Classification results on supergravity vacua

José Figueroa-O'Farrill Edinburgh Mathematical Physics Group School of Mathematics



Based on work in collaboration with

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George Papadopoulos (King's College, London)

- * hep-th/0211089 (JHEP 03 (2003) 048)
- * math.AG/0211170

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• Ali Chamseddine and Wafic Sabra (CAMS, Beirut)

 \star in preparation

A geometric motivation

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These are in one-to-one correspondence with parallel sections of the bundle

 $\mathcal{E}(M) = TM \oplus \Lambda^2 T^* M$

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AdS_n
$$\subset \mathbb{E}^{2,n-1}$$
: $-t_1^2 - t_2^2 + x_1^2 + \dots + x_{n-1}^2 = \frac{-1}{\kappa^2}$

Note: the $\kappa \neq 0$ spaces are *quadrics* in a flat space in one dimension higher

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Only the spherical case is solved (culminating in the work of Wolf in the 1970s), but there are many partial results in other cases.

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In other words, we will classify vacua up to local isometry.

Supergravities

| | 32 | | | | 24 | | 20 | 16 | | 12 | 8 | 4 |
|----|-------|-------|-------|--------|-------|--------|-------|--------|-------|-------|-------|-------|
| 11 | М | М | | | | | | | | | | |
| 10 | IIA | IIB | | | | | | 1 | | | | |
| 9 | N=2 | | | | | | | N = 1 | | | | |
| 8 | N=2 | | - | | - | | | N = 1 | - | | | |
| 7 | N = 4 | | | | | 1 | | N=2 | | | | |
| 6 | (2,2) | | (3,1) | (4, 0) | (2,1) | (3, 0) | | (1, 1) | (2,0) | | (1,0) | |
| 5 | | N=8 | | | N = 6 | | | N=4 | | | N=2 | |
| 4 | | N = 8 | | | N = 6 | | N = 5 | N = 4 | | N = 3 | N=2 | N = 1 |

[Van Proeyen, hep-th/0301005]

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Any theory in the table can be dimensionally reduced down its column : symmetric backgrounds can be reduced à la Kaluza–Klein

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To classify vacua, one can therefore

• classify vacua of theories at the top of each column, and

• investigate their possible Kaluza–Klein reductions.

Let (M, g, Φ, S) be a supergravity background:

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- fermions have been put to zero

(M, g, Φ, S) is supersymmetric if it admits Killing spinors

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defined by the supersymmetric variation of the gravitino:

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[Chamseddine–FO–Sabra]

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- $D = 11 \,\,\mathrm{M}$

[FO–Papadopoulos]

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• spinors are Majorana; that is, associated to one of the two irreducible real 32-dimensional representations of $C\ell(1, 10)$. Therefore the gravitino also has 128 physical degrees of freedom.

$$D_{\mu} = \nabla_{\mu} - \frac{1}{288} F_{\nu\rho\sigma\tau} \left(\Gamma^{\nu\rho\sigma\tau}{}_{\mu} + 8\Gamma^{\nu\rho\sigma}\delta^{\tau}_{\mu} \right)$$

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where *I* is an index labeling the following elements

$$\Gamma_a \quad \Gamma_{ab} \quad \Gamma_{abc} \quad \Gamma_{abcd} \quad \Gamma_{abcde}$$

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with T quadratic in F. This means that $R_{\mu\nu\rho\sigma}$ is parallel; equivalently, that g is locally symmetric.

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- if the plane is lorentzian, we can choose a pseudo-orthonormal frame in which the only nonzero component of F is F_{0123}
- if the plane is null, we can choose a pseudo-orthonormal frame in which the only nonzero component of F is F_{-123}

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• F lorentzian: a one parameter R < 0 family of vacua

 $AdS_4(8R) \times S^7(-7R)$ $F = \sqrt{-6R} dvol(AdS_4)$

$$g = 2dx^{+}dx^{-} - \frac{1}{36}\mu^{2} \left(4\sum_{i=1}^{3} (x^{i})^{2} + \sum_{i=4}^{9} (x^{i})^{2}\right) (dx^{-})^{2} + \sum_{i=1}^{9} (dx^{i})^{2}$$

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All vacua embed isometrically in $\mathbb{E}^{2,11}$ as the intersections of two quadrics.

