

Supersymmetry
&
lie algebra deformations

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works in progress

Part I

Lie algebras of Killing vectors

Question

What is the structure of the Lie algebra of isometries of (M, g) ?

$$\mathfrak{g} = \{ \xi \in \mathfrak{X}(M) \mid \mathcal{L}_\xi g = 0 \}$$



$$\nabla \xi \in \underline{\text{so}}(TM)$$

$$(\nabla_a \xi_b = -\nabla_b \xi_a)$$

$$(\nabla \xi)(Y) := \nabla_Y \xi$$

Notation

$$A_\xi = -\nabla \xi$$

Killing identity

$$\nabla_x A_\xi = R(\xi, x)$$

$\Rightarrow \xi \in \mathfrak{g}$ determined by $\xi_p, (\nabla \xi)_p$

Killing transport

[Kostant '55, Geroch '69]

$$E := TM \oplus \underline{\text{so}}(TM)$$

connection D

$$D_x \begin{pmatrix} \xi \\ A \end{pmatrix} = \begin{pmatrix} \nabla_x \xi + A(x) \\ \nabla_x A + R(x, \xi) \end{pmatrix}$$

$$D \begin{pmatrix} \xi \\ A \end{pmatrix} = 0 \quad \Leftrightarrow \quad \xi \text{ Killing}, \quad A = -\nabla \xi$$

$$\begin{array}{ll} p \in M & \xi \text{ KV} \quad (v, A) \in E_p \\ & \eta \text{ KV} \quad (w, B) \end{array}$$

What about $[\xi, \eta]$? $(Aw - Bv, AB - BA + R(v, w))$

$$\mathfrak{g} \subset E_p = T_p M \oplus \underline{\text{so}}(T_p M)$$

$$[(v, A), (w, B)] = (Aw - Bv, [A, B] + R(v, w))$$

Exercise: Prove directly the Jacobi identity.

(Hint: you will need both Bianchi identities.)

Grading:

$T_p M$	\oplus	$\underline{\text{so}}(T_p M)$	$R: \Lambda^2 T_p M \rightarrow \underline{\text{so}}(T_p M)$
-1	0	2	

$R = 0 \Rightarrow$ euclidean Lie algebra \mathfrak{e} \mathbb{Z} -graded

$R \neq 0 \Rightarrow$ isometry Lie algebra \mathfrak{g} filtered

$$\mathfrak{g} = F^{-1}\mathfrak{g} \supset F^0\mathfrak{g} \supset F^1\mathfrak{g} = 0 \quad \Rightarrow \quad \text{gr}(\mathfrak{g}) < \mathfrak{e}$$

$F^1 \underline{\text{so}}(T_p M)$

Filtered deformations of graded Lie algebras

Lie algebra on $V \Leftrightarrow \varphi \in \Lambda^2 V^* \otimes V \quad [\![\varphi, \varphi]\!] = 0$

\uparrow Nijenhuis-Richardson
bracket

Deformations $\bar{\varphi}(t) = \varphi + \underbrace{\sum_{n \geq 1} t^n \varphi_n}_{\varphi_+(t)}$ $\varphi_n \in \Lambda^2 V^* \otimes V$

$$0 = [\![\bar{\varphi}, \bar{\varphi}]\!] = [\![\varphi, \varphi]\!] + \underbrace{2[\![\varphi, \varphi_+]\!] + [\![\varphi_+, \varphi_+]\!]}$$

$$\partial \varphi_+ + \frac{1}{2} [\![\varphi_+, \varphi_+]\!] = 0 \quad \text{Maurer-Cartan}$$

$$\partial = [\![\varphi, -]\!] \quad \text{Chevalley-Eilenberg}$$

$$\partial : C^p(\mathfrak{g}; \mathfrak{g}) \rightarrow C^{p+1}(\mathfrak{g}; \mathfrak{g}) \quad \partial^2 = 0$$

$$\Lambda^p \mathfrak{g}^* \otimes \mathfrak{g}$$

$$\partial \varphi_+ + \frac{1}{2} [\bar{I}\varphi_+, \varphi_+] = 0 \quad \varphi_+ = \sum_{n \geq 1} t^n \varphi_n$$

