Killing spinors

 $\widehat{M} = \mathbb{R}^+ \times M$

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[Bär (1993), Kath (1999)]

$$\widehat{M} = \mathbb{R}^+ imes M$$
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admits parallel spinors: $\nabla \hat{\varepsilon} = 0$

[Bär (1993), Kath (1999)]

• equivariant under the isometry group G of (M, g)[hep-th/9902066]



• (M,g) riemannian $\implies (\widehat{M},\widehat{g})$ riemannian

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- $(M^{1,n-1},g)$ lorentzian

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again the problem reduces to one of flat spaces!

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• one-parameter subgroups \leftrightarrow projectivised adjoint orbits of $\mathfrak{so}(2,p)$ under SO(2,p)







We can still use the lorentzian elementary blocks



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• (0,2)



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• (0,2) and also (2,0)



• (0,2) and also (2,0), $\mu(x) = x^2 + \varphi^2$





 $B^{(0,2)}(\varphi)$

We play again but with a bigger set! We can still use the lorentzian elementary blocks: • (0,2) and also (2,0), $\mu(x) = x^2 + \varphi^2$, rotation

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$$B^{(0,2)}(\varphi) = B^{(2,0)}(\varphi) = \begin{bmatrix} 0 & \varphi \\ -\varphi & 0 \end{bmatrix}$$

• (1,1)

59

• (1,1),
$$\mu(x) = x^2 - \beta^2$$

• (1,1),
$$\mu(x)=x^2-\beta^2$$
, boost

•
$$(1,1)$$
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 $B^{(1,1)}$

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$$B^{(1,2)} = B^{(2,1)} = \begin{bmatrix} 0 & -1 & 0 \\ 1 & 0 & -1 \\ 0 & 1 & 0 \end{bmatrix}$$



•
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, $\mu(x) = x^2$

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 $B_{\pm}^{(2,2)}$

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$$B_{\pm}^{(2,2)} = \begin{bmatrix} 0 & \mp 1 & 1 & 0 \\ \pm 1 & 0 & 0 & \mp 1 \\ -1 & 0 & 0 & 1 \\ 0 & \pm 1 & -1 & 0 \end{bmatrix}$$



•
$$(2,2)$$
 , $\mu(x)=(x^2\!-\!\beta^2)^2$

• (2, 2), $\mu(x) = (x^2 - \beta^2)^2$, deformation of $B_{\pm}^{(2,2)}$ by a (anti)selfdual boost

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The associated discrete quotient of AdS_3
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The associated discrete quotient of AdS_3 yields the extremal BTZ black hole; the non-extremal black hole

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The associated discrete quotient of AdS_3 yields the extremal BTZ black hole; the non-extremal black hole is obtained from $B^{(1,1)}(\beta_1) \oplus B^{(1,1)}(\beta_2)$, for $|\beta_1| \neq |\beta_2|$



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$$B_{\pm}^{(2,2)}(\varphi) = \begin{bmatrix} 0 & \mp 1 \pm \varphi & 1 & 0 \\ \pm 1 \mp \varphi & 0 & 0 & \mp 1 \\ -1 & 0 & 0 & 1 + \varphi \\ 0 & \pm 1 & -1 - \varphi & 0 \end{bmatrix}$$



• (2,2),
$$\mu(x) = (x^2 + \beta^2 + \varphi^2) - 4\beta^2 x^2$$

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•
$$(2,3)$$
, $\mu(x) = x^{\mathrm{s}}$

• (2,3), $\mu(x) = x^5$, deformation of $B^{(2,2)}_+$ by a null rotation in a perpendicular direction

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 $B^{(2,3)}$

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$$B^{(2,3)} = \begin{bmatrix} 0 & 1 & -1 & 0 & -1 \\ -1 & 0 & 0 & 1 & 0 \\ 1 & 0 & 0 & -1 & 0 \\ 0 & -1 & 1 & 0 & -1 \\ 1 & 0 & 0 & 1 & 0 \end{bmatrix}$$



•
$$(2,4)$$
, $\mu(x) = (x^2 + \varphi^2)^3$

 $B^{(2,4)}_{\pm}(arphi)$

$$B_{\pm}^{(2,4)}(\varphi) = \begin{bmatrix} 0 & \mp \varphi & 0 & 0 & -1 & 0 \\ \pm \varphi & 0 & 0 & 0 & 0 & \mp 1 \\ 0 & 0 & 0 & \varphi & -1 & 0 \\ 0 & 0 & -\varphi & 0 & 0 & -1 \\ 1 & 0 & 1 & 0 & 0 & \varphi \\ 0 & \pm 1 & 0 & 1 & -\varphi & 0 \end{bmatrix}$$

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• and that's all!