Vacua are related by Penrose limits

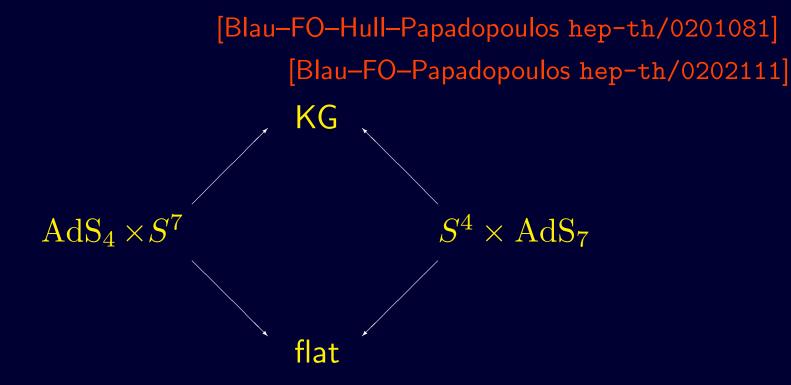
[Blau-FO-Hull-Papadopoulos hep-th/0201081]

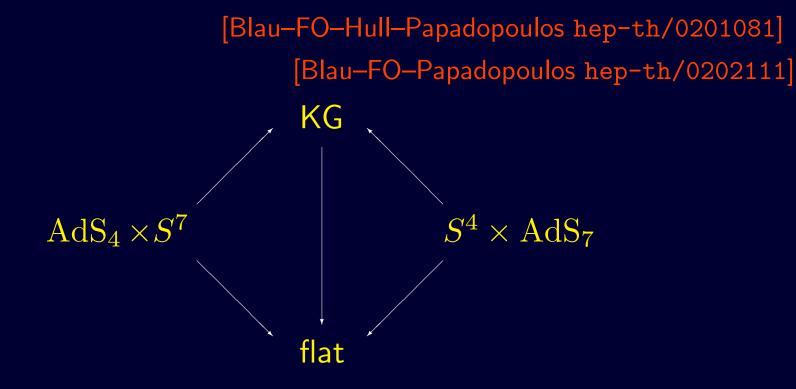
[Blau-FO-Papadopoulos hep-th/0202111]



 $S^4 \times \overline{\mathrm{AdS}_7}$

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[Back]

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• spinors are positive-chirality symplectic Majorana–Weyl; i.e., associated to the 8-dimensional real representation of $\frac{\text{Spin}(1,5) \times \text{Sp}(1)}{\text{Sp}(1)}$ having positive six-dimensional chirality.

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- \star anti-selfdual closed 3-form F

for a total of 9 + 3 = 12 physical bosonic degrees of freedom

• spinors are positive-chirality symplectic Majorana–Weyl; i.e., associated to the 8-dimensional real representation of $Spin(1,5) \times Sp(1)$ having positive six-dimensional chirality. The gravitino has therefore also 12 physical degrees of freedom.

$$D_{\mu} = \nabla_{\mu} + \frac{1}{8} F_{\mu}{}^{ab} \Gamma_{ab}$$

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Maximal supersymmetry $\implies D$ is flat.

A pseudoriemannian manifold admitting a flat metric connection with torsion

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The solution to this problem is known.

Which Lie algebras have an invariant metric?

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But there is a more general construction.

• g a Lie algebra with an invariant metric

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- h a Lie algebra acting on g via *antisymmetric derivations*

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• since \mathfrak{h} preserves the metric on \mathfrak{g} , there is a linear map

 $\mathfrak{h} \to \Lambda^2 \mathfrak{g}$

$$\omega:\Lambda^2\mathfrak{g} o\mathfrak{h}^*$$

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relative to bases X_a , H_i and H^i for \mathfrak{g} , \mathfrak{h} and \mathfrak{h}^* , respectively.

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• \mathfrak{h} acts on $\mathfrak{g} \times_{\omega} \mathfrak{h}^*$ preserving the Lie bracket, so we can form the *double extension*

 $\mathfrak{d}(\mathfrak{g},\mathfrak{h})=\mathfrak{h}\ltimes(\mathfrak{g} imes_\omega\mathfrak{h}^*)$

$$egin{array}{cccc} X_b & H_j & H^j \ X_a & \left(egin{array}{cccc} g_{ab} & 0 & 0 \ 0 & B_{ij} & \delta^j_i \ 0 & \delta^i_j & 0 \end{array}
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This construction is due to Medina and Revoy who proved an important structure theorem.

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The structure theorem of Medina and Revoy

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of indecomposables.