$O(t)$ $\partial \varphi_+ = 0 \Rightarrow [\varphi_+] \in H^2(\mathfrak{g}; \mathfrak{g})$

$O(t^2)$ $\partial \varphi_2 + \frac{1}{2} [\varphi_+, \varphi_+] = 0 \Rightarrow [\bar{I}\varphi_+, \varphi_+] = 0 \in H^3(\mathfrak{g}; \mathfrak{g})$

etc...

$\therefore H^2(\mathfrak{g}; \mathfrak{g}) = \{ \text{infinitesimal deformations} \}$

$H^3(\mathfrak{g}; \mathfrak{g}) \ni \{ \text{obstructions} \}$

Remark

$$\mathfrak{g} \text{ Z-graded} \Leftrightarrow \deg \varphi = 0 \Leftrightarrow \deg \partial = 0$$

Can refine by degree

$$C^*(\mathfrak{g}; \mathfrak{g}) = \bigoplus_{\text{degree } d} C^{d,*}(\mathfrak{g}; \mathfrak{g})$$

Filtred deformations \equiv deformations in positive degree

e.g., $\deg R = 2$ in LA of isometries

Fact Infinitesimal filtred deformations are governed by

$$H^{d,2}(\mathfrak{g}_-; \mathfrak{g}) \quad d > 0$$

↪ negative degree
subalgebra

This is called (generalised) Spencer cohomology

Example

$$\mathfrak{g} = \underset{-1}{V} \oplus \underset{0}{\underline{\text{so}}(V)}$$

$$\mathfrak{g}_- = V$$

(V, η)

inner product space

e.g., $(T_p M, g_p)$

degree 1

$$C^{1,1} = V^* \otimes \underline{\text{so}}(V)$$

$$C^{1,2} = \Lambda^2 V^* \otimes V$$

$$\partial : C^{1,1} \xrightarrow{\cong} C^{1,2} \Rightarrow H^{1,2} = 0$$

degree 2

$$C^{2,2} = \Lambda^2 V^* \otimes \underline{\text{so}}(V)$$

$$R : \Lambda^2 V \rightarrow \underline{\text{so}}(V)$$

$$C^{2,3} = \Lambda^3 V^* \otimes V$$

$$(\partial R)(u, v, w) = R(u, v)w + R(v, w)u + R(w, u)v = 0$$

algebraic Bianchi identity

$$\therefore H^{2,2} = \{ \text{algebraic curvature tensors} \}$$

Summary

(M, g)

1. Killing vectors \longleftrightarrow parallel sections of $E = TM \oplus \text{so}(TM)$
2. Lie algebra of KVs is filtered and associated graded LA $< \epsilon$ Lie algebra of KVs of flat space
 \Rightarrow a filtered deformation
3. Infinitesimal filtered deformations of \mathfrak{g} are governed by "Spencer cohomology" $H^{q,2}(\mathfrak{g}; \mathfrak{g})$ $d > 0$
4. $H^{2,2}(\mathfrak{g}; \mathfrak{g}) \cong \{\text{algebraic curvature tensors}\}$
(but Jacobi requires differential Bianchi identity)

Part II

Supersymmetry

A "square root" of Killing transport

(M, g) spin $\$ \rightarrow M$ spinor bundle $(Cl(TM)\text{-modules})$

Promote $E = TM \oplus \underline{so}(TM)$ to a super VB

$$\mathcal{E} = \mathcal{E}_{\bar{s}} \oplus \mathcal{E}_s = E \oplus \$$$

Extend the Killing transport connection to a superconnection

$$D_x \begin{pmatrix} \xi \\ A \end{pmatrix} = \begin{pmatrix} \nabla_X \xi + A(X) \\ \nabla_X A + R(X, \xi) \end{pmatrix} \quad \sigma \in \Gamma(\$), \quad D_x \sigma = ?$$

Want: D -parallel sections of \mathcal{E} define a lie (super)algebra?