• Killing vectors on $AdS_{1+p} \times S^q$ decompose

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$$\xi = \xi_A + \xi_S$$

whose norms add

 $\|\xi\|^2 = \|\xi_A\|^2 + \|\xi_S\|^2$





$R^2 M^2 \ge \|\xi_S\|^2$

•
$$S^q$$
 is compact \Longrightarrow

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• ξ can be everywhere spacelike on $AdS_{1+p} \times S^{2k+1}$, even if ξ_A is not spacelike everywhere
• S^q is compact \Longrightarrow

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• ξ can be everywhere spacelike on $AdS_{1+p} \times S^{2k+1}$, even if ξ_A is not spacelike everywhere, provided that $\|\xi_A\|^2$ is <u>bounded below</u>

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- it is convenient to distinguish Killing vectors according to norm

everywhere non-negative norm

• everywhere non-negative norm:

 $\star \oplus_i B^{(0,2)}(\varphi_i)$

- everywhere non-negative norm:
 - $\star \oplus_i B^{(0,2)}(\varphi_i)$ $\star B^{(1,1)}(\beta_1) \oplus B^{(1,1)}(\beta_2) \oplus_i B^{(0,2)}(\varphi_i)$

- $\begin{array}{l} \star \oplus_{i} B^{(0,2)}(\varphi_{i}) \\ \star B^{(1,1)}(\beta_{1}) \oplus B^{(1,1)}(\beta_{2}) \oplus_{i} B^{(0,2)}(\varphi_{i}), \text{ if } |\beta_{1}| = |\beta_{2}| \end{array}$
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 - $\star \oplus_i B^{(0,2)}(\varphi_i)$
 - $\star B^{(1,1)}(eta_1) \oplus B^{(1,1)}(eta_2) \oplus_i B^{(0,2)}(arphi_i)$, if $|eta_1| = |eta_2|$
 - $\star B^{(1,2)} \oplus_i B^{(0,2)}(\varphi_i)$
 - $\star B^{(1,2)} \oplus \overline{B^{(1,2)} \oplus_i B^{(0,2)}}(\varphi_i)$

- everywhere non-negative norm:

• everywhere non-negative norm:

norm bounded below

- everywhere non-negative norm:
- norm bounded below:
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- everywhere non-negative norm:
- norm bounded below:
 - $\star B^{(2,0)}(\varphi) \oplus_i \overline{B^{(0,2)}(\varphi_i)}$, if p is even and $|\varphi_i| \ge \varphi > 0$ for all i

- everywhere non-negative norm:
- norm bounded below:
 - $\begin{array}{l} \star \ B^{(2,0)}(\varphi) \oplus_i \overline{B^{(0,2)}(\varphi_i)}, \ \overline{\text{if } p \text{ is even and } |\varphi_i| \ge \varphi > 0 \text{ for all } i \\ \star \ B^{(2,2)}_{\pm}(\varphi) \oplus_i \overline{B^{(0,2)}(\varphi_i)} \end{array}$

- everywhere non-negative norm:
 - $\begin{array}{l} \star \oplus_{i} B^{(0,2)}(\varphi_{i}) \\ \star B^{(1,1)}(\beta_{1}) \oplus B^{(1,1)}(\beta_{2}) \oplus_{i} B^{(0,2)}(\varphi_{i}), \text{ if } |\beta_{1}| = |\beta_{2}| \\ \star B^{(1,2)} \oplus_{i} B^{(0,2)}(\varphi_{i}) \\ \star B^{(1,2)} \oplus B^{(1,2)} \oplus_{i} B^{(0,2)}(\varphi_{i}) \\ \star B^{(2,2)}_{\pm} \oplus_{i} B^{(0,2)}(\varphi_{i}) \end{array}$
- norm bounded below:
 - * $B^{(2,0)}(\varphi) \oplus_i B^{(0,2)}(\varphi_i)$, if p is even and $|\varphi_i| \ge \varphi > 0$ for all i* $B^{(2,2)}_{\pm}(\varphi) \oplus_i B^{(0,2)}(\varphi_i)$, if $|\varphi_i| \ge |\varphi| \ge 0$ for all i

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- everywhere non-negative norm:
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 - $\begin{array}{l} \star \ B^{(2,0)}(\varphi) \oplus_i B^{(0,2)}(\overline{\varphi_i}), \text{ if } p \text{ is even and } |\varphi_i| \geq \overline{\varphi} > 0 \text{ for all } i \\ \star \ B^{(2,2)}_{\pm}(\varphi) \oplus_i B^{(0,2)}(\varphi_i), \text{ if } |\varphi_i| \geq |\varphi| \geq 0 \text{ for all } i \end{array}$
- arbitrarily negative norm: the rest!

