# Plane wave limits in gauge theory

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Recall: ['t Hooft (1974)]

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$$\sum_{g=0}^{\infty} N^{2-2g} f_g(\lambda)$$

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Which string theory?



$$S_{
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in the low energy limit decouples into

IIB supergravity in bulk

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IIB supergravity in bulk + IIB supergravity near brane horizon

 $\mathcal{N}{=}4$   $\mathrm{SU}(N)$  Yang-Mills at  $g_{\mathrm{YM}}$  on  $\mathbb{R} \times S^3$ 

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• *spectrum* of, say, single trace operators; e.g., chiral primaries (in short multiplets)

$$\operatorname{Tr}' \phi \otimes \cdots \otimes \phi$$

But...

But...further checks are hindered by lack of clear dictionary between *massive* string modes

unprotected single trace operators

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However...recent progress in this direction has been made using a different large N limit, the so-called plane wave limit.

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Boost along  $\gamma$  and zoom in on  $\gamma$  while rescaling g

$$g = 2dUdV + AdV^2 + B_i dY^i dV + C_{ij} dY^i dY^j$$

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rescale

$$g_{\mathrm{plane\ wave}} = \lim_{\Omega \to 0} \Omega^{-2} g(\Omega)$$

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• "lightlike gauge"

$$i_{\partial/\partial U}A^{(p)} = 0$$

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**Note**: field strength is *null*:

$$F_{\text{plane wave}} = dA_{\text{plane wave}} = dx^+ \wedge \Theta(x^+)$$

diffeomorphisms

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It is useful in classifying the different plane wave limits of a given background.

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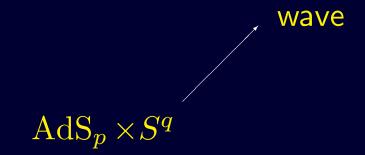
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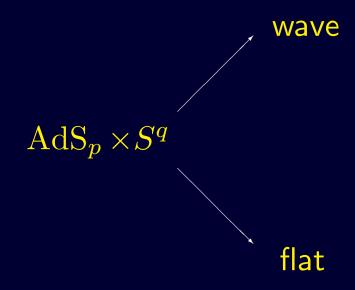
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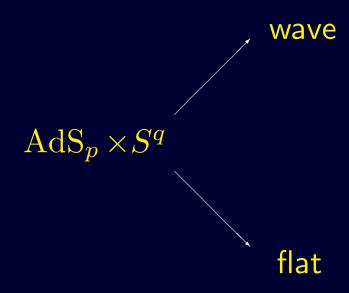
PWL : Vacua → Vacua

Vacua of the form  $\mathrm{AdS}_p \times S^q$ 

$$AdS_p \times S^q$$



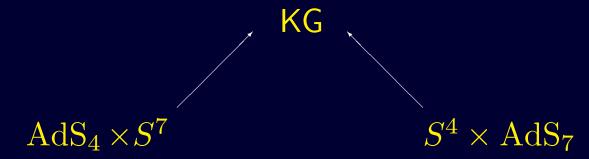


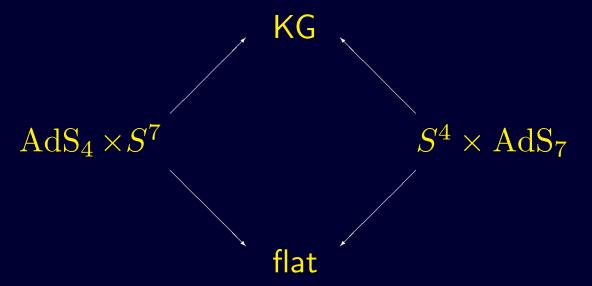


depending on whether or not the component of  $\dot{\gamma}$  tangent to  $S^q$  vanishes. [Blau-FO-Hull-Papadopoulos (2002)]

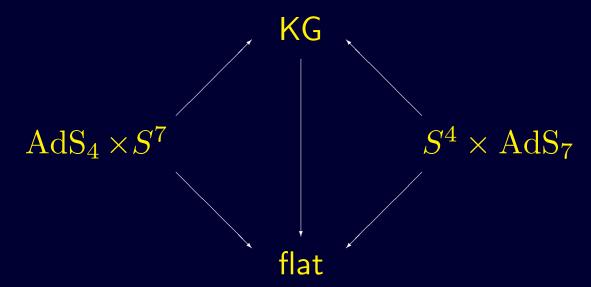
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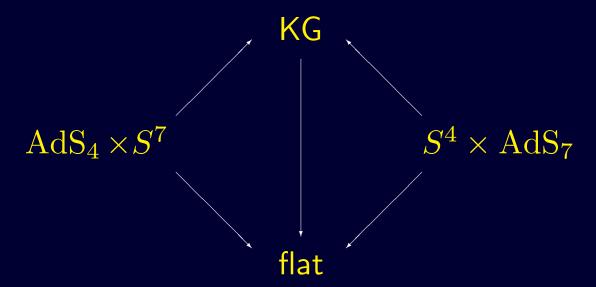


Eleven-dimensional supergravity vacua:



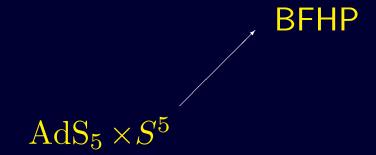
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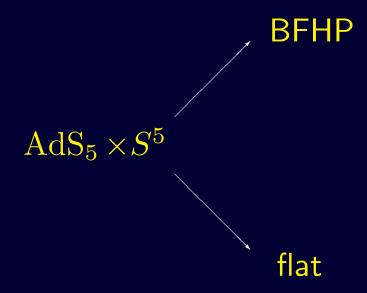
[FO-Papadopoulos (2002)]

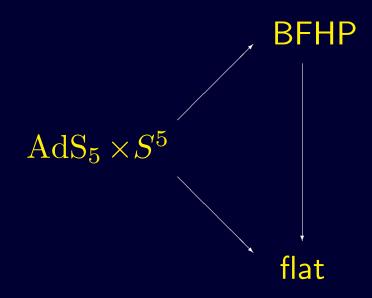


KG: 11-dimensional symmetric plane wave [Kowalski-Glikman (1984)]

$$AdS_5 \times S^5$$

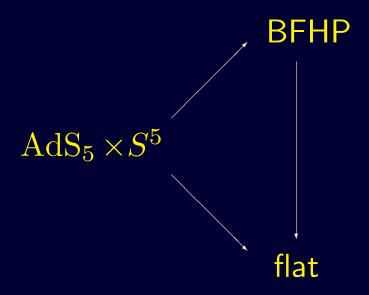






In ten-dimensional IIB supergravity:

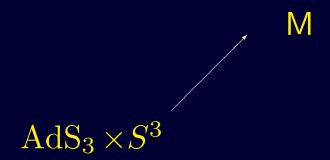
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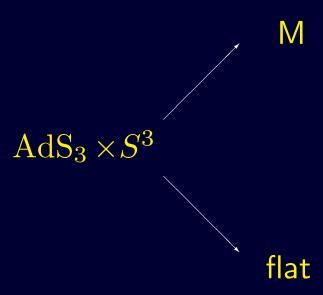


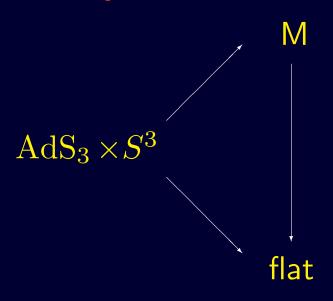
BFHP: 10-dimensional symmetric plane wave

[Blau-FO-Hull-Papadopoulos (2001)]

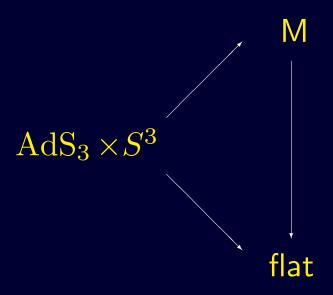
$$AdS_3 \times S^3$$







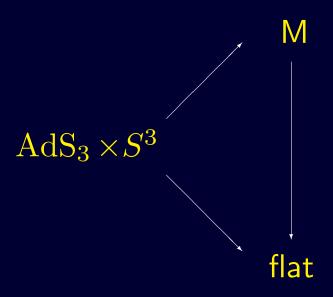
[Chamseddine–FO–Sabra (to appear in 2003)]



M: 6-dimensional symmetric plane wave

[Meessen (2001)]

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All solutions are Lie groups and all plane wave limits are group contractions.

[Stanciu-FO (2003)]

Symmetric plane waves

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where  $A_{ij}$  is *constant*, and is determined by its eigenvalues, up to order and scale.

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- $\mu$  inessential, but  $\mu \to 0$  recovers flat solution

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- $\bullet$   $\partial/\partial x^-$  is a parallel null vector  $\Longrightarrow$  natural light-cone gauge
- gauge-fixed Green–Schwarz action is *quadratic* [Metsaev (2001)]

$$\frac{1}{2\pi\alpha'} \int d\tau \int_{0}^{2\pi\alpha'p^{+}} d\sigma \left[ \frac{1}{2}\dot{x}^{2} - \frac{1}{2}x'^{2} - \frac{1}{2}\mu^{2}x^{2} + i\bar{\psi}(\partial \!\!\!/ + \mu M)\psi \right]$$

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- $(\dot{\gamma}_{AdS} \dot{\gamma}_S)/R^2 \to \partial/\partial x^- \implies p^+ \sim (\Delta + J)/R^2$

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**Notice**: As  $N \to \infty$  we focus on states with larger and larger J. Thus observables are not held fixed in this limit.

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Now...

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- they are dual to the free string excitations on the plane wave background

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BMN operators are single trace operators consisting in a string of J Z's and a finite number of *impurities*.

$$\frac{1}{\sqrt{(N/2)^{J+1}}} \operatorname{Tr} Z^J \longleftrightarrow |p^+\rangle$$

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etc...

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How about interactions?

# **Interacting string theory**

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### **Interacting string theory**

A detailed study of BMN correlators shows:

[Kristjansen et al., Gross et al., Constable et al. (2002)]

• theory develops a different effective coupling constant

$$\lambda' = \frac{g_{\text{YM}}^2 N}{J^2} = \frac{1}{(\mu p^+ \alpha')^2}$$

• a different genus-counting parameter

$$g_2^2 = \left(\frac{J^2}{N}\right)^2 = 16\pi^2 g_s^2 (\mu p^+ \alpha')^4$$

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string field theory

[..., Chu-Khoze, Gomis et al. (2003)]

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• the correspondence has been extended to theories with less supersymmetry; but as usual QCD remains elusive.

Thank you.