Lie superalgebras

$$\mathfrak{g} = \mathfrak{g}_{\bar{0}} \oplus \mathfrak{g}_{\bar{1}}$$

$$[-, -]: \mathfrak{g} \times \mathfrak{g} \rightarrow \mathfrak{g} \text{ (even)}$$

$$[x, y] = -(-1)^{|x||y|} [y, x]$$

$$[x, [y, z]] = [[x, y], z] + (-1)^{|x||y|} [y, [x, z]]$$

$[-, -]$ has three components :

$$\mathfrak{g}_{\bar{0}} \times \mathfrak{g}_{\bar{0}} \rightarrow \mathfrak{g}_{\bar{0}}$$

$$\mathfrak{g}_{\bar{0}} \times \mathfrak{g}_{\bar{1}} \rightarrow \mathfrak{g}_{\bar{1}}$$

$$\mathfrak{g}_{\bar{1}} \times \mathfrak{g}_{\bar{1}} \rightarrow \mathfrak{g}_{\bar{0}}$$

Jacobi has four components :

$$[\mathfrak{g}_{\bar{0}}, \mathfrak{g}_{\bar{0}}, \mathfrak{g}_{\bar{0}}] \Leftrightarrow \mathfrak{g}_{\bar{0}} \text{ is a LA}$$

$$[\mathfrak{g}_{\bar{0}}, \mathfrak{g}_{\bar{0}}, \mathfrak{g}_{\bar{1}}] \Leftrightarrow \mathfrak{g}_{\bar{1}} \text{ is } \mathfrak{g}_{\bar{0}}\text{-module}$$

$$[\mathfrak{g}_{\bar{0}}, \mathfrak{g}_{\bar{1}}, \mathfrak{g}_{\bar{1}}] \Leftrightarrow \mathfrak{g}_{\bar{1}} \times \mathfrak{g}_{\bar{1}} \rightarrow \mathfrak{g}_{\bar{0}} \text{ is } \mathfrak{g}_{\bar{0}}\text{-equivariant}$$

$$[\mathfrak{g}_{\bar{1}}, \mathfrak{g}_{\bar{1}}, \mathfrak{g}_{\bar{1}}]$$

Example

The Poincaré superalgebra

(V, η) Lorentzian vector space

$\underline{\text{so}}(V) \subset Cl(V)$ Clifford algebra $v^2 = -\eta(v, v)\mathbf{1}$

S a Clifford module

$\langle \cdot, \cdot \rangle$ an inner product on S (could be symplectic)

$$\langle s_1, v \cdot s_2 \rangle = \langle s_2, v \cdot s_1 \rangle \quad \forall v \in V, s_1, s_2 \in S$$

$$\Rightarrow K: S \times S \rightarrow V \quad \text{Dirac current}$$

$$\eta(K(s_1, s_2), v) = \langle s_1, v \cdot s_2 \rangle$$

$$\square = \underline{\text{so}}(V) \oplus S \oplus V \quad \square_{\bar{o}} = \underline{\text{so}}(V) \oplus V \quad \square_T = S$$

$$\mathbb{Z}\text{-graded} \quad \circ \quad -1 \quad -2$$

$$s_1, s_2 \in S \quad [s_1, s_2] = K(s_1, s_2)$$

Killing superalgebras

(M, g)

$$\mathcal{E} = E \oplus \mathbb{S}$$

D

Suppose

1) \mathbb{S} has an inner product $\langle \cdot, \cdot \rangle$ satisfying

$$\langle \sigma_1, x \cdot \sigma_2 \rangle = \langle \sigma_2, x \cdot \sigma_1 \rangle \quad \forall \sigma_1, \sigma_2 \in \Gamma(\mathbb{S}), x \in \mathfrak{X}(M)$$

2) the Dirac current $\kappa: \mathbb{S} \times \mathbb{S} \rightarrow TM$ obeys

$$\underbrace{D\sigma_1 = D\sigma_2 = 0}_{\text{Killing spinors}} \Rightarrow \kappa(\sigma_1, \sigma_2) \text{ is a KV}$$

3) $[\mathcal{L}_{\kappa(\sigma_1, \sigma_2)}, D] = 0$

spinorial Lie derivative $\mathcal{L}_\xi \sigma = \nabla_\xi \sigma + \rho(A_\xi) \cdot \sigma$

$$\rho: \mathfrak{so}(TM) \rightarrow \mathfrak{gl}(\mathbb{S})$$

Then let

$$\mathfrak{h} = \mathfrak{h}_{\bar{0}} \oplus \mathfrak{h}_{\bar{1}}$$

$$\mathfrak{h}_{\bar{0}} = \{ \xi \text{ KV} \mid [\mathcal{L}_\xi, D] = 0 \}$$

$$\mathfrak{h}_{\bar{1}} = \{ \sigma \in \Gamma(\$) \mid D\sigma = 0 \}$$

$[-, -]$

$$\mathfrak{h}_{\bar{0}} \times \mathfrak{h}_{\bar{0}} \rightarrow \mathfrak{h}_{\bar{0}}$$

Lie bracket of vector fields

$$\mathfrak{h}_{\bar{0}} \times \mathfrak{h}_{\bar{1}} \rightarrow \mathfrak{h}_{\bar{1}}$$

spinorial lie derivative

$$\mathfrak{h}_{\bar{1}} \times \mathfrak{h}_{\bar{1}} \rightarrow \mathfrak{h}_{\bar{0}}$$

Dirac current

Jacobi

$$(\mathfrak{h}_{\bar{0}}, \mathfrak{h}_{\bar{0}}, \mathfrak{h}_{\bar{0}}) \quad \checkmark$$

Jacobi of lie bracket of VFs

$$(\mathfrak{h}_{\bar{0}}, \mathfrak{h}_{\bar{0}}, \mathfrak{h}_{\bar{1}}) \quad \checkmark$$

$$[\mathcal{L}_\xi, \mathcal{L}_\eta] = \mathcal{L}_{[\xi, \eta]}$$

$$(\mathfrak{h}_{\bar{0}}, \mathfrak{h}_{\bar{1}}, \mathfrak{h}_{\bar{1}}) \quad \checkmark$$

Equivariance of Dirac current

$$(\mathfrak{h}_{\bar{1}}, \mathfrak{h}_{\bar{1}}, \mathfrak{h}_{\bar{1}}) \quad ?$$

Needs to be checked

If all Jacobis are satisfied, then

1. \mathfrak{h} is a lie superalgebra
2. \mathfrak{h} is filtered
3. $\text{gr}(\mathfrak{h}) < \square$

e: Killing superalgebras are filtered deformations of graded subalgebras of the Poincaré superalgebra

Infinitesimal deformations computable from the generalised Spencer cohomology

$$H^{2,2}(\mathfrak{g}_-; \mathfrak{g}) \quad \mathfrak{g} < \square$$

Part III

Applications & Calculations

Supersymmetric supergravity backgrounds

Fact

Every supergravity theory gives D on S and every background has an associated Killing superalgebra \mathfrak{h} .

Example

$D=11$ Supergravity [Nahm, Cremmer+Julia+Scherk '78]

$$D_X \sigma = \nabla_X \sigma + \frac{1}{8} F \cdot X \cdot \sigma + \frac{1}{24} X \cdot F \cdot \sigma$$

$$\begin{aligned} F &\in \Omega^4(M) \\ dF &= 0 \end{aligned}$$

(Bosonic) field equations \Leftrightarrow vanishing of Clifford trace of curvature of D

Theorem

[JMF + Meessen + Philip '09]

$$(M, g) \quad \text{11-dimensional spin} \quad F \in \Omega^4(M) \quad dF = 0$$

$$k = k_{\bar{0}} \oplus k_{\bar{1}}$$
$$k_{\bar{0}} = \{ \xi \text{ } kv \mid \alpha \xi F = 0 \}$$
$$k_{\bar{1}} = \{ \sigma \in \Gamma(\$) \mid D\sigma = 0 \}$$

is a Killing superalgebra.