 $\star B^{(1,1)}(eta_1) \oplus B^{(1,1)}(eta_2) \oplus_i B^{(0,2)}(arphi_i)$

 $\star B^{(1,1)}(\beta_1) \oplus B^{(1,1)}(\beta_2) \oplus_i B^{(0,2)}(\varphi_i)$, unless $|\beta_1| = |\beta_2| > 0$

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Some of these give rise to higher-dimensional BTZ-like black holes

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Some of these give rise to higher-dimensional BTZ-like black holes: quotient only a part of AdS and check that the boundary thus introduced lies behind a horizon.

Discrete quotients with CTCs

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 $\gamma = \exp(LX)$

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 $\|\dot{c}\|^2$

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 - $\star N$ is Einstein with positive cosmological constant and Einstein
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- geometrical CTCs are also natural in certain kinds of supersymmetric Freund–Rubin backgrounds $M \times N$, where M is lorentzian Einstein–Sasaki: timelike circle bundles over Kähler manifolds

- there are many families of smooth supersymmetric reductions of ${\rm AdS}_4 \times S^7$

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- $\frac{3}{4}$ -BPS $AdS_4 \times \mathbb{CP}^3$ background of IIA

[Duff-Lü-Pope, hep-th/9704186]

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[Duff-Lü-Pope, hep-th/9704186]

• $\frac{9}{16}$ -BPS IIA backgrounds

- there are many families of smooth supersymmetric reductions of $AdS_4 \times S^7$, $S^4 \times AdS_7$, $AdS_5 \times S^5$, and $AdS_3 \times S^3 \times \mathbb{R}^4$.
- $\frac{3}{4}$ -BPS AdS₄ × \mathbb{CP}^3 background of IIA [Duff-Lü-Pope, hep-th/9704186]
- $\frac{9}{16}$ -BPS IIA backgrounds: reductions of $AdS_4 \times S^7$

- there are many families of smooth supersymmetric reductions of $AdS_4 \times S^7$, $S^4 \times AdS_7$, $AdS_5 \times S^5$, and $AdS_3 \times S^3 \times \mathbb{R}^4$.
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[Duff-Lü-Pope, hep-th/9704186]

• $\frac{9}{16}$ -BPS IIA backgrounds: reductions of $AdS_4 \times S^7$ by

 $B^{(2,2)}_{+}$

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- $\frac{3}{4}$ -BPS $AdS_4 \times \mathbb{CP}^3$ background of IIA [Duff-Lü-Pope, hep-th/9704186]
- $\frac{9}{16}$ -BPS IIA backgrounds: reductions of $AdS_4 \times S^7$ by

$$B_{+}^{(2,2)} \oplus \varphi(R_{12} + R_{34} + R_{56} - R_{78})$$



• a half-BPS IIA background: reduction of $S^4 imes { m AdS}^7$

 $B^{(1,2)}\oplus B^{(1,2)}$

 $B^{(1,2)} \oplus B^{(1,2)}$

• a family of half-BPS IIA backgrounds

 $B^{(1,2)} \oplus B^{(1,2)}$

• a family of half-BPS IIA backgrounds: reductions of $S^4 imes {
m AdS}^7$

$$B^{(1,2)}\oplus B^{(1,2)}$$

• a family of half-BPS IIA backgrounds: reductions of $S^4 imes ext{AdS}^7$ by $B^{(1,1)}(eta) \oplus B^{(1,1)}(eta) \oplus B^{(0,2)}(arphi) \oplus B^{(0,2)}(-arphi)$

$$B^{(1,2)}\oplus B^{(1,2)}$$

 a family of half-BPS IIA backgrounds: reductions of S⁴ × AdS⁷ by
 B^(1,1)(β) ⊕ B^(1,1)(β) ⊕ B^(0,2)(φ) ⊕ B^(0,2)(-φ)
 Both these half-BPS quotients are of the form S⁴ × (AdS₇ /Γ)

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 Both these half-BPS quotients are of the form S⁴ × (AdS₇ /Γ)
- a number of maximally supersymmetric reductions of $AdS_3 \times S^3$: near-horizon geometries of the supersymmetric rotating black holes, including over-rotating cases

reductions which break no supersymmetry are rare

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[Chamseddine-FO-Sabra, hep-th/0306278]
Thank you.