# **Quotients of M-theory vacua**

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.: can use orbifolds to study time-dependent phenomena in string theory; e.g., the nullbrane [Liu-Moore-Seiberg, hep-th/0206182]

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- (M, g, F, ...) with symmetry group G, Lie algebra  $\mathfrak{g}$
- $X, X' \in \mathfrak{g}$  give rise to equivalent quotients if and only if

$$X' = \lambda g X g^{-1} \qquad g \in G \quad \lambda \in \mathbb{R}^{\times}$$

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purely geometric backgrounds

purely geometric backgrounds, with product geometry

 $(M^4 imes N^7, g \oplus h)$ 

purely geometric backgrounds, with product geometry

 $(M^4 imes N^7, g \oplus h)$  and  $F \propto \operatorname{dvol}_g$ 

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• (M,g) admits geometric Killing spinors

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[Bär (1993), Kath (1999)]

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[Bär (1993), Kath (1999)]

• equivariant under the isometry group G of (M, g)[hep-th/9902066]



# • (M,g) riemannian $\implies (\widehat{M},\widehat{g})$ riemannian

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- the problem reduces to one of flat spaces!

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$$\mathbb{R}^{p+q} = \bigoplus_i \mathbb{V}_i$$

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• we need to determine the elementary blocks

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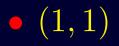
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$$B^{(1,1)}(\beta) = \begin{bmatrix} 0 & -\beta \\ \beta & 0 \end{bmatrix}$$

# • (1,2)

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• (1,2) and also (2,1)

• (1,2) and also (2,1), null rotation

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 $B^{(1,2)}$ 

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• (1,2) and also (2,1), null rotation

$$B^{(1,2)} = B^{(2,1)} = \begin{bmatrix} 0 & -1 & 0 \\ 1 & 0 & -1 \\ 0 & 1 & 0 \end{bmatrix}$$



• (2,2), "rotation" in a totally null plane

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 $B_{\pm}^{(2,2)}$ 

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$$B_{\pm}^{(2,2)} = \begin{bmatrix} 0 & \mp 1 & 1 & 0 \\ \pm 1 & 0 & 0 & \mp 1 \\ -1 & 0 & 0 & 1 \\ 0 & \pm 1 & -1 & 0 \end{bmatrix}$$



$$B_{\pm}^{(2,2)}(\beta > 0)$$

$$B_{\pm}^{(2,2)}(\beta > 0) = \begin{bmatrix} 0 & \mp 1 & 1 & -\beta \\ \pm 1 & 0 & \pm \beta & \mp 1 \\ -1 & \mp \beta & 0 & 1 \\ \beta & \pm 1 & -1 & 0 \end{bmatrix}$$

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[Bañados–Henneaux–Teitelboim–Zanelli, gr-qc/9302012]



$$B_{\pm}^{(2,2)}(\varphi)$$

$$B_{\pm}^{(2,2)}(\varphi) = \begin{bmatrix} 0 & \mp 1 \pm \varphi & 1 & 0 \\ \pm 1 \mp \varphi & 0 & 0 & \mp 1 \\ -1 & 0 & 0 & 1 + \varphi \\ 0 & \pm 1 & -1 - \varphi & 0 \end{bmatrix}$$



• (2,2), self-dual boost + antiself-dual rotation

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$$B^{(2,3)} = \begin{bmatrix} 0 & 1 & -1 & 0 & -1 \\ -1 & 0 & 0 & 1 & 0 \\ 1 & 0 & 0 & -1 & 0 \\ 0 & -1 & 1 & 0 & -1 \\ 1 & 0 & 0 & 1 & 0 \end{bmatrix}$$



• (2, 4), double null rotation + simultaneous rotation

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# $B^{(2,4)}_{\pm}(arphi)$

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$$B_{\pm}^{(2,4)}(\varphi) = \begin{bmatrix} 0 & \mp \varphi & 0 & 0 & -1 & 0 \\ \pm \varphi & 0 & 0 & 0 & 0 & \mp 1 \\ 0 & 0 & 0 & \varphi & -1 & 0 \\ 0 & 0 & -\varphi & 0 & 0 & -1 \\ 1 & 0 & 1 & 0 & 0 & \varphi \\ 0 & \pm 1 & 0 & 1 & -\varphi & 0 \end{bmatrix}$$