Theorem

[JMF + Hustler '13]

$$\dim k_{\bar{1}} > \frac{1}{2} \text{ rank } \$ \quad \xrightarrow{\text{(sharp)}}$$

$[k_{\bar{1}}, k_{\bar{1}}] \subset k_{\bar{0}}$ acts locally
transitively on M

A cohomological rederivation of D=11 SUGRA?

$$\mathfrak{g} = \square = \underline{\text{so}}(V) \oplus S \oplus V$$

(V, η) 11-dim'l Lorentzian
 S real, 32-dim'l

Theorem

[JMF + Santi '16]

$$H^{2,2}(\mathfrak{g}_-; \mathfrak{g}) \cong \wedge^4 V$$

($\underline{\text{so}}(V)$ -mod)

$$[\beta + \gamma] \leftrightarrow F$$

$$\beta: V \rightarrow \text{End}(S)$$

$$\gamma: \Theta^2 S \rightarrow \underline{\text{so}}(V)$$

$$\beta_X s = \frac{1}{8} F \cdot X \cdot s + \frac{1}{24} X \cdot F \cdot s$$

Fables of the Reconstruction

The Killing superalgebra (when $\dim \mathfrak{h}_T > 16$) contains all the information to reconstruct (M, g, F) ... locally.

Theorem

[JMF + Santi '18]

If $\dim \mathfrak{h}_T > \frac{1}{2} \text{rank } \$$ ($= 16$), one can reconstruct a solution (M, g, F) of the supergravity field equations (up to local isometry) with Killing superalgebra \mathfrak{h} .

This allows a purely algebraic reformulation of the classification problem of " $>\frac{1}{2}$ -BPS" supergravity backgrounds.

Supersymmetry in curved space

A Killing superalgebra $\mathfrak{h} = \{\text{D-parallel sections of } E \oplus \$\}$ defines a supersymmetry algebra on (M, g) .

⇒ (rigidly) supersymmetric field theories on (M, g) .

These are key ingredients in the application of supersymmetric localisation to computing exact quantities in QFT.

Spencer cohomology $H^{2,2}(g_-; g)$ for $g = \square$ (in diverse dimensions, possibly extended) suggests connections D on $\$$ whose parallel spinors generate KSAs.

Supergauges

off-shell SUGRA

"conformal" SUGRA

Some calculations

$$D = 11$$

$$\wedge^4 V$$

[JMF + Saetti '16]

$$D = 10 \quad (I)$$

$$\wedge^3 V$$

[de Medeiros + JMF]

$$D = 6 \quad (1,0)$$

$$\wedge^3_+ V \oplus (V \otimes \mathfrak{sp}(1))$$

[de Medeiros + JMF + Saetti '18]

$$D = 6 \quad (1,0) + R$$

$$\wedge^3 V \oplus (V \otimes \mathfrak{sp}(1))$$

[de Medeiros + JMF + Saetti '18]

$$D = 5$$

$$\wedge^2 V \oplus (V \otimes \mathfrak{sp}(1))$$

[Bechtel + JMF '20]

$$D = 4 \quad N = 1$$

$$\wedge^\circ V \oplus V \oplus \wedge^4 V$$

[de Medeiros + JMF + Saetti '16]

Future & work in progress

Non-lorentzian : galilean + Bargmann ($D=3$)

Kaluza-Klein reduction

Extended $D=4$, $D=3$

Superconformal $D=3$

Exotic $D=6$