[See also FO-Stanciu hep-th/9506152]

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where \mathfrak{a} is abelian with euclidean metric and \mathfrak{h} is one-dimensional. (Any semisimple factors in \mathfrak{a} factor out of the double extension. [FO-Stanciu hep-th/9402035])



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- ∂(ℝ⁴, ℝ), actually a family of Lie algebras parametrised by homomorphisms

$$\mathbb{R} \to \Lambda^2 \mathbb{R}^4 \cong \mathfrak{so}(4)$$



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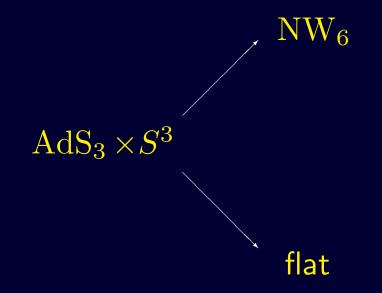
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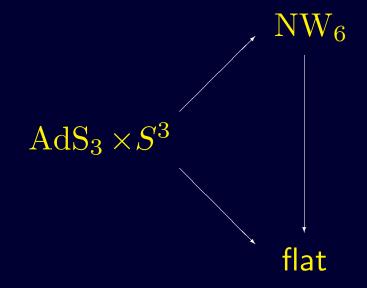
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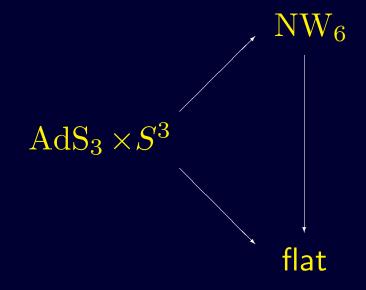
The third case is a six-dimensional version of the Nappi-Witten spacetime, NW_6 , discovered by Meessen. [Meessen hep-th/0111031]

$AdS_3 \times S^3$

 NW_6 $\mathrm{AdS}_3 \times S^3$







which in this case are *group contractions* à la Inönü–Wigner. [Stanciu–FO hep-th/0303212]

[Back]

• bosonic fields:

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where we have written the two real spinors as a complex spinor, and α depends on the constant value of τ .

Expanding the curvature of D into antisymmetric products of Γ -matrices

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This equation defines a generalisation of a Lie algebra known as a4-Lie algebra (with an invariant metric).[Filippov (1985)]

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(Unfortunate notation: a 2-Lie algebra is a Lie algebra.)

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A Lie algebra is a vector space **g** together with an antisymmetric *bilinear* map

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 $[]: \Lambda^2 \mathfrak{g} \to \mathfrak{g}$

is a derivation over []; that is,

 $ad_X[Y, Z] = [ad_X Y, Z] + [Y, ad_X Z]$

An *n-Lie algebra*

An *n-Lie algebra* is a vector space n together with an antisymmetric *n*-linear map

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[FO-Papadopoulos math.AG/0211170]

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Plugging these expressions back into the relation between the

curvature tensor to F and g

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Notice that *g* is a bi-invariant metric on a Lie group: a ten-dimensional version of the Nappi–Witten spacetime. [Stanciu–FO hep-th/0303212] These vacua again embed isometrically in $\mathbb{E}^{2,10}$ as intersections of quadrics

[Blau-FO-Hull-Papadopoulos hep-th/0201081]

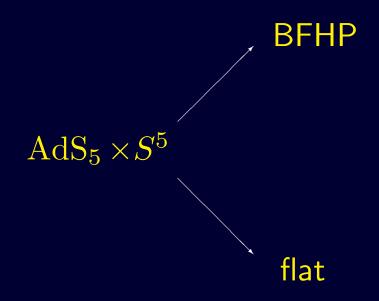
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[Blau-FO-Hull-Papadopoulos hep-th/0201081]

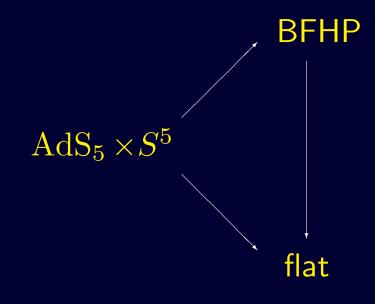
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- D = 6 (2,0) supergravity: all (1,0) vacua are also vacua of (2,0) and early indications show that there are no others. (1,0) vacua do have reductions preserving all supersymmetry.

[Gauntlett-Gutowsky-Hull-Pakis-Reall hep-th/0209114] [Lozano-Tellechea-Meessen-Ortín hep-th/0206200]

Thank you.