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and now we simply play



• Killing vectors on  $AdS_{1+p} \times S^q$  decompose

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whose norms add

 $\|\xi\|^2 = \|\xi_A\|^2 + \|\xi_S\|^2$ 





# $R^2 M^2 \ge \|\xi_S\|^2$

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$$S^q$$
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 $R^2 \overline{M^2} \ge \|\xi_S\|^2 \ge R^2 \overline{m^2}$ 

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- it is convenient to distinguish Killing vectors according to norm

everywhere non-negative norm

• everywhere non-negative norm:

 $\star \oplus_i B^{(0,2)}(\varphi_i)$ 

- everywhere non-negative norm:
  - $\star \oplus_i B^{(0,2)}(\varphi_i)$  $\star B^{(1,1)}(\beta_1) \oplus B^{(1,1)}(\beta_2) \oplus_i B^{(0,2)}(\varphi_i)$

- $\begin{array}{l} \star \oplus_{i} B^{(0,2)}(\varphi_{i}) \\ \star B^{(1,1)}(\beta_{1}) \oplus B^{(1,1)}(\beta_{2}) \oplus_{i} B^{(0,2)}(\varphi_{i}), \text{ if } |\beta_{1}| = |\beta_{2}| \end{array}$
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  - $\star B^{(1,2)} \oplus_i B^{(0,2)}(\varphi_i)$
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  - $\star B^{(2,0)}(\varphi) \oplus_i B^{(0,2)}(\varphi_i)$

- everywhere non-negative norm:
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  - $\star B^{(2,0)}(arphi) \oplus_i B^{(0,2)}(arphi_i)$ , if p is even

- everywhere non-negative norm:
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  - $\begin{array}{l} \star \ B^{(2,0)}(\varphi) \oplus_i \overline{B^{(0,2)}(\varphi_i)}, \ \overline{\text{if } p \text{ is even and } |\varphi_i| \ge \varphi > 0 \text{ for all } i \\ \star \ B^{(2,2)}_{\pm}(\varphi) \oplus_i \overline{B^{(0,2)}(\varphi_i)} \end{array}$

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- everywhere non-negative norm:
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- arbitrarily negative norm: the rest!

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Some of these give rise to higher-dimensional BTZ-like black holes

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Some of these give rise to higher-dimensional BTZ-like black holes: quotient only a part of AdS and check that the boundary thus introduced lies behind a horizon.

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 $\gamma = \exp(LX)$ 

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• geometrical CTCs are also natural in certain kinds of supersymmetric Freund–Rubin backgrounds  $M \times N$ , where M is lorentzian Einstein–Sasaki: timelike circle bundles over Kähler-Einstein spin manifolds

[FO-Leitner-Simón, to appear]

- there are many families of smooth supersymmetric reductions of  ${\rm AdS}_4 \times S^7$ 

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[Duff-Lü-Pope, hep-th/9704186]

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- $\frac{3}{4}$ -BPS AdS<sub>4</sub> ×  $\mathbb{CP}^3$  background of IIA [Duff-Lü-Pope, hep-th/9704186]
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- $\frac{3}{4}$ -BPS  $AdS_4 \times \mathbb{CP}^3$  background of IIA [Duff-Lü-Pope, hep-th/9704186]
- $\frac{9}{16}$ -BPS IIA backgrounds: reductions of  $AdS_4 \times S^7$  by

$$B_{+}^{(2,2)} \oplus \varphi(R_{12} + R_{34} + R_{56} - R_{78})$$



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[Chamseddine-FO-Sabra, hep-th/0306278]

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equations of motion

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## $\frac{1}{8}$ -BPS quotients

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- all are amenable to conformal field theory!

### Go forth and calculate!

# Thank you.